

Characterizing Magnetized Turbulence in Molecular Clouds (and Galaxies)

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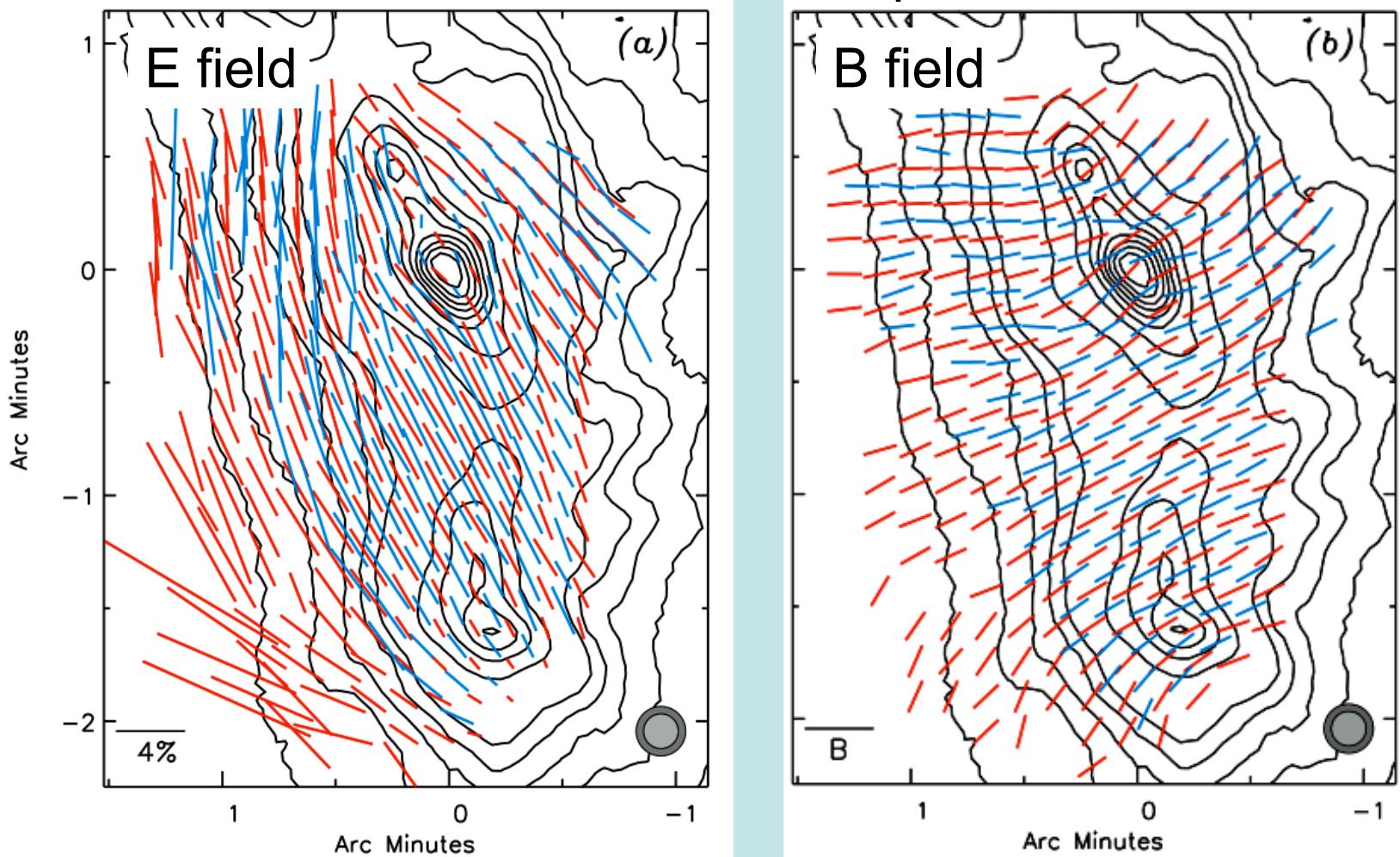
C. Darren Dowell (Caltech/NASA JPL)

Outline

- Dispersion of magnetic fields
 - Separation of turbulent and large-scale fields through structure functions
 - Example: the Chandrasekhar-Fermi technique
- Application/results
 - Single-dish - OMC-1, CSO/SHARP
 - Turbulence correlation length
 - Turbulent/ordered field energy ratio (CF equation)
 - Interferometry - SMA
 - Magnetized turbulent power spectrum
 - Ambipolar diffusion scale
 - Single-dish (Effelsberg) + Interferometry (VLA)
 - M51 - Anisotropic turbulence

Polarization Maps - what are they good for?

OMC-1 - SHARP, 350 and 450 μ m



Structure Functions

- Common for studying turbulence
 - Nice properties for power-law power spectra with stationary signals
- Have been used in astrophysics for some time
 - Molecular clouds
 - Dotson (1996, ApJ, 470, 566) → M17 SW with KAO at 100 μm (polarization angles)
 - Falceta-Gonçalves et al. (2008, ApJ,) → simulations
 - Radio Astronomy
 - Beck et al. (1999) → Intensity maps (Stokes I, Q, and U)

Structure Functions

Given a polarization map

Angle $\Phi(\mathbf{r}) \rightarrow \mathbf{B}$ (plane of the sky)

The Angular Structure Function (stationarity and isotropy)

$$\langle \Delta\Phi^2(\ell) \rangle = \frac{1}{N(\ell)} \sum_{N(\ell) \text{ pairs}} [\Phi(\mathbf{r}) - \Phi(\mathbf{r} + \ell)]^2$$

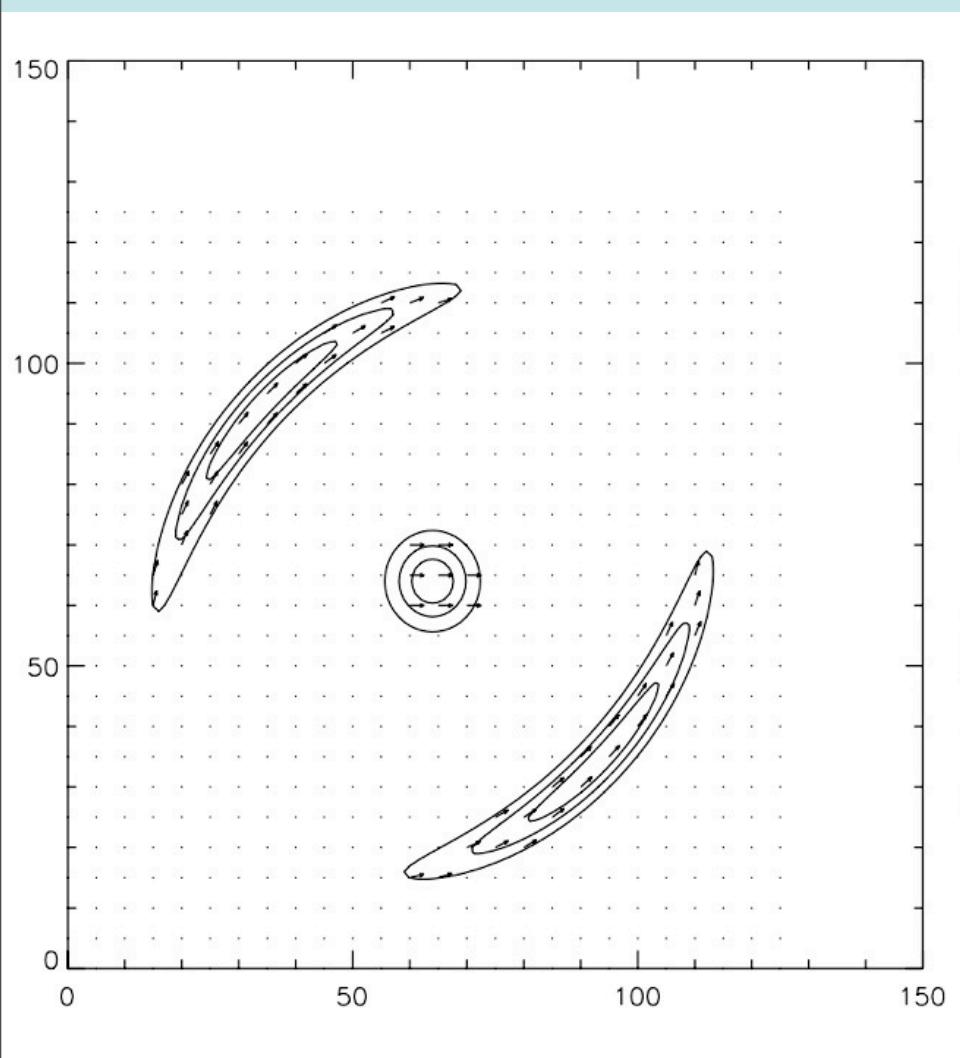
If $\mathbf{B} = \mathbf{B}_t + \mathbf{B}_0$ (turbulent and ordered (large-scale) components)

$$\Rightarrow \langle \Delta\Phi^2(\ell) \rangle = \langle \Delta\Phi_t^2(\ell) \rangle + \langle \Delta\Phi_0^2(\ell) \rangle$$

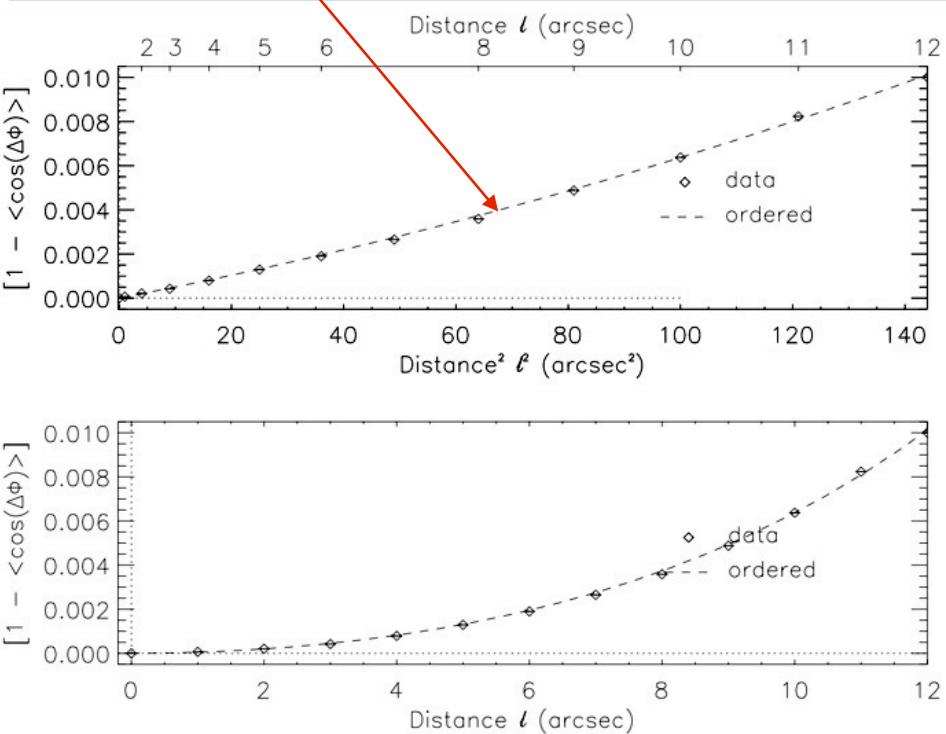
with statistical independence.

$$\Rightarrow 1 - \langle \cos[\Delta\Phi(\ell)] \rangle \simeq \frac{\langle \Delta\Phi^2(\ell) \rangle}{2} \Leftarrow$$

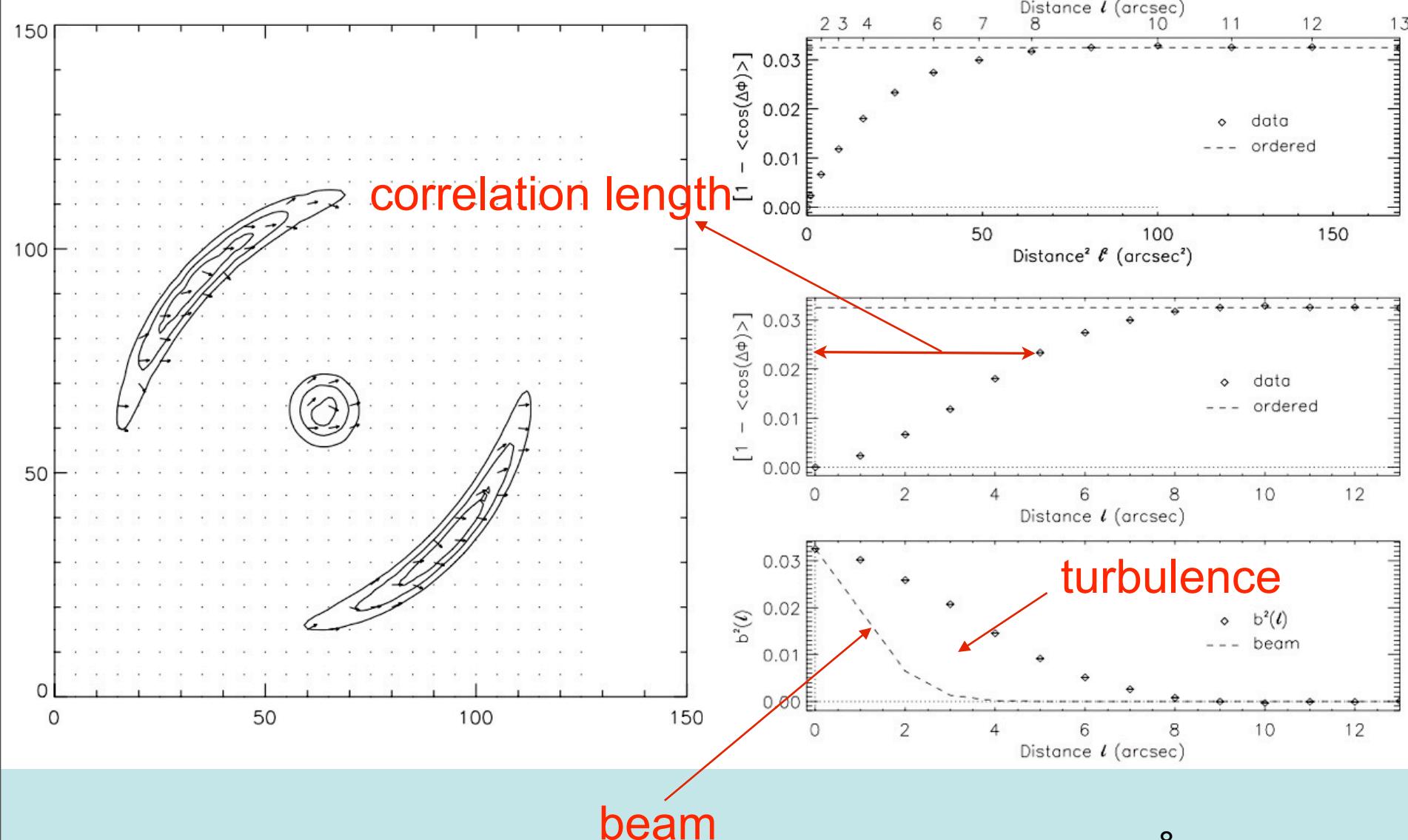
Structure Functions - Large-scale



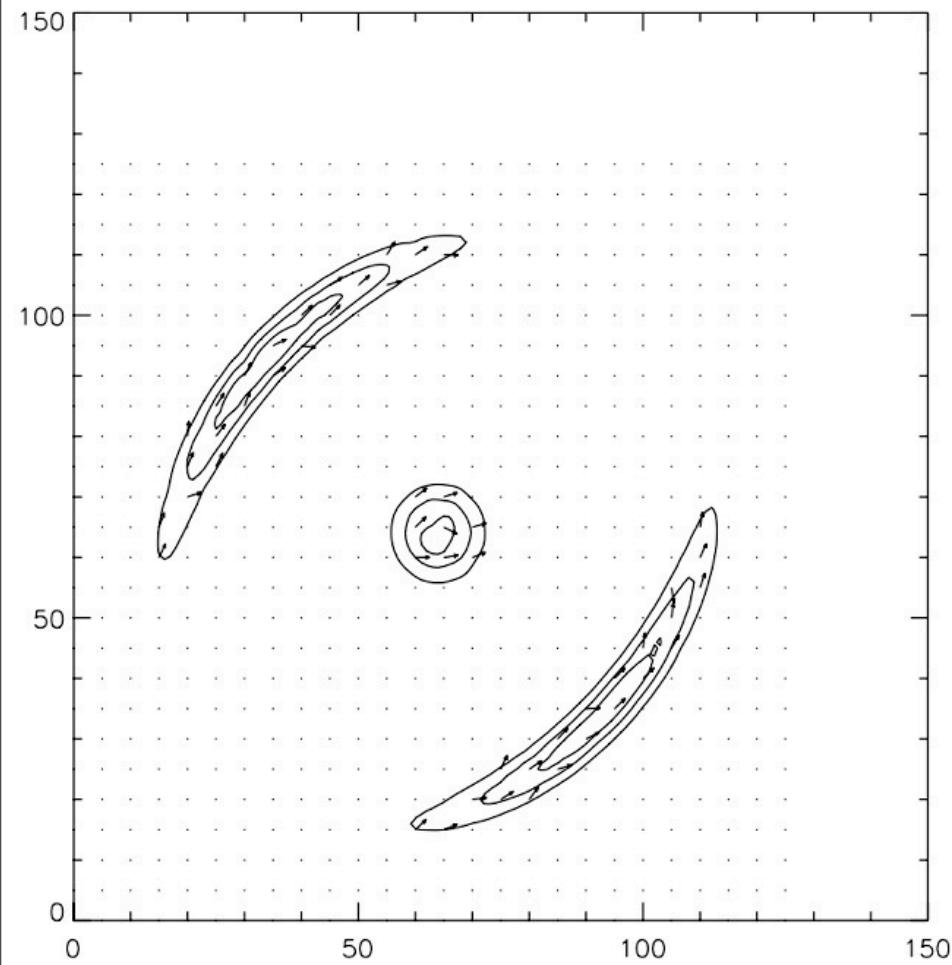
polynomial fit (even powers)



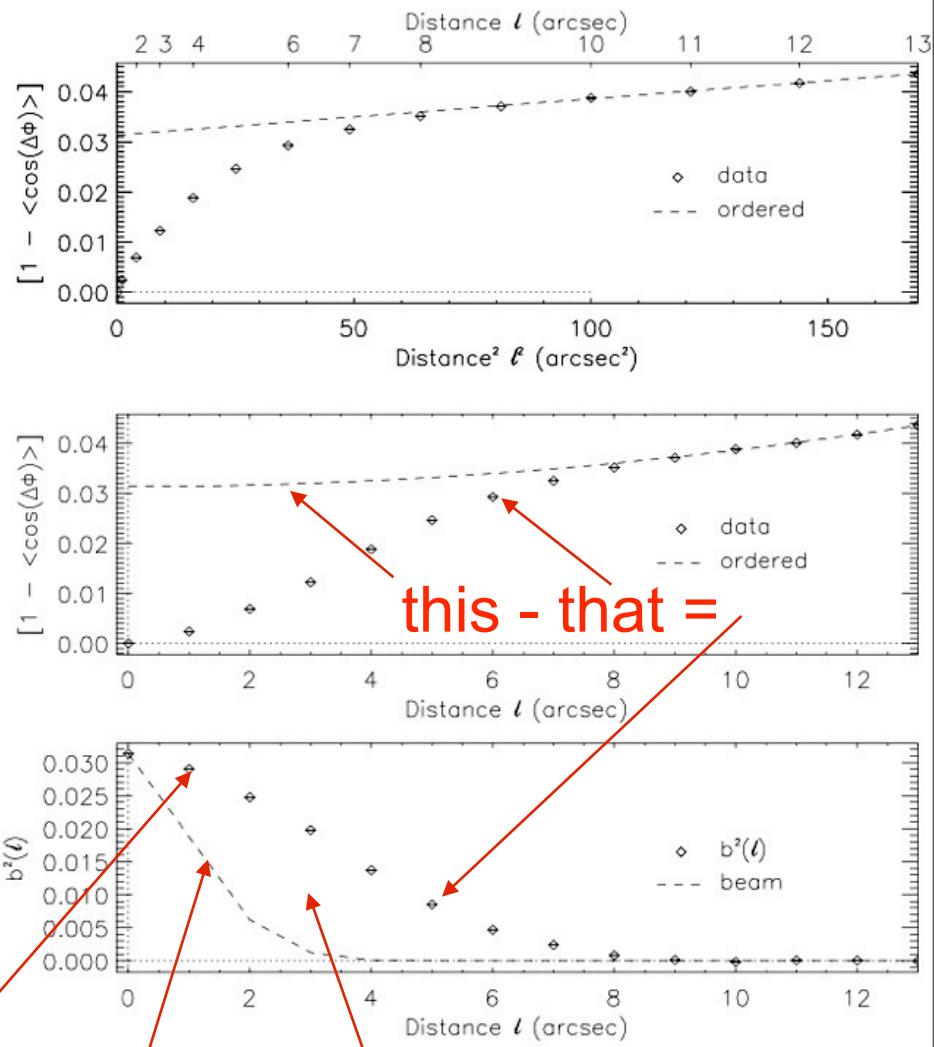
Structure Functions - Turbulence



Structure Functions - Turb.+large-scale



beam-broadened autocorrelation



beam

turbulence

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Example - Chandra-Fermi Equation

$$B_0 \simeq \sqrt{4\pi\rho}\sigma(v) \left[\frac{\langle B_t^2 \rangle}{\langle B_0^2 \rangle} \right]^{-1/2}$$

turbulent

(Chandrasekhar-Fermi 1953)

ordered field

ρ : mass density

$\sigma(v)$: velocity dispersion (one-dimension)

But the angular dispersion $\delta\Phi$ relative to the ordered field determined with polarization maps is

$$\delta\Phi \approx \left[\frac{\langle B_t^2 \rangle}{\langle B_0^2 \rangle} \right]^{1/2}$$

or is it really the case?

Example - Chandra-Fermi Equation

Problems with the CF method

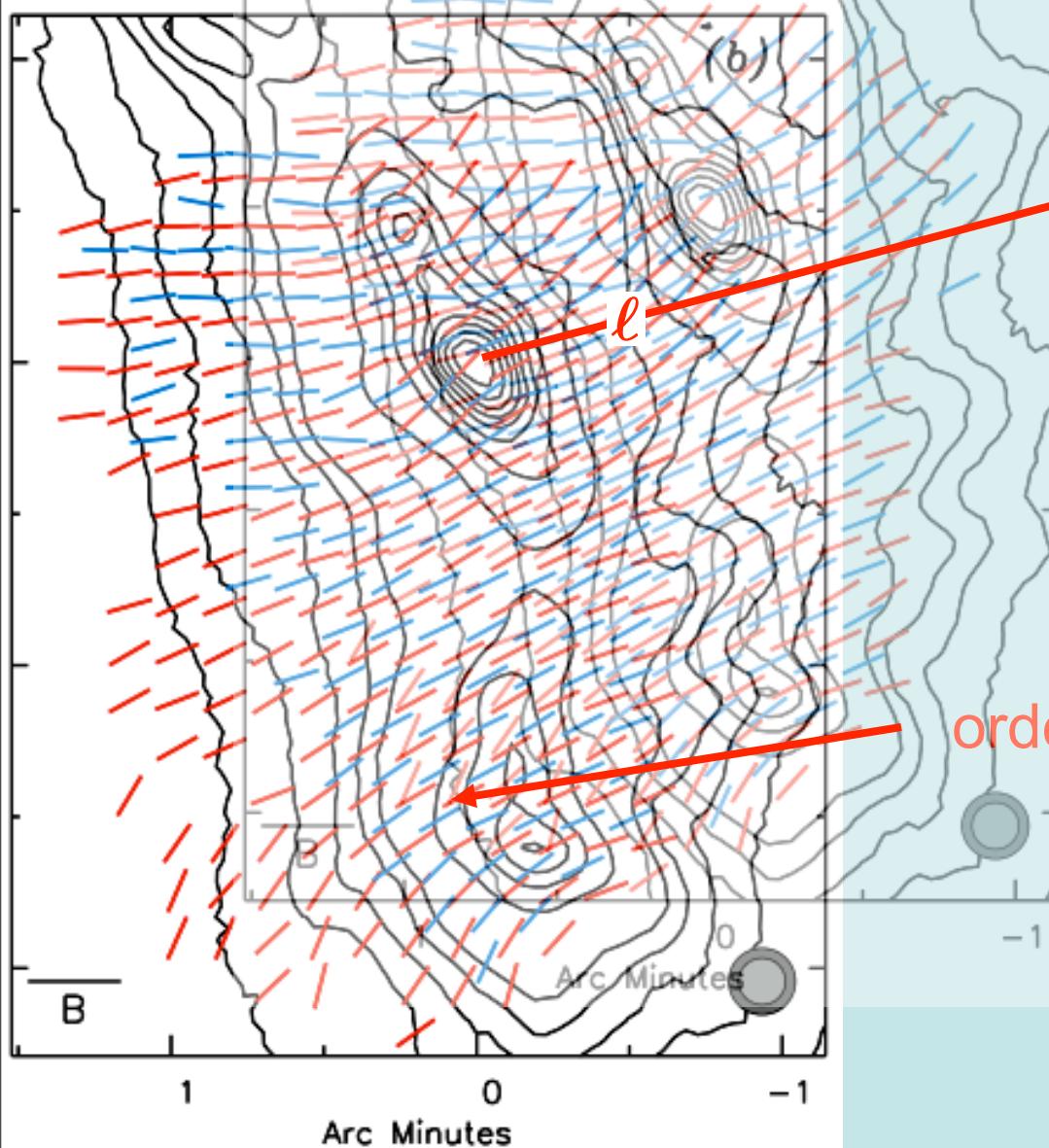
1. The models for \mathbf{B}_0 are imperfect and introduce more errors in the determination of $\delta\Phi$. **This is solved with the structure function.**

Moreover

2. Signal integration along the line of sight and across the telescope beam
 - $\langle \mathbf{B}_t^2 \rangle$ is underestimated due to averaging process
 - \mathbf{B}_0 is therefore overestimated

OMC-1 with SHARP at 350 μ m

OMC-1 - SHARP/CSO, 350 and 450 μ m



$$1 - \langle \cos[\Delta\Phi(\ell)] \rangle \approx \frac{\langle \Delta\Phi^2(\ell) \rangle}{2}$$

ordered + turbulent fields
 $\mathbf{B} = \mathbf{B}_0 + \mathbf{B}_t$

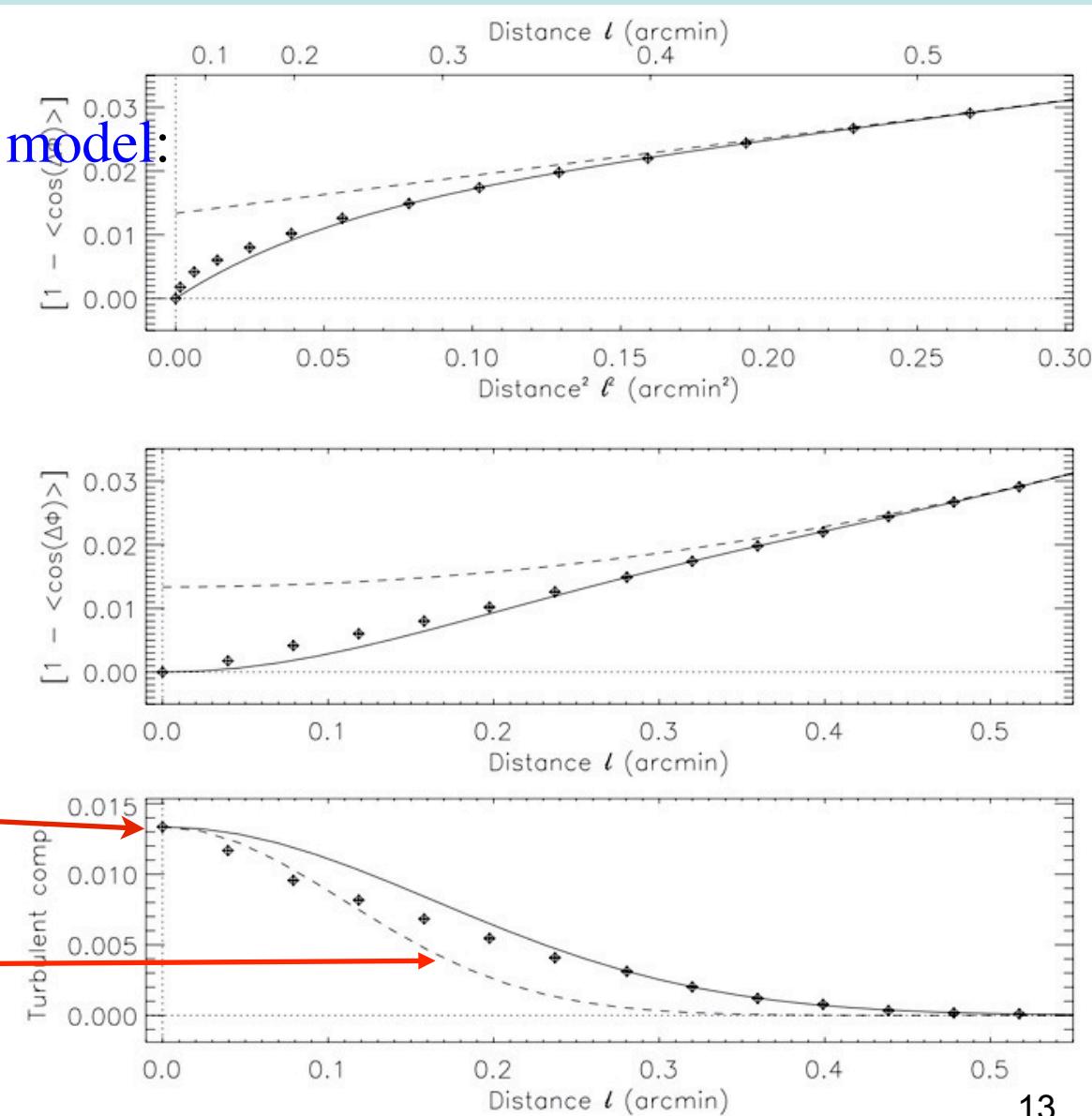
Vaillancourt et al., 2008, ApJ, 679, L25

OMC-1 with SHARP at 350 μm

χ^2 fit to Gaussian model:
 $\delta, \langle \bar{B}_t^2 \rangle / \langle \bar{B}_0^2 \rangle$.

$$(\delta\Phi)^2 \approx \frac{\langle \bar{B}_t^2 \rangle}{\langle \bar{B}_0^2 \rangle}$$

beam



OMC-1 / SHARP - Results

$\delta \simeq 7.3'' = 16$ mpc turbulent correlation length

$$N = \frac{(\delta^2 + 2W^2)\Delta'}{\sqrt{2\pi}\delta^3} \simeq 21 \quad \text{number of turbulent cells}$$

$$\frac{\langle \bar{B}_t^2 \rangle}{\langle \bar{B}_0^2 \rangle} = \frac{1}{N} \frac{\langle B_t^2 \rangle}{\langle B_0^2 \rangle} \simeq 0.013$$

$$\frac{\langle B_t^2 \rangle}{\langle B_0^2 \rangle} \simeq 0.28 \quad \text{turbulent/ordered field energy ratio}$$

with Chandrasekhar-Fermi equation

$$B_0 \simeq \sqrt{4\pi\rho}\sigma(v) \left[\frac{\langle B_t^2 \rangle}{\langle B_0^2 \rangle} \right]^{-1/2} \simeq 760 \mu\text{G} \quad \text{plane of the sky}$$

with $n = 10^5 \text{ cm}^{-3}$, $A = 2.3$, and $\sigma(v) = 1.85 \text{ km s}^{-1}$

Turbulent Power Spectrum

$$1 - \langle \cos[\Delta\Phi(\ell)] \rangle \simeq \frac{\langle \Delta\Phi^2(\ell) \rangle}{2}$$

but

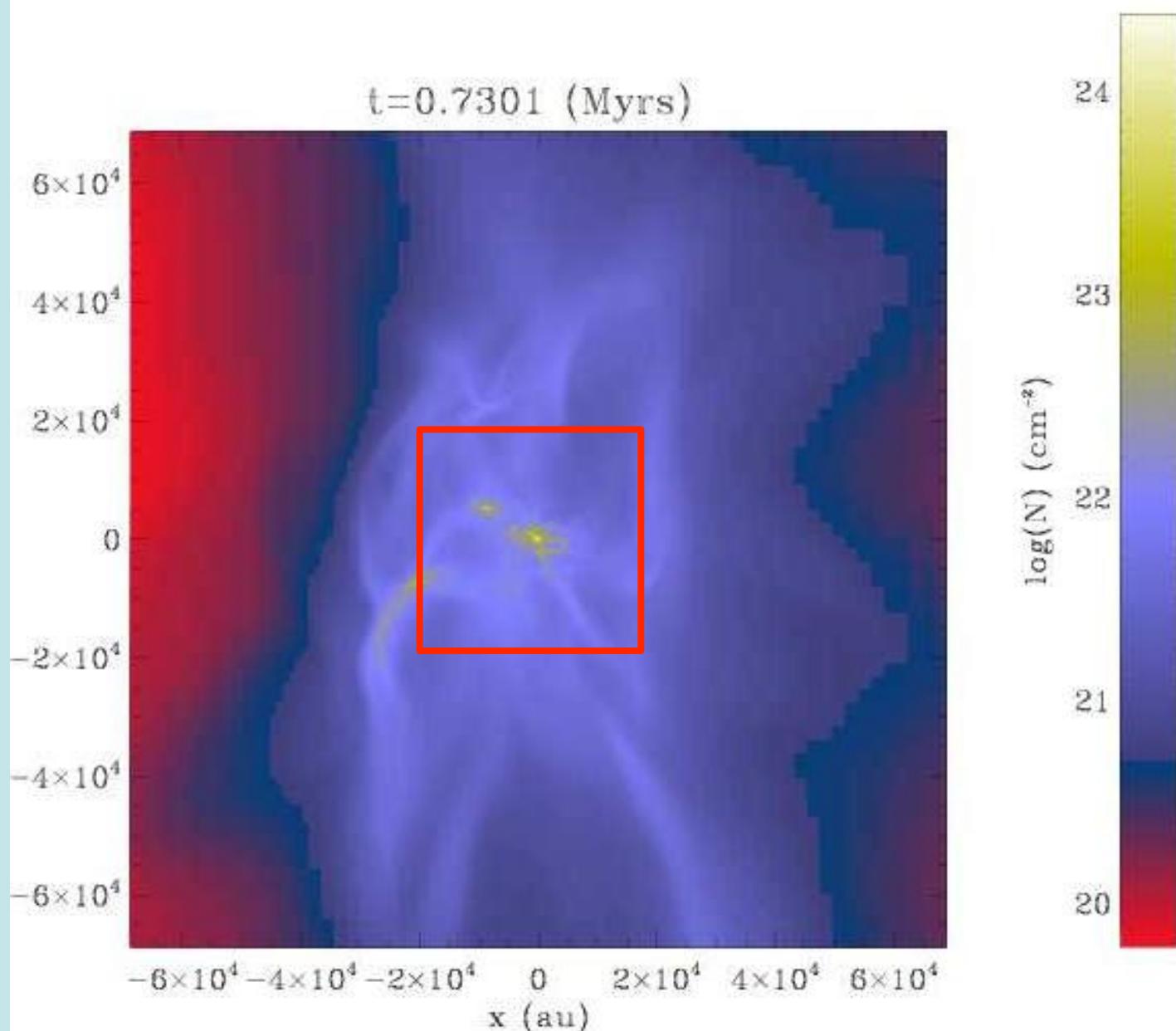
$$\Rightarrow \langle \cos[\Delta\Phi(\ell)] \rangle \equiv \frac{\langle \bar{\mathbf{B}} \cdot \bar{\mathbf{B}}(\ell) \rangle}{\langle \bar{\mathbf{B}} \cdot \bar{\mathbf{B}}(0) \rangle} \leftarrow$$

With a Fourier transform on the turbulent component

$$\frac{\langle \bar{\mathbf{B}} \cdot \bar{\mathbf{B}}(\ell) \rangle}{\langle \bar{B}^2 \rangle} \rightleftharpoons \frac{1}{\langle \bar{B}^2 \rangle} \|H(k_v)\|^2 R_t(k_v) \left[\equiv b^2(k_v) \right]$$

We can determine the turbulent power spectrum $R_t(k_v)$
by deconvolution of the beam $H(k_v)$

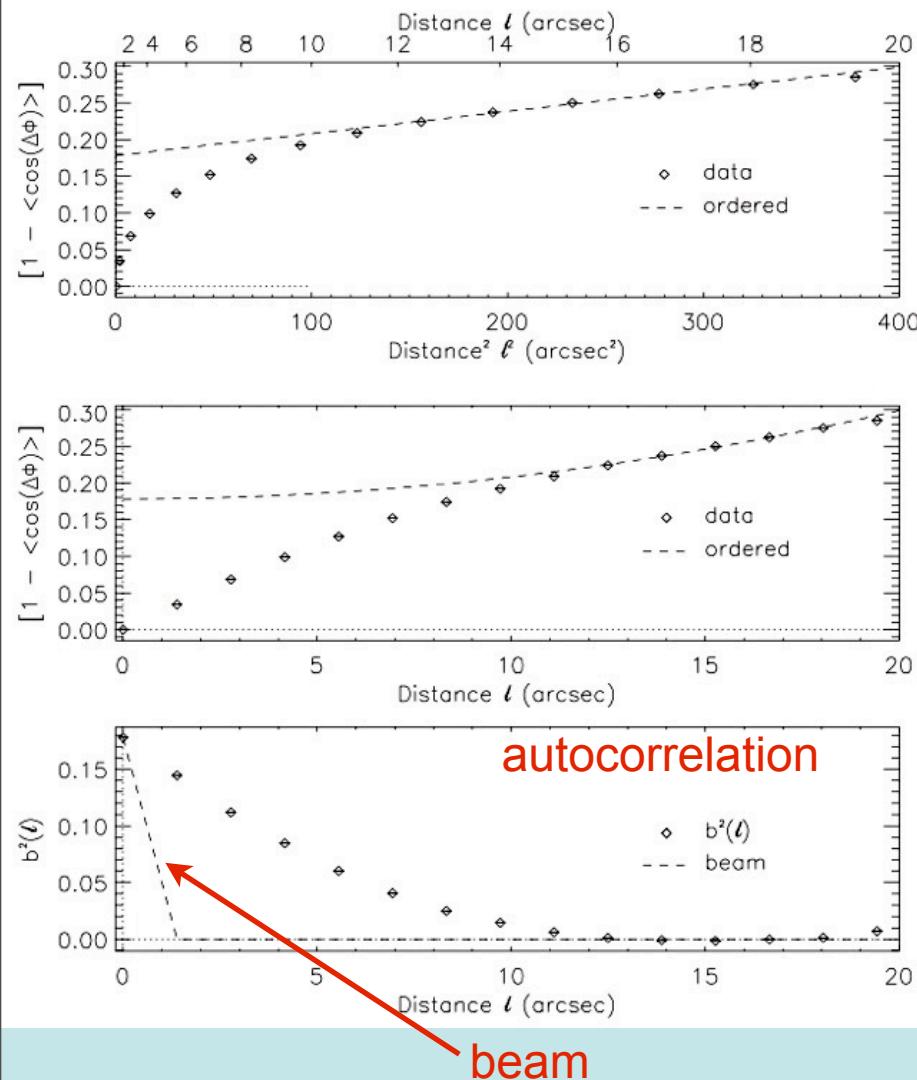
Turbulent Power Spectrum - simulations



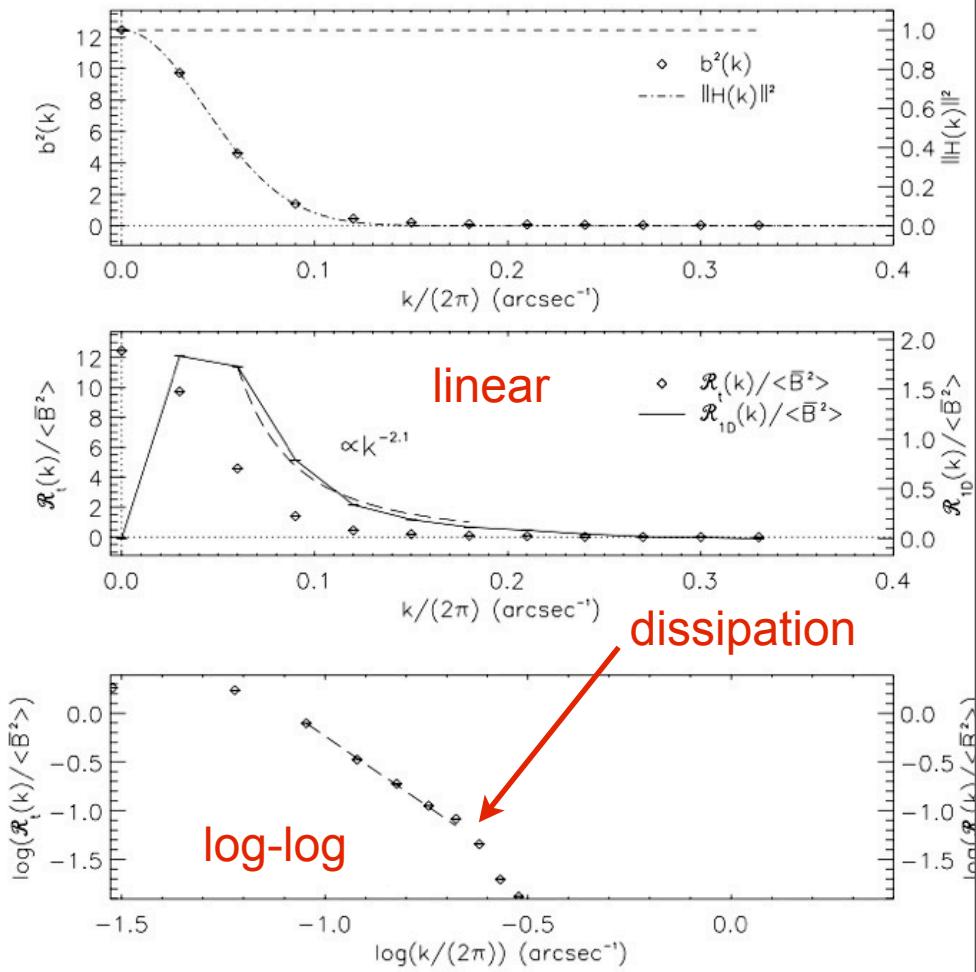
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Turbulent Power Spectrum - simulations

Structure Function

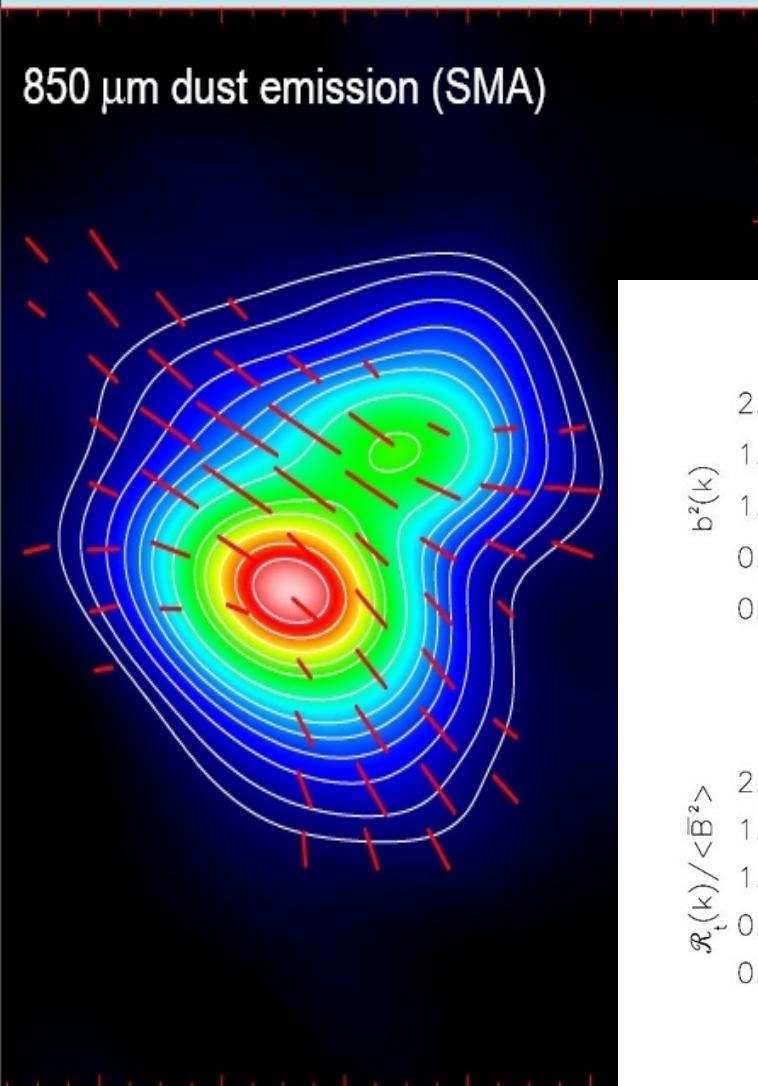


Power Spectrum



Turbulent Power Spectrum - NGC 1333/SMA

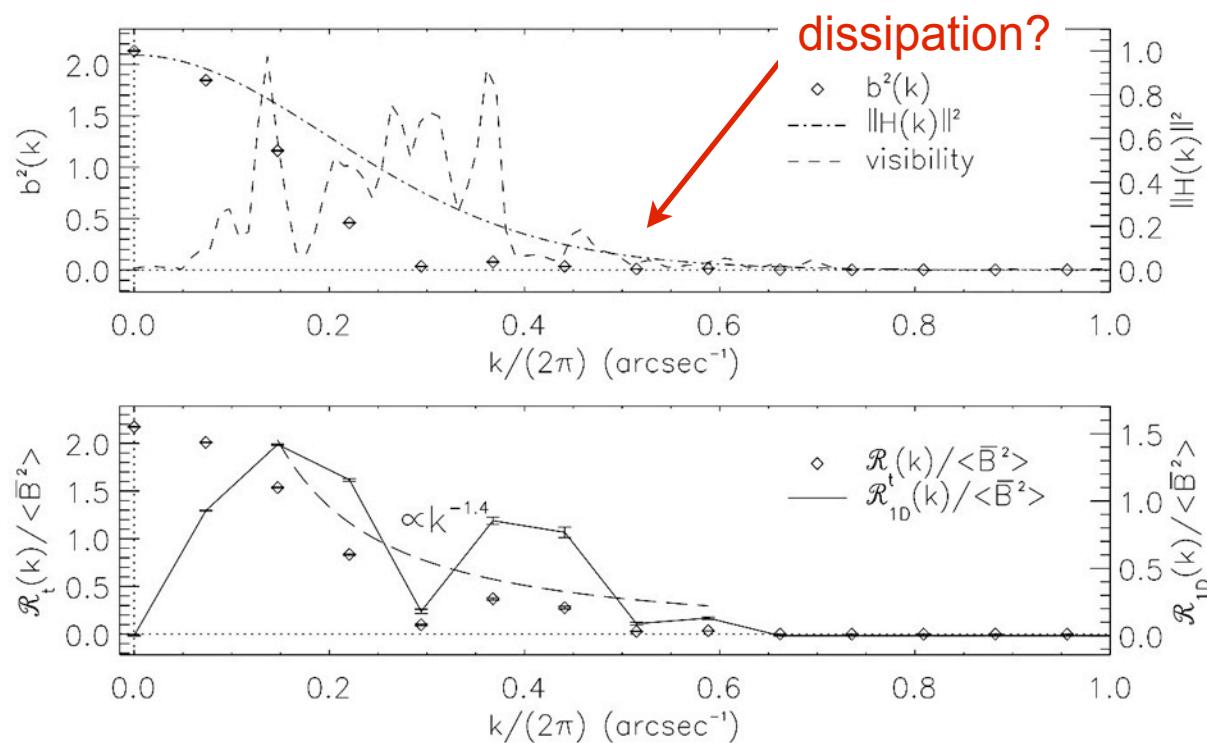
850 μm dust emission (SMA)



B-vectors

beam: 1.6" x 1.0"

sampling: 0.2"



Houde et al. 2011

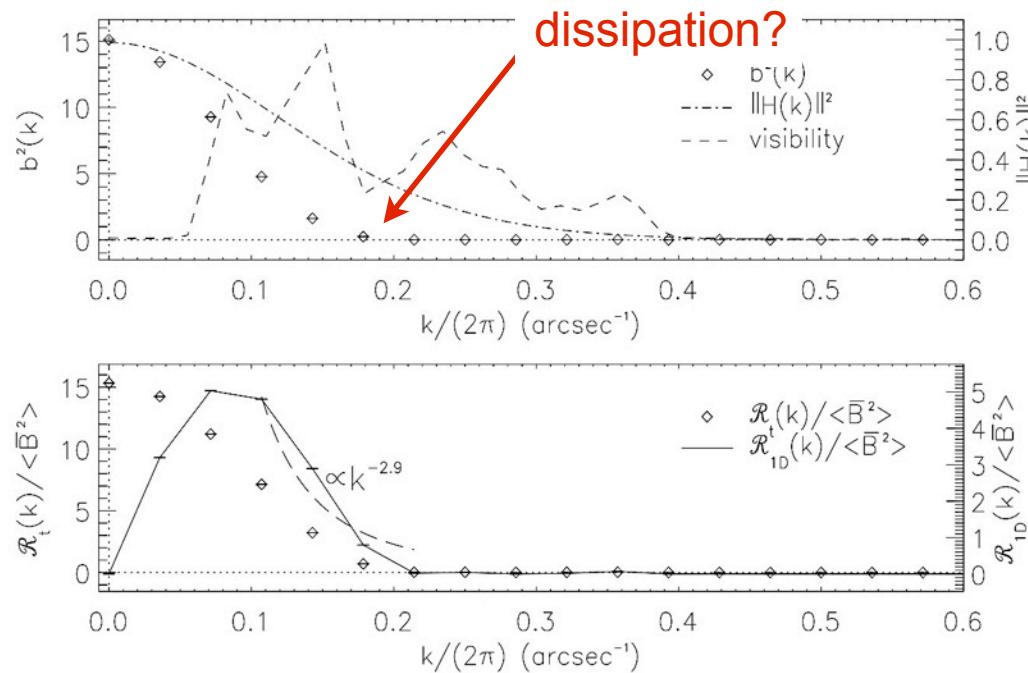
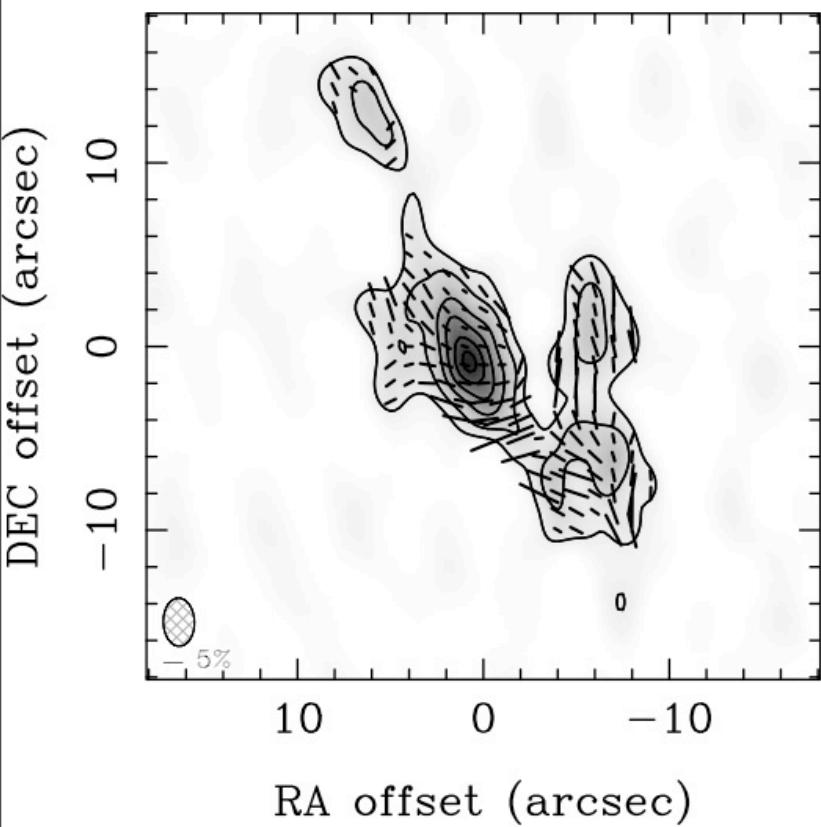
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Turbulent Power Spectrum - Orion KL/SMA

B-vectors

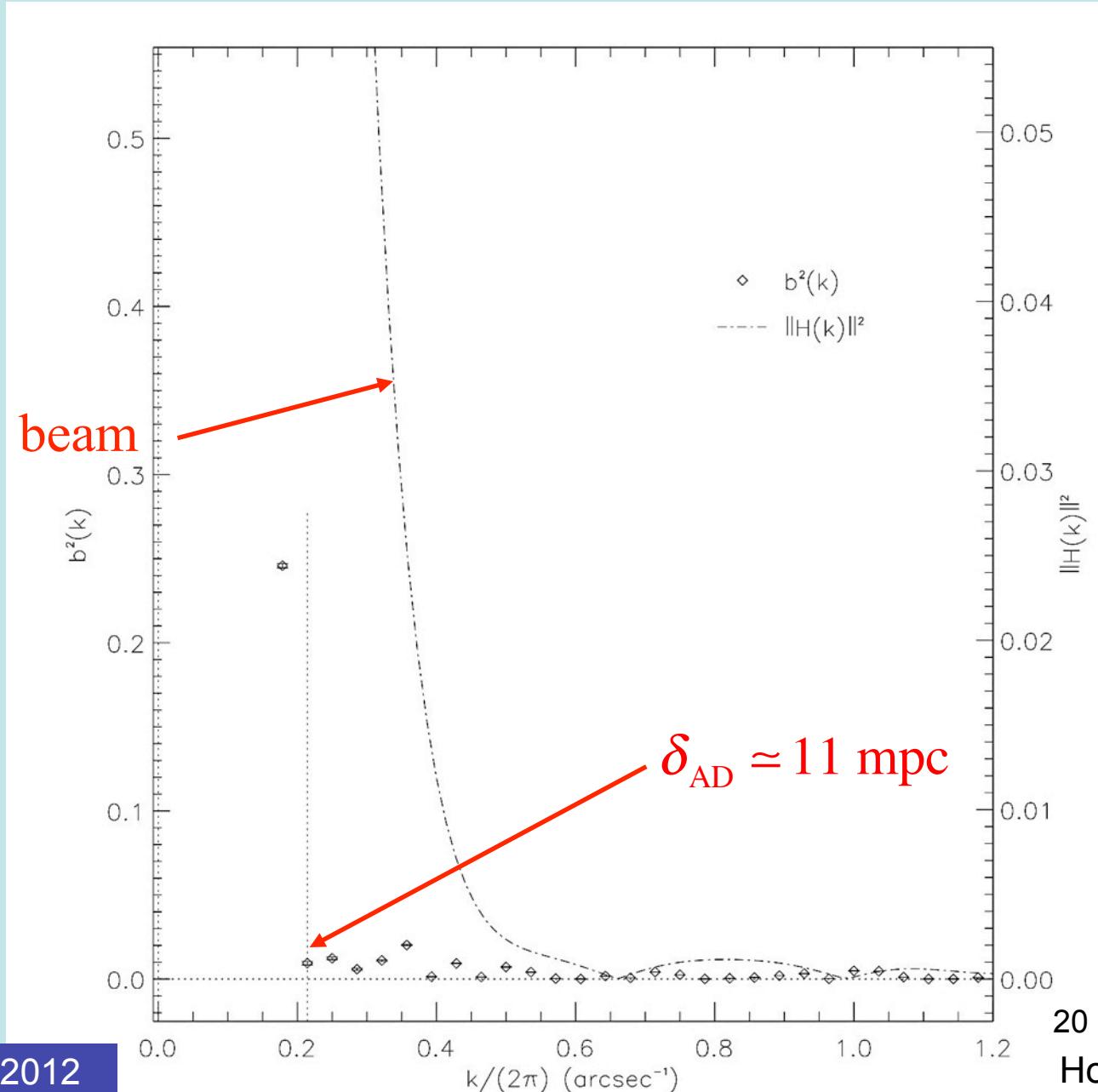
beam: $2.6'' \times 1.7''$

sampling: $0.25''$

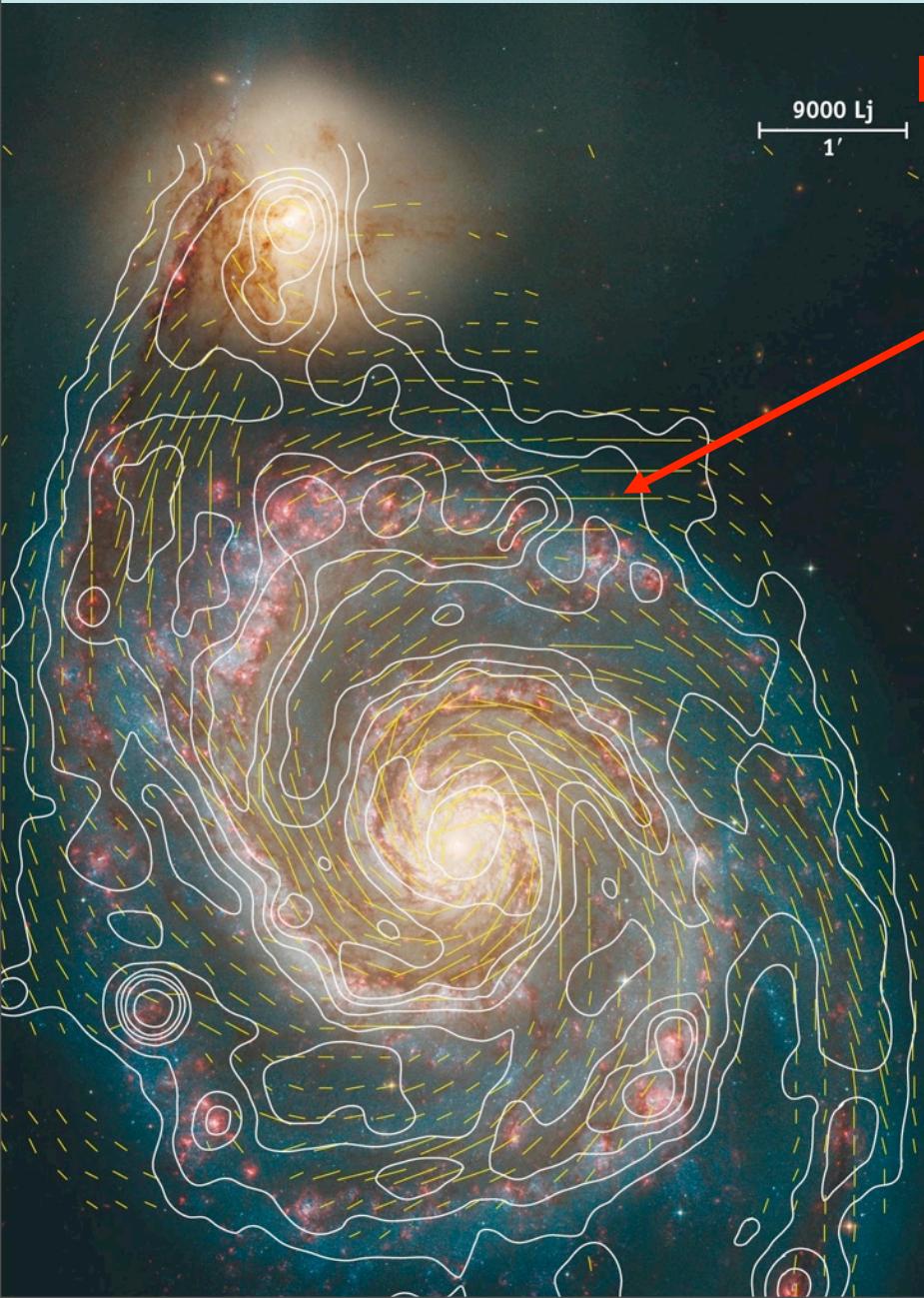


Houde et al. 2011

Ambipolar Diffusion - Orion KL/SMA



Magnetized Turbulence in Galaxies



M51 with Effelsberg (100m) + VLA

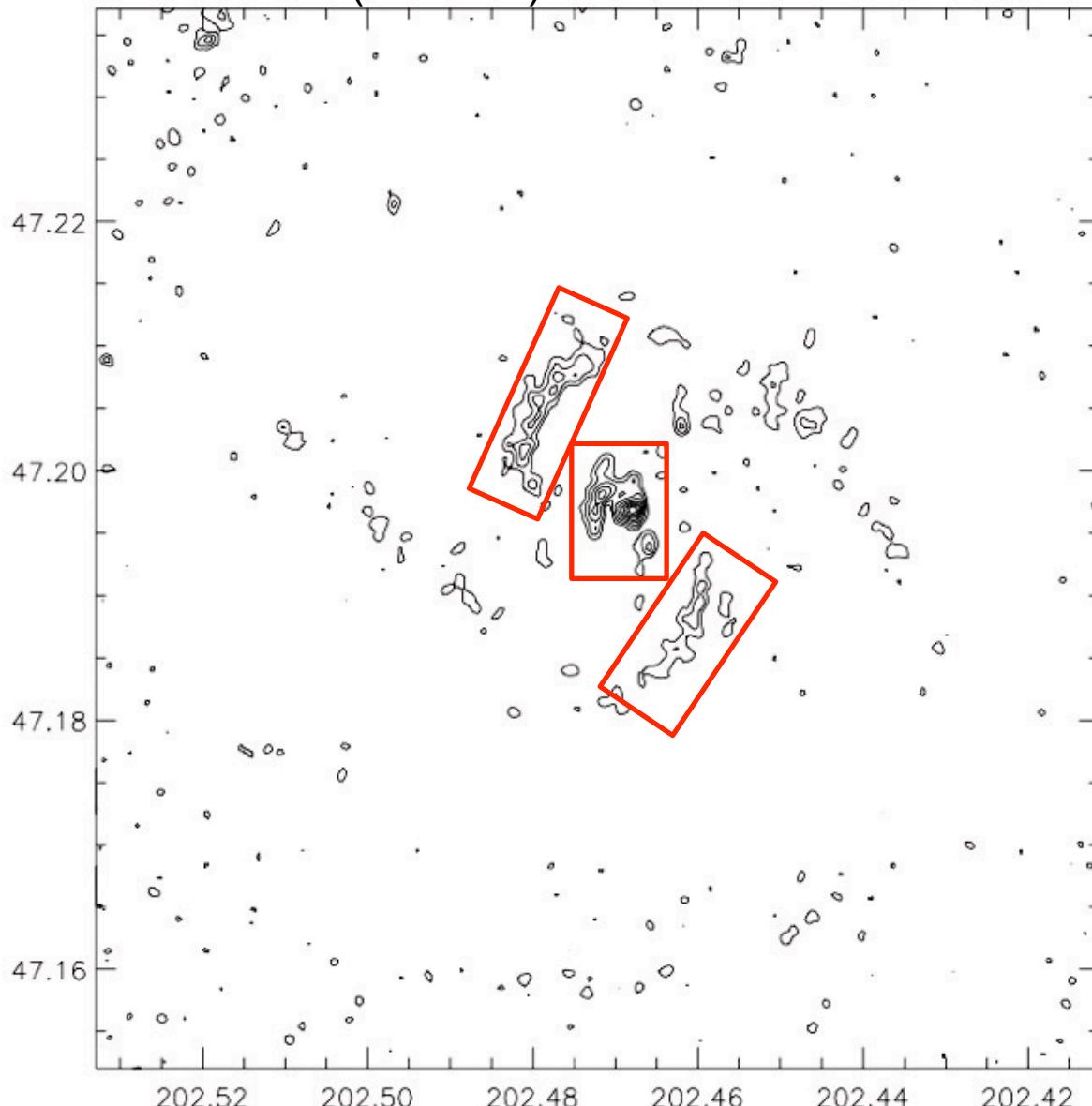
ordered + turbulent fields

$$\mathbf{B} = \mathbf{B}_0 + \mathbf{B}_t$$

Fletcher et al. 2011 (MNRAS)

M51 - Polarized Flux

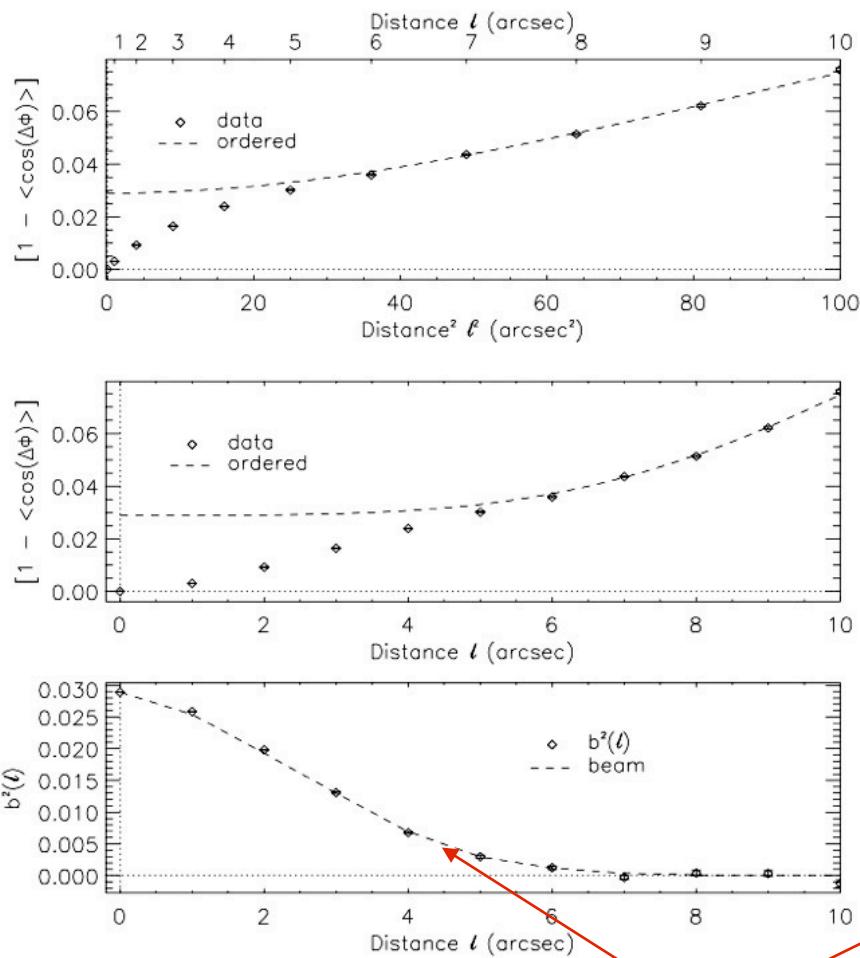
Fletcher et al. 2011 (MNRAS)



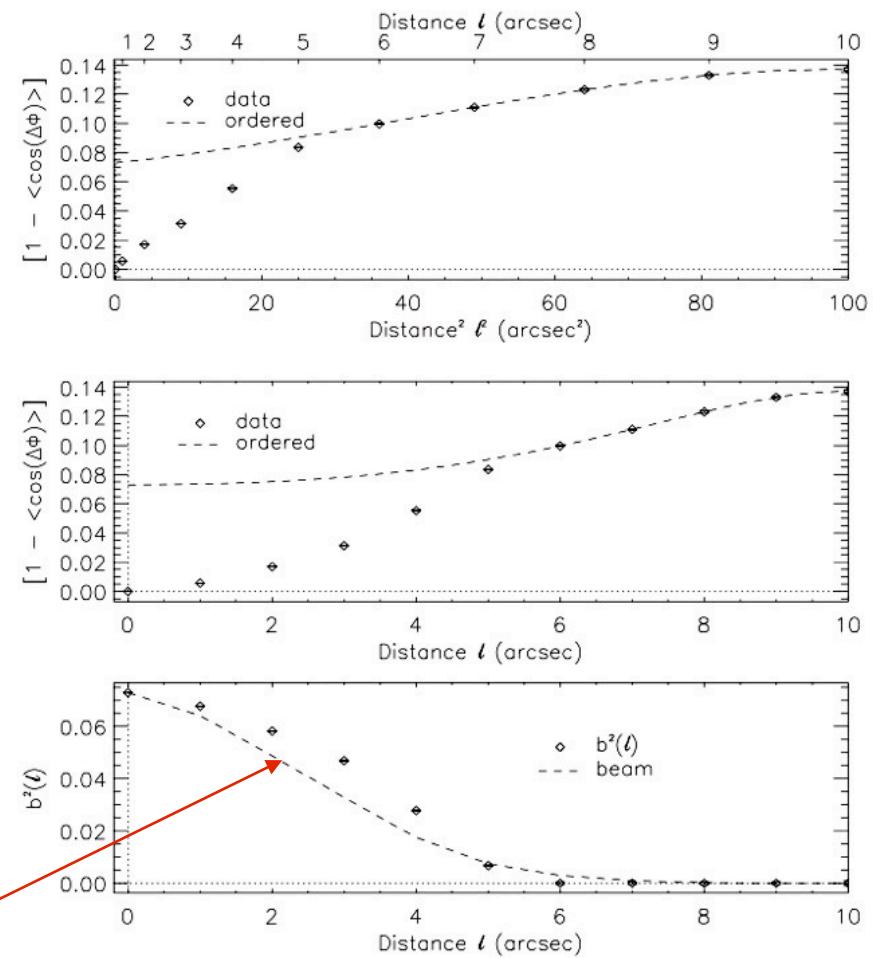
$d = 7.6 \text{ Mpc}$
 $1'' = 37 \text{ pc}$
 $\lambda = 6.2 \text{ cm}$
4" beam
1" sampling

M51 - Structure Functions

Northeast



Southwest



Houde et al. 2012
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M51- Isotropic Turbulence

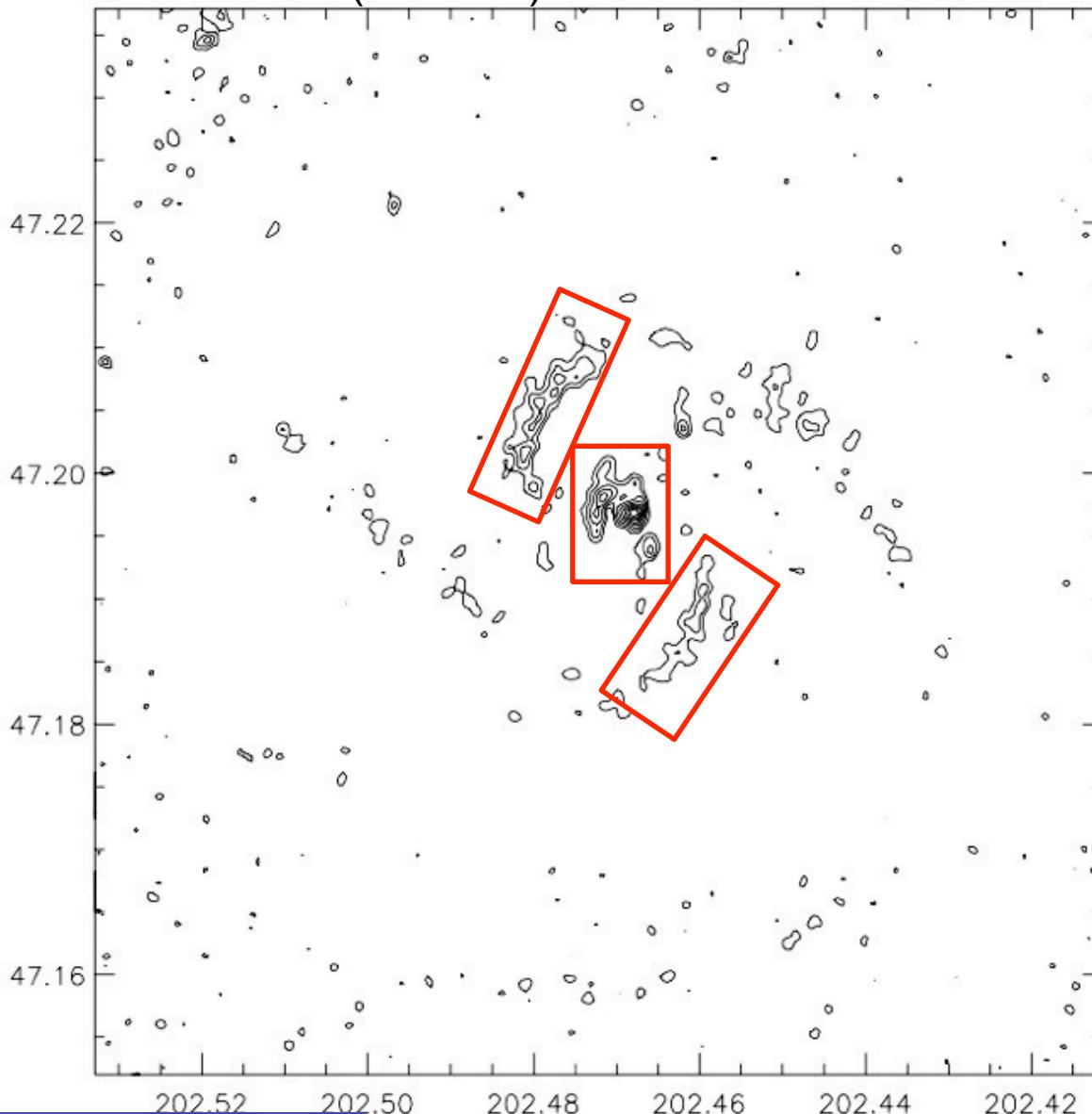
	Northeast	Centre	Southwest
δ (pc)	...	67 ± 7	66 ± 8
N	...	13 ± 3	14 ± 4
\bar{B}_t^2 / \bar{B}^2	0.028 ± 0.002	0.088 ± 0.026	0.072 ± 0.025
B_t^2 / B_0^2	...	1.28 ± 0.29	1.08 ± 0.29
B_t / B_0	...	1.13 ± 0.13	1.04 ± 0.14

From σ_{RM} analysis, Fletcher et al. get

$$\delta \simeq 50 \text{ pc} \quad \text{and} \quad \frac{B_t}{B_0} \simeq 1.2 - 1.5$$

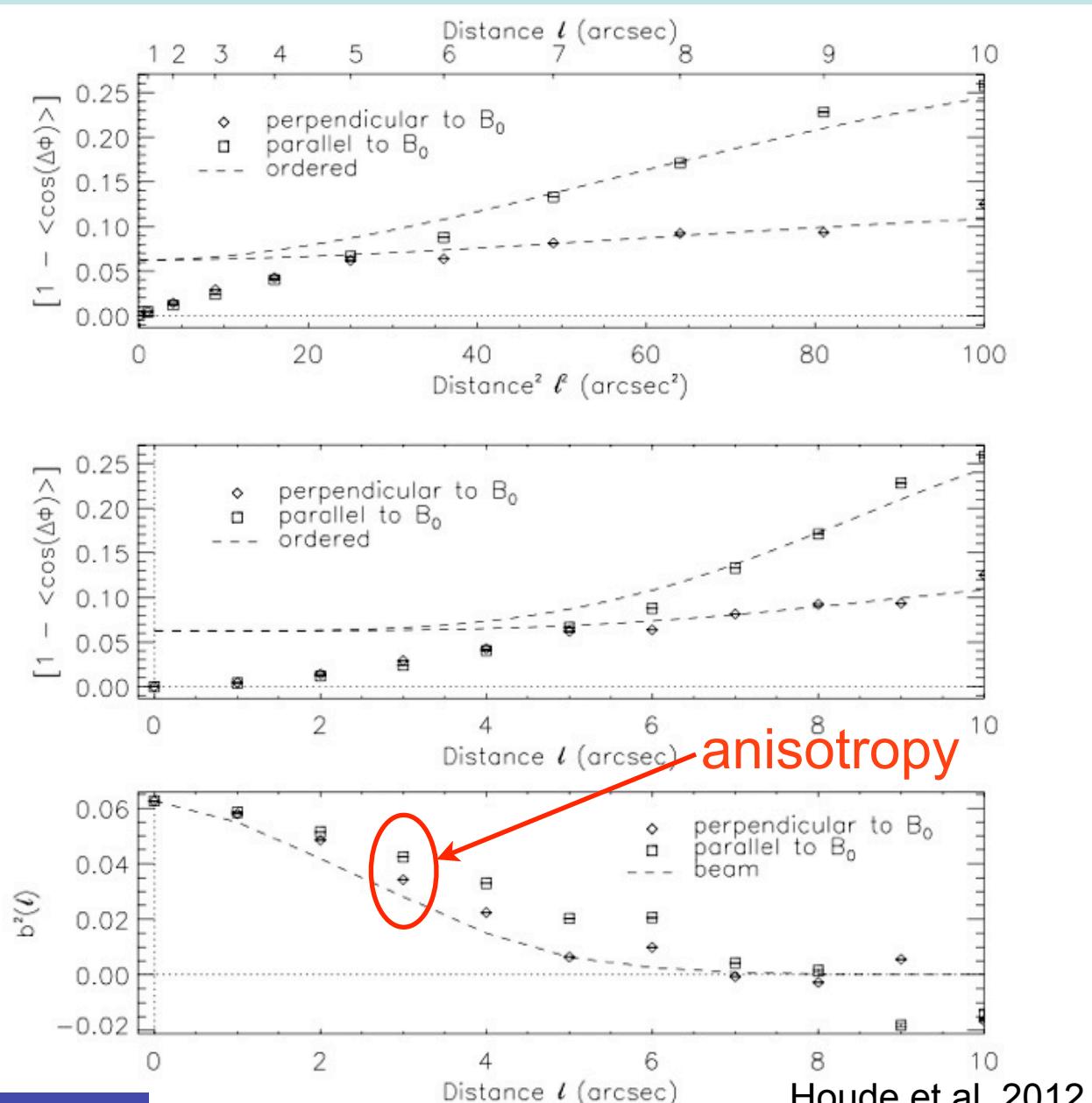
M51- Anisotropic Turbulence

Fletcher et al. 2011 (MNRAS)

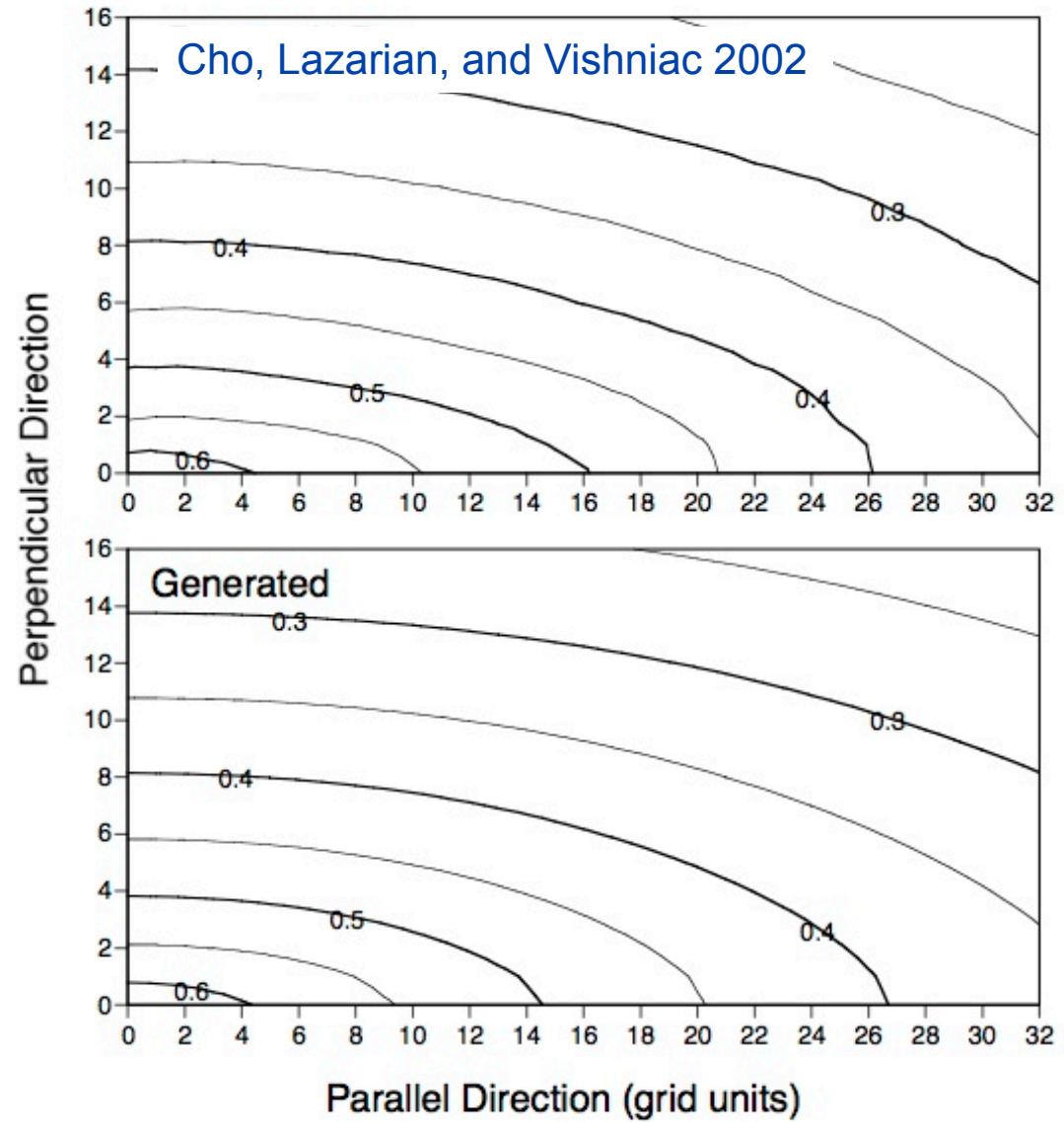
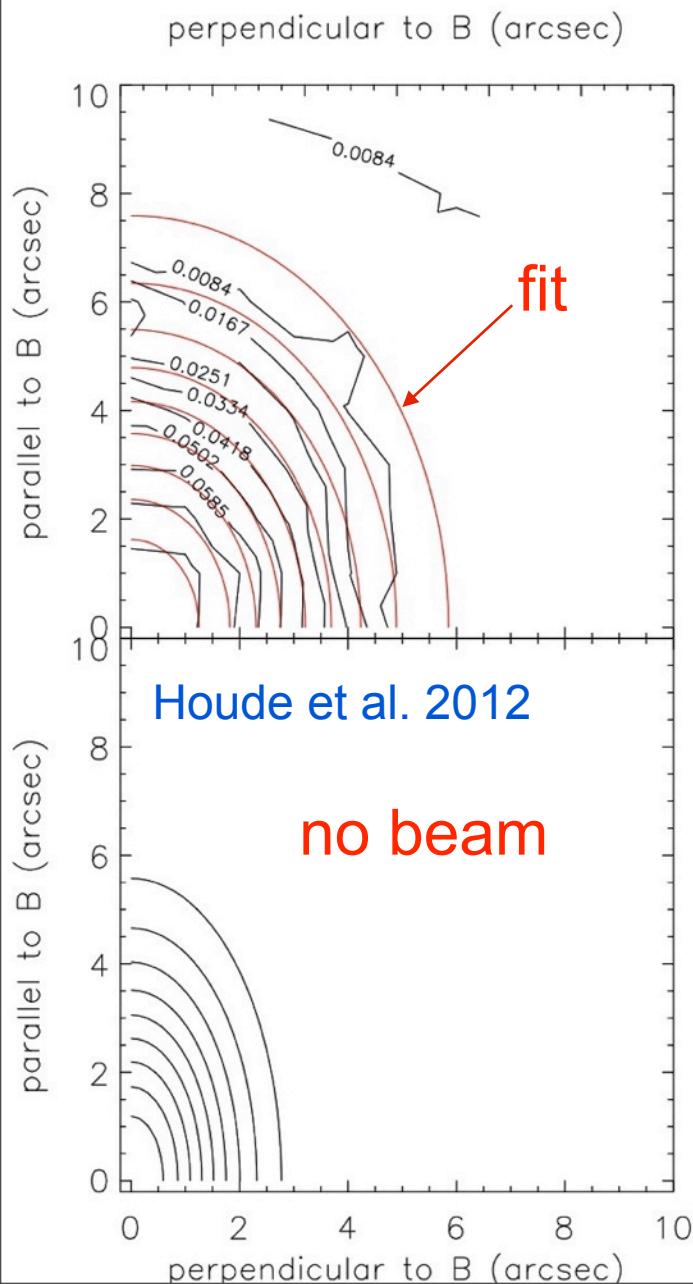


Consider all three regions at once
→ more vectors

M51- Anisotropic Turbulence



M51- Anisotropic Turbulence



M51- Anisotropic Turbulence

$$\delta_{\parallel} \simeq 98 \pm 5 \text{ pc}$$

$$\delta_{\perp} \simeq 54 \pm 3 \text{ pc}$$

$$\delta_{\parallel}/\delta_{\perp} \simeq 1.87 \pm 0.14$$

$$N \simeq 15 \pm 2$$

$$\bar{B}_{\text{t}}^2 / \bar{B}_0^2 \simeq 0.06 \pm 0.01$$

$$B_{\text{t}}^2 / B_0^2 \simeq 1.02 \pm 0.08$$

$$B_{\text{t}} / B_0 \simeq 1.01 \pm 0.04$$

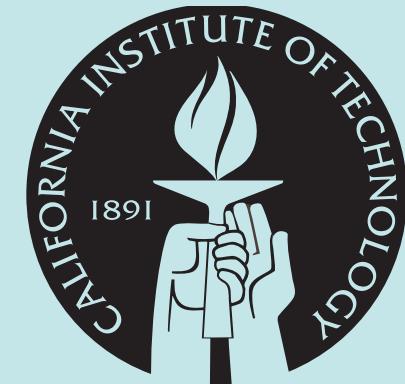
Summary

- Angular dispersion function allows the separation of the turbulent and ordered components of the magnetic field without assuming any model for the latter.
- We can also account for the signal integration process along the line of sight and across the telescope beam.
- With high-enough resolution data → determination of the magnetized turbulent power spectrum (e.g., correlation length, inertial range index, dissipation scale).
- But we need even higher resolution (ALMA) and “larger” single-dish observatories, as well as an increase in the number of “vectors” (SOFIA and CCAT) for anisotropy measurements.

Merci!



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