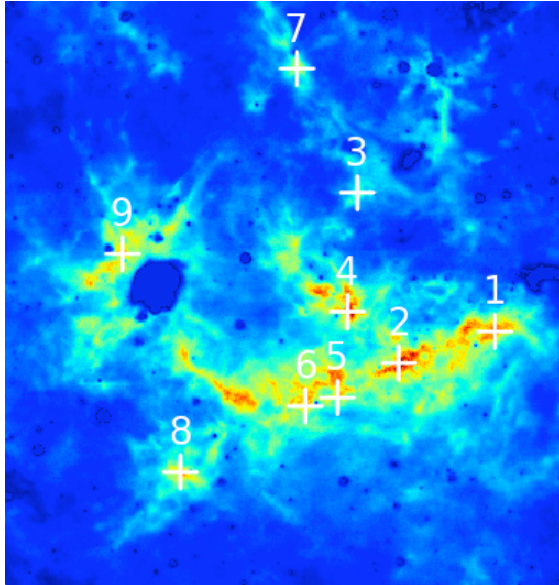
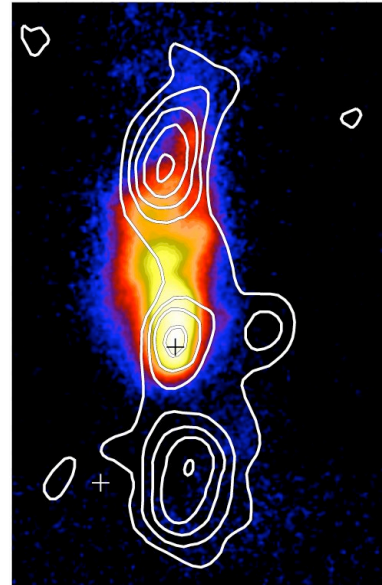


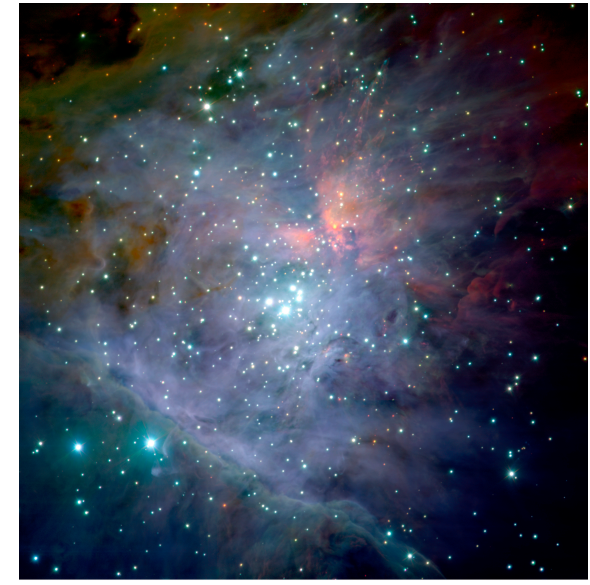
# Massive Star Formation: From Testing Basic Formation Scenarios to Quantitative Constraints on Core Accretion Theory



IRDC G28.4 (Cloud C) (Butler & Tan 09)



MYSO G35.2N (De Buizer 06)



ONC (McCauchrean 01)

Michael Butler  
Audra Hernandez (Wisconsin)  
Yichen Zhang  
Sourav Chatterjee  
Peter Barnes

Paola Caselli (Leeds),  
Francesco Fontani (Arcetri)  
Izaskun Jimenez-Serra (CfA)  
Christopher McKee (UCB)  
Francesco Palla (Arcetri)  
Leonardo Testi (ESO)

Jonathan Tan  
(University of Florida)

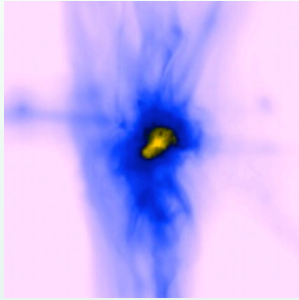
SOFIA Massive Star Formation Team:  
Maite Beltran (Arcetri)  
James De Buizer (NASA/SOFIA)  
Ed Churchwell (Wisconsin)  
Göran Sandell (NASA/SOFIA)  
Ralph Shuping (NASA/SOFIA)  
Jan Staff (LSU)  
Charles Telesco (UF)  
Barbara Whitney (SSI)  
Yichen Zhang (UF)

# Outline

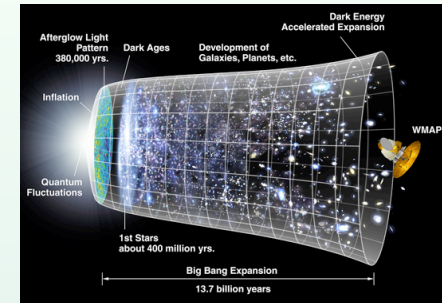
- **Introduction**
  - why study massive star formation?
  - why is it difficult?
  - open questions: formation theories
- **Initial Conditions**
  - Infra-Red Dark Clouds (IRDCs)
  - Massive Starless Cores?
- **Massive Protostars**
  - Overview of expected properties
  - Radiative transfer modeling
  - SOFIA FORCAST observations of G35.2N



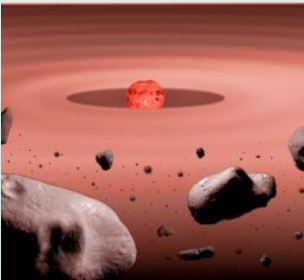
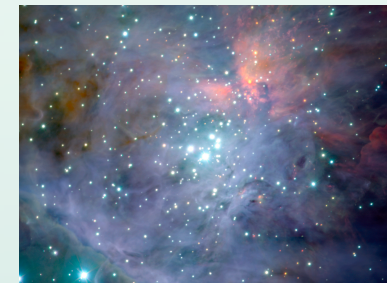
# Why study massive star formation?



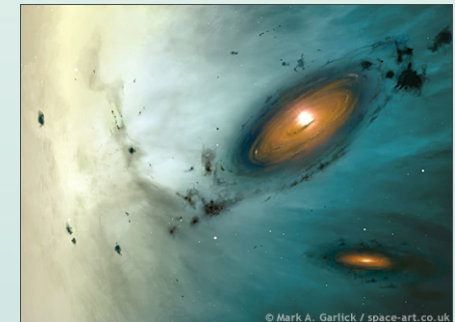
The **First (Pop III) Stars** were likely massive, some potentially supermassive stars, reionizing the universe and producing the first metals.



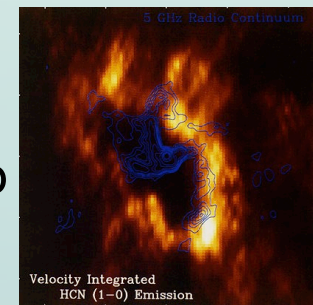
**Galaxies form and evolve** by forming **star clusters**, where the influence of massive stars is paramount. Massive stars are what tend to be seen in distant galaxies.



**Planets form** from the crumbs left over from star formation. Planet & star formation in star clusters can be influenced by massive star feedback.



**Supermassive black hole formation** may be via massive star clusters or Pop III stars. **Supermassive black hole accretion** is likely to be regulated by star formation.



# Why not to work on massive star formation...

A complicated, nonlinear process

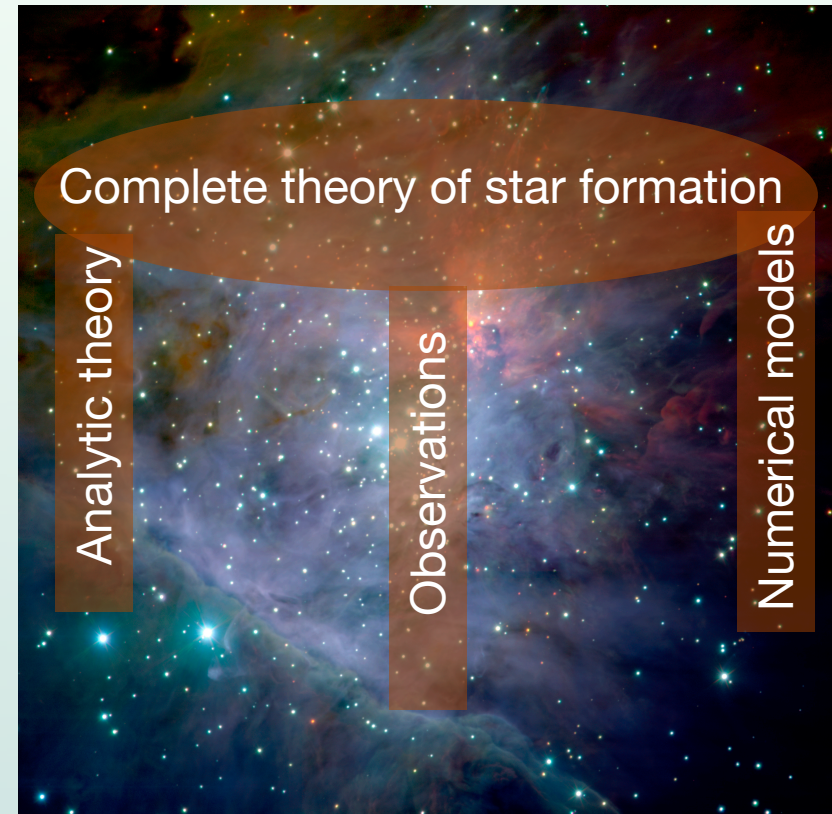
Physics:

Gravity vs pressure (thermal, magnetic, turbulence, radiation, cosmic rays) and shear.

Heating and cooling, generation and decay of turbulence, generation (dynamo) and diffusion of B-fields, etc.  
Chemical evolution of dust and gas.

Wide range of scales (~12 dex in space, time) and multidimensional.

Uncertain/unconstrained initial conditions/boundary conditions.



Some notation:

Core -> star or close binary

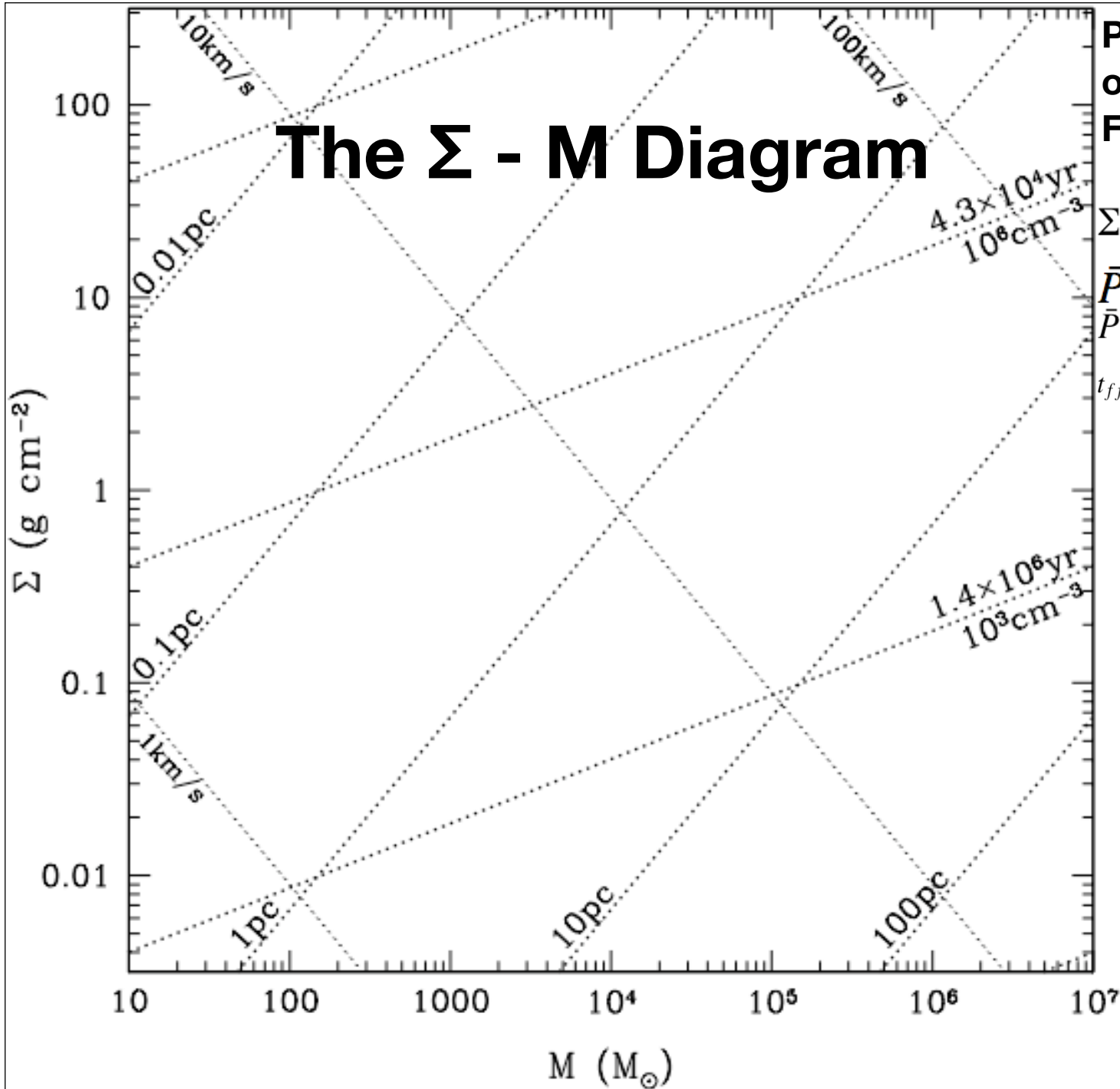
Clump -> star cluster

# Massive Star Formation: Open Questions

- **Causation:** external triggering or spontaneous gravitational instability?
- **Initial conditions:** how close to equilibrium?
- **Accretion mechanism:** [turbulent/magnetic/thermal-heat]-regulated fragmentation vs competitive accretion
- **Timescale:** fast or slow (# of dynamical times)?
- **End result**
  - Initial mass function (IMF)
  - Binary fraction and properties
  - Initial cluster mass function (ICMF)
  - Efficiency and Rate (& relation to galaxy-scale)

How do these properties vary with environment?

# The $\Sigma$ - M Diagram



## Physical Properties of Massive Star-Forming Regions

$$\Sigma \equiv \frac{M}{\pi R^2}$$

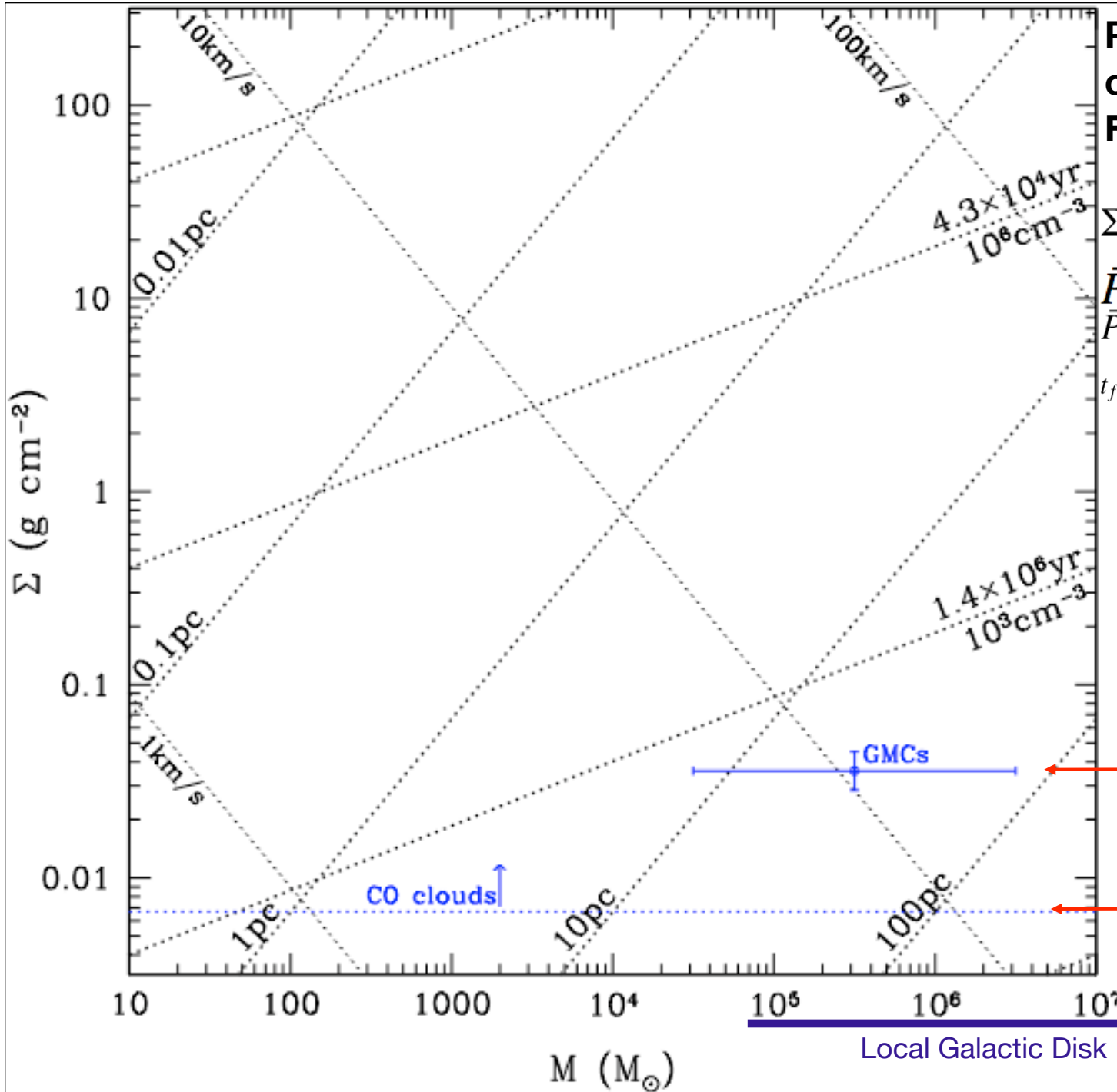
$$\bar{P} \simeq G \Sigma^2$$

$$\bar{P}/k = 4.3 \times 10^8 \Sigma^2 \text{ K cm}^{-3}$$

$$t_{ff} = \left( \frac{3\pi}{32G\rho} \right)^{1/2}$$



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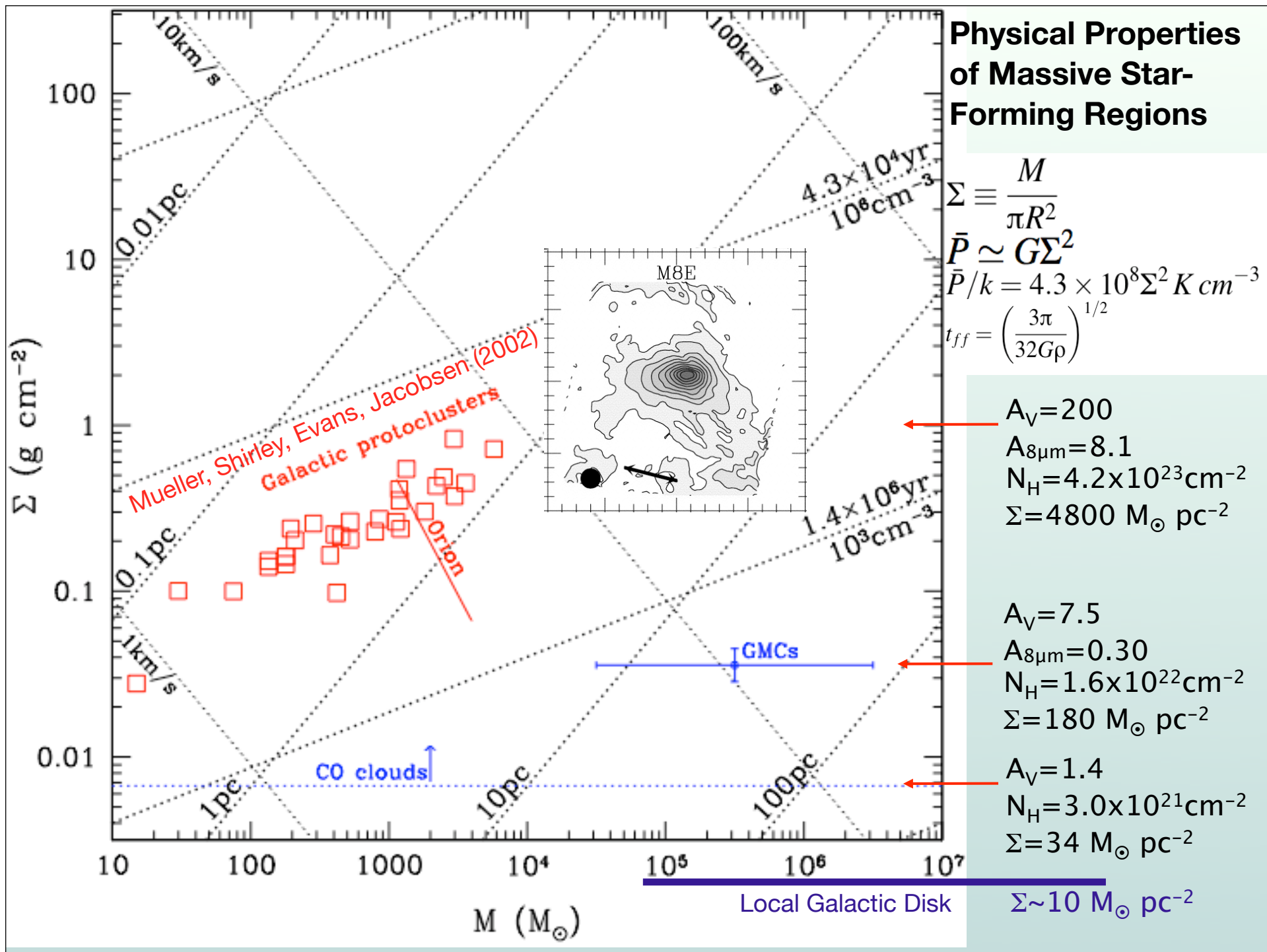
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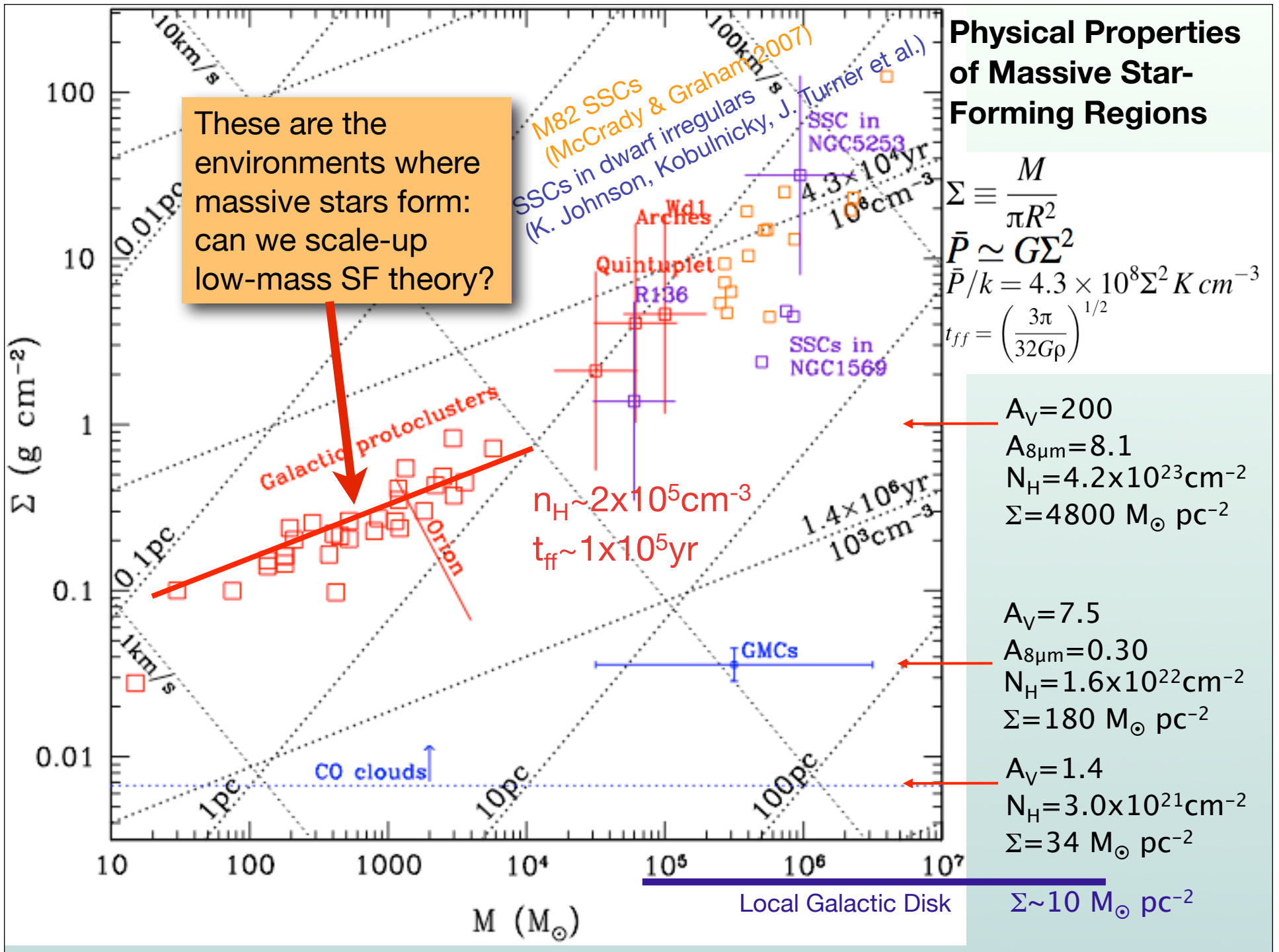
$A_V = 7.5$   
 $A_{8\mu\text{m}} = 0.30$   
 $N_H = 1.6 \times 10^{22} \text{ cm}^{-2}$   
 $\Sigma = 180 M_\odot \text{ pc}^{-2}$

$A_V = 1.4$   
 $N_H = 3.0 \times 10^{21} \text{ cm}^{-2}$   
 $\Sigma = 34 M_\odot \text{ pc}^{-2}$

Local Galactic Disk

$\Sigma \sim 10 M_\odot \text{ pc}^{-2}$





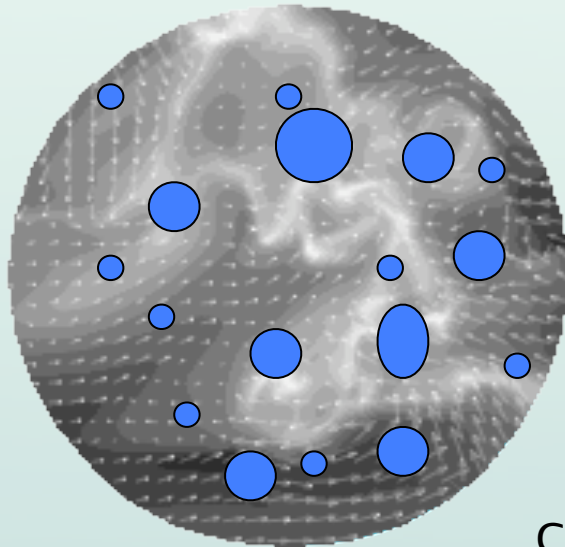
# Massive Star Formation Theories

Stahler, Palla, Ho 2000; Beuther, Churchwell, McKee, Tan 2007; Zinnecker & Yorke 2007

## Core Accretion

Myers & Fuller 1992; Caselli & Myers 1995;  
McLaughlin & Pudritz 1997;  
Yorke & Sonnhalter 2002;  
McKee & Tan 2002, 2003;

Stars form from “cores”,  $M_{\text{core}} \sim 2m_*$ ,  
which fragment from the clump



$$\bar{P} = \phi_P G \Sigma^2$$

If in equilibrium,  
then self-gravity  
is balanced by  
internal pressure:  
B-field, turbulence,  
radiation pressure  
(thermal P is small)

Cores form from this  
turbulent medium: at any given time there  
is a small mass fraction in unstable cores.  
These cores collapse quickly to form  
individual stars or binaries.

## Equilibrium Star Cluster Formation

Tan, Krumholz, McKee (2006)

Small SFE per free-fall time  $\sim 0.02$

Formation time  $\geq$  several to many  $t_{\text{ff}}$

-> Age spreads  $\sim$  Myr in clusters

Turbulence maintained by protostellar winds  
(see also Nakamura & Li 2007)

## Core continues to accrete from clump

$$\dot{m}_{\text{acc}} = 2.50 \times 10^{-4} \left( \frac{A \phi_{\rho, \text{core}} f_g^2 \alpha_{\text{vir}} \phi_{\text{grav}}^2}{k_P^2 \epsilon_{\text{core}}^2 \phi_{\bar{P}}} \right)^{1/2} \quad (54)$$

$$\times \left( \frac{m_{*f}}{30 M_{\odot}} \right) \left( \frac{M_{\text{cl}}}{4000 M_{\odot}} \right)^{-1/4} \Sigma_{\text{cl}}^{3/4} M_{\odot} \text{ yr}^{-1}$$

$$\rightarrow 7.9 \times 10^{-4} \left( \frac{\phi_{\text{grav}}}{1.6} \right) \left( \frac{m_{*f}}{30 M_{\odot}} \right) \left( \frac{M_{\text{cl}}}{4000 M_{\odot}} \right)^{-1/4} \\ \times \Sigma_{\text{cl}}^{3/4} M_{\odot} \text{ yr}^{-1} . \quad (55)$$

Thus, we see that a massive core will tend to interact with  
clump material at a rate that is comparable to the rate at  
which it is collapsing or being eroded by star formation,  
 $\dot{m}_* / \epsilon_{\text{core}}$ . This rate is a factor of several smaller in SSCs with  
 $M_{\text{cl}} \sim 10^6 M_{\odot}$ .



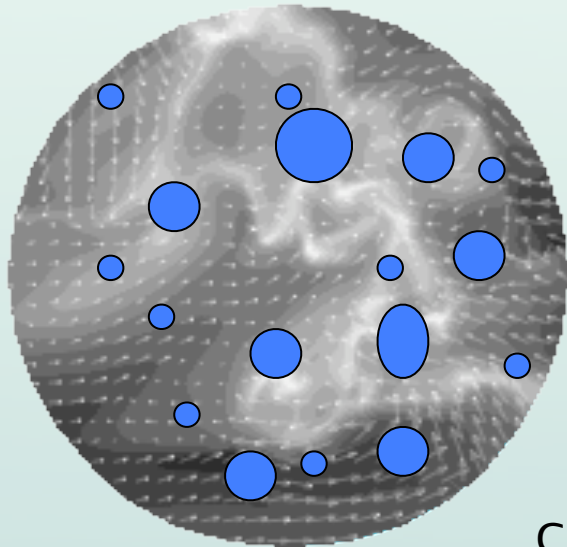
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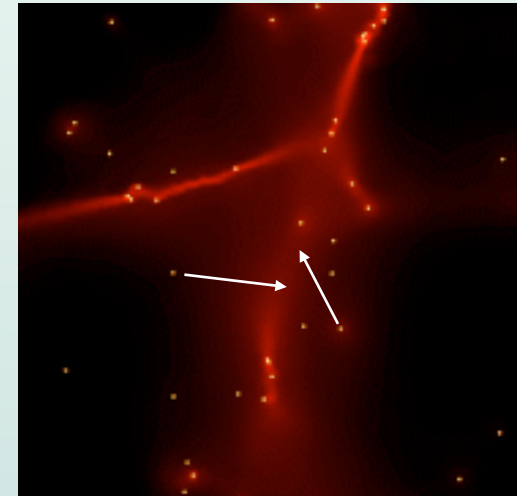
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then self-gravity  
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radiation pressure  
(thermal P is small)

Cores form from this  
turbulent medium: at any given time there  
is a small mass fraction in unstable cores.  
These cores collapse quickly to form  
individual stars or binaries.

## Competitive Accretion

Bonnell, Clarke, Bate, Pringle 2001;  
Bonnell, Vine, & Bate 2004;  
Schmeja & Klessen 2004;  
Wang, Li, Abel, Nakamura 2010;

Stars gain most mass by Bondi-  
Hoyle accretion of ambient gas

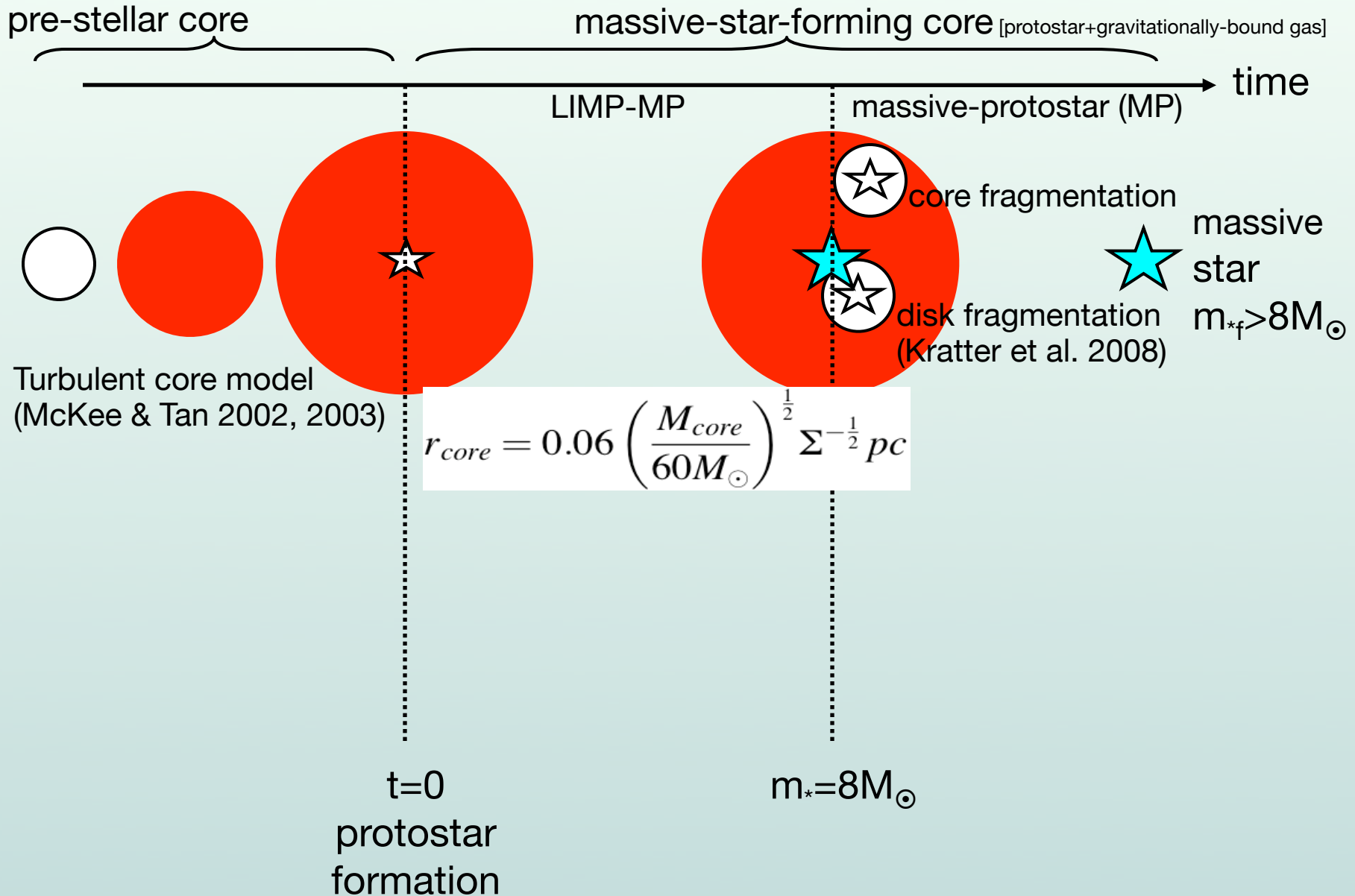


Based on simulations including  
only thermal pressure.

Requires global collapse of clump  
(Krumholz, McKee & Klein 2005)

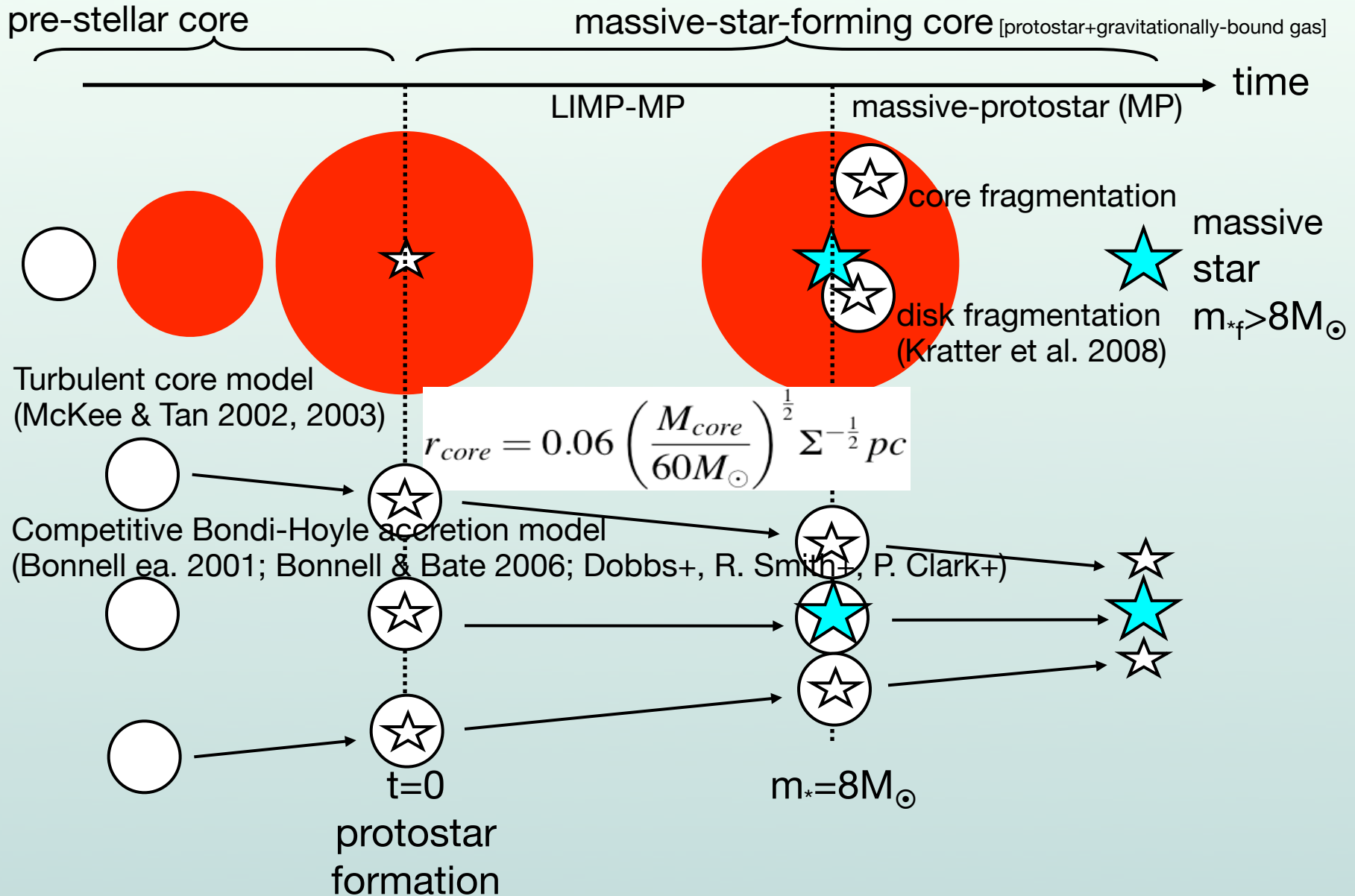
# Schematic Differences Between Massive Star Formation Theories

Beuther, Churchwell, McKee, Tan (2007); Tan (2008)



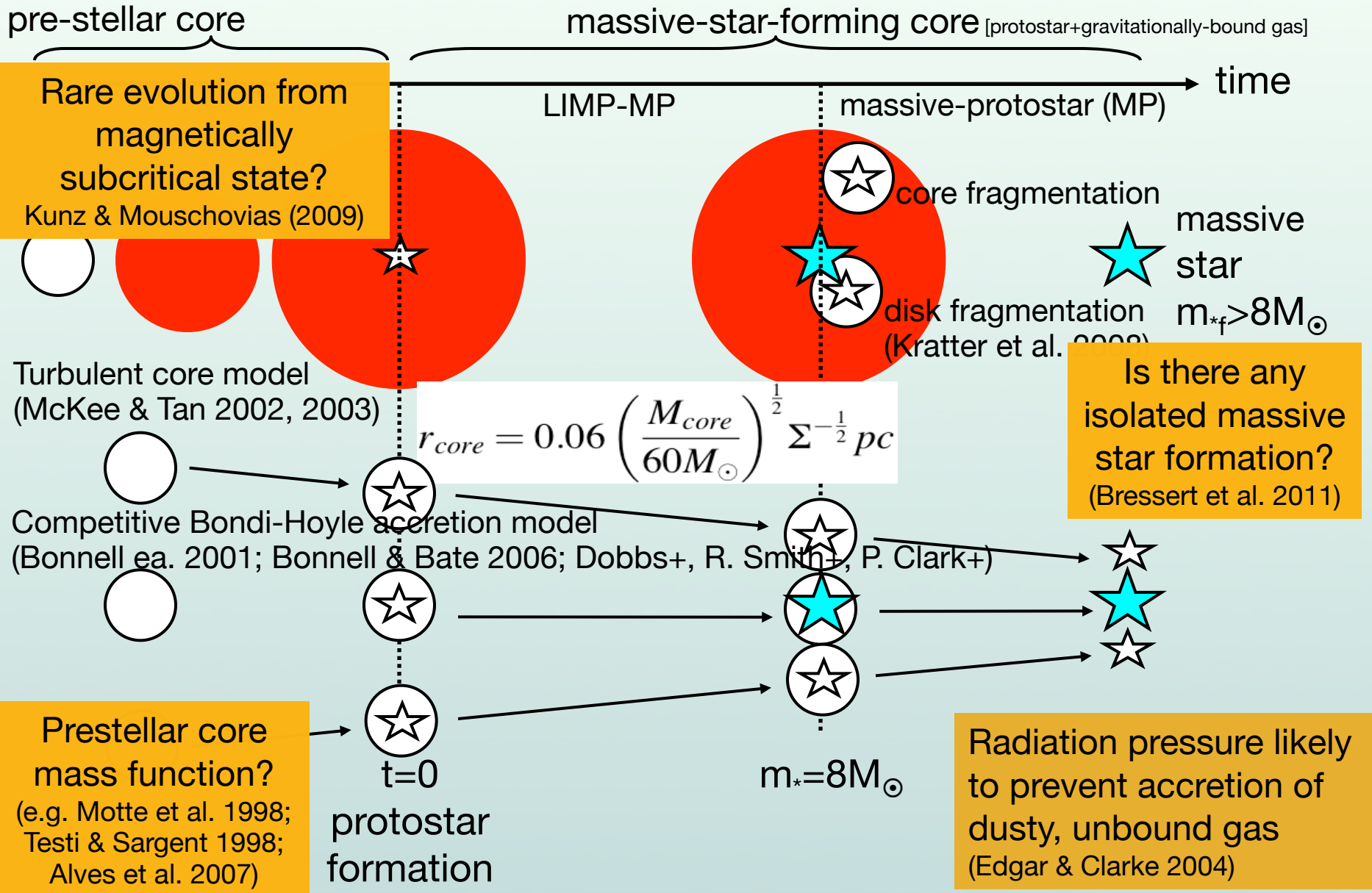
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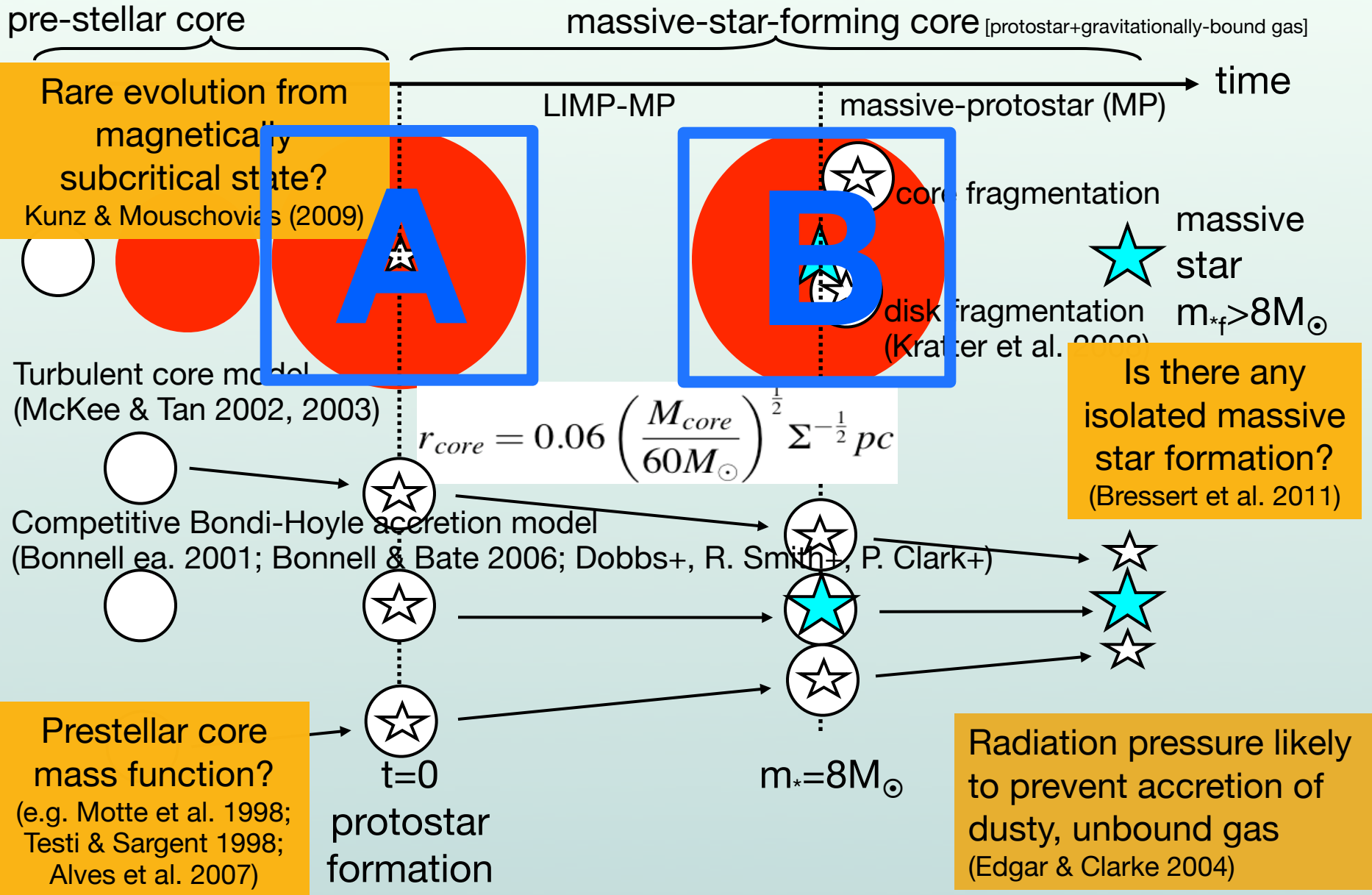
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Tan (2008)

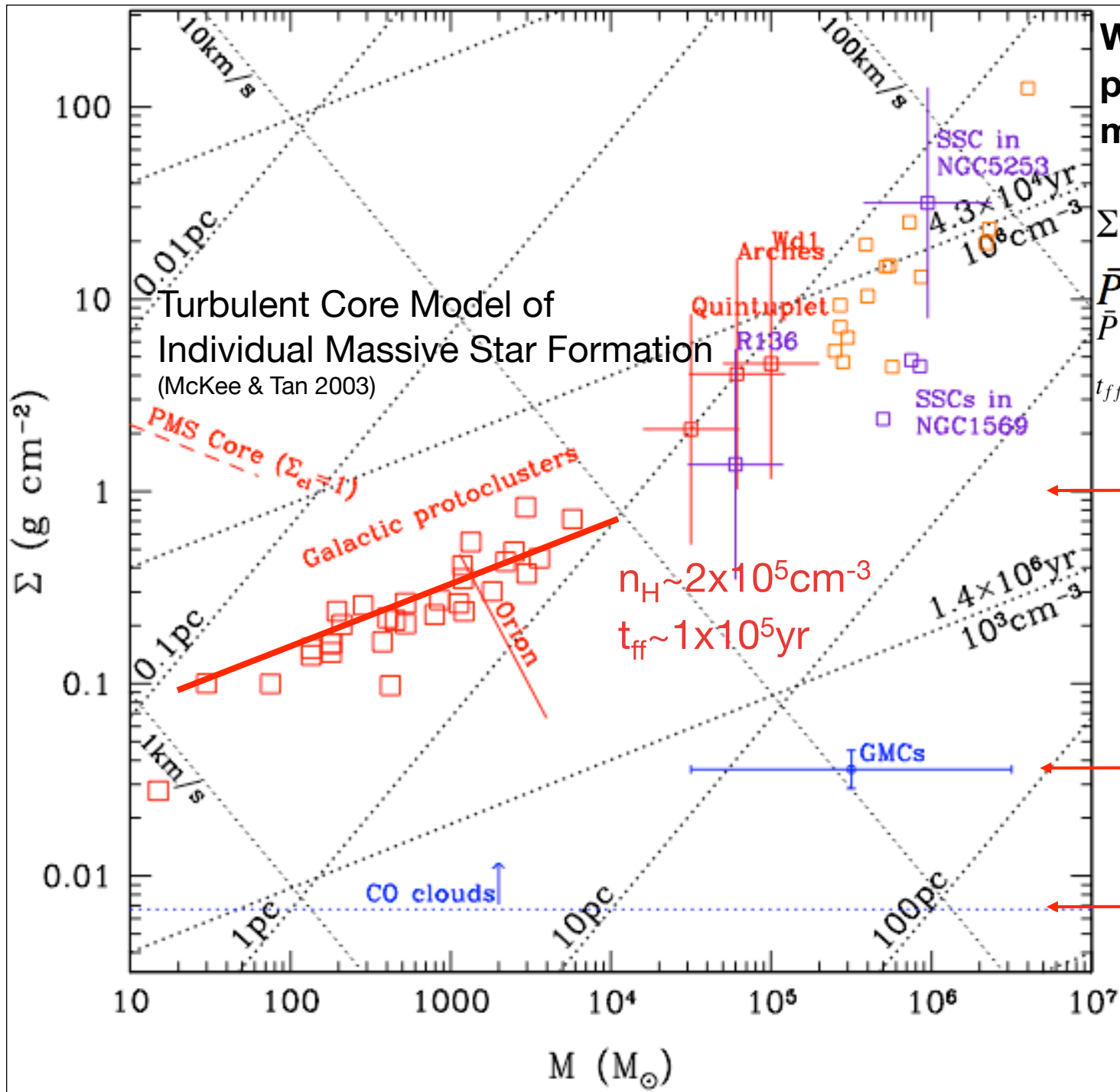




# Schematic Differences Between Massive Star Formation Theories

Beuther, Churchwell, McKee, Tan (2007);  
Tan (2008)





What are the pressures where massive stars form?

$$\Sigma \equiv \frac{M}{\pi R^2}$$

$$\bar{P} \simeq G \Sigma^2$$

$$\bar{P}/k = 4.3 \times 10^8 \Sigma^2 K \text{ cm}^{-3}$$

$$t_{ff} = \left( \frac{3\pi}{32G\rho} \right)^{1/2}$$

$A_V=200$   
 $A_{8\mu\text{m}}=8.1$   
 $N_H=4.2 \times 10^{23} \text{ cm}^{-2}$   
 $\Sigma=4800 M_{\odot} \text{ pc}^{-2}$

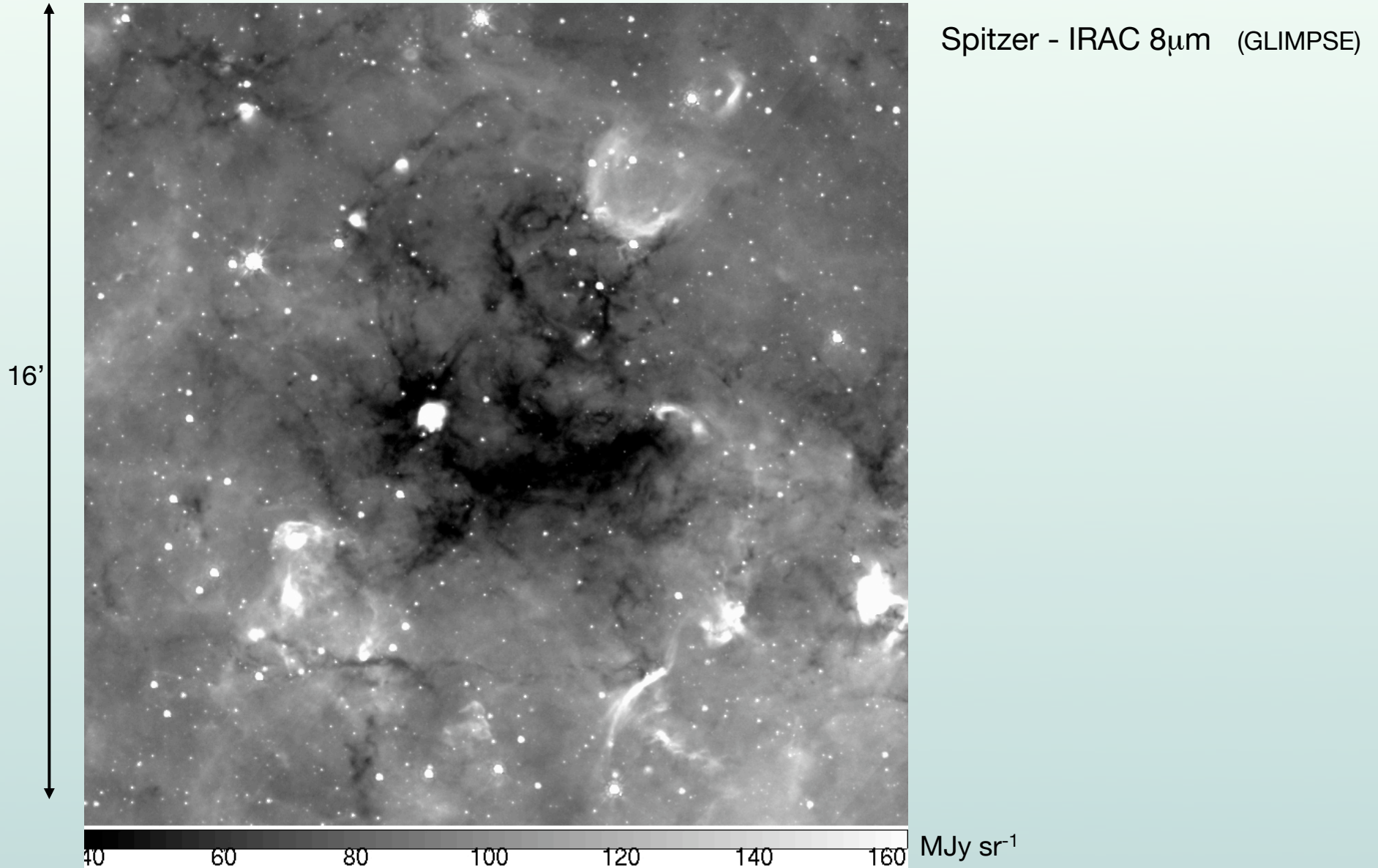
$A_V=7.5$   
 $A_{8\mu\text{m}}=0.30$   
 $N_H=1.6 \times 10^{22} \text{ cm}^{-2}$   
 $\Sigma=180 M_{\odot} \text{ pc}^{-2}$

$A_V=1.4$   
 $N_H=3.0 \times 10^{21} \text{ cm}^{-2}$   
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# Mid-IR Extinction Mapping of Infrared Dark Clouds

(Butler & Tan 2009, 2011; see also Peretto & Fuller 2009; Ragan et al. 2009; Battersby et al. 2010)

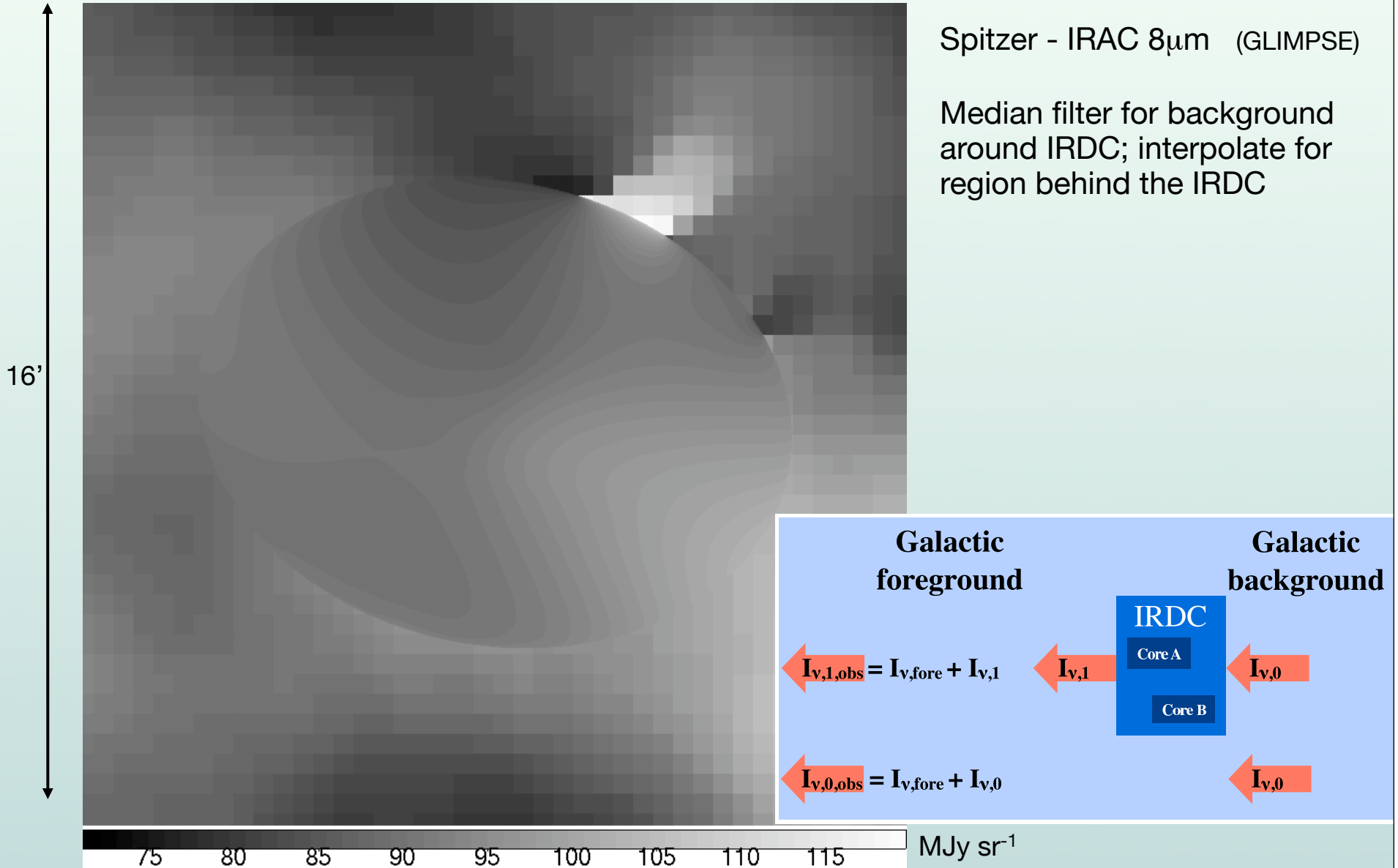
**G28.37+00.07**



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**G28.37+00.07**

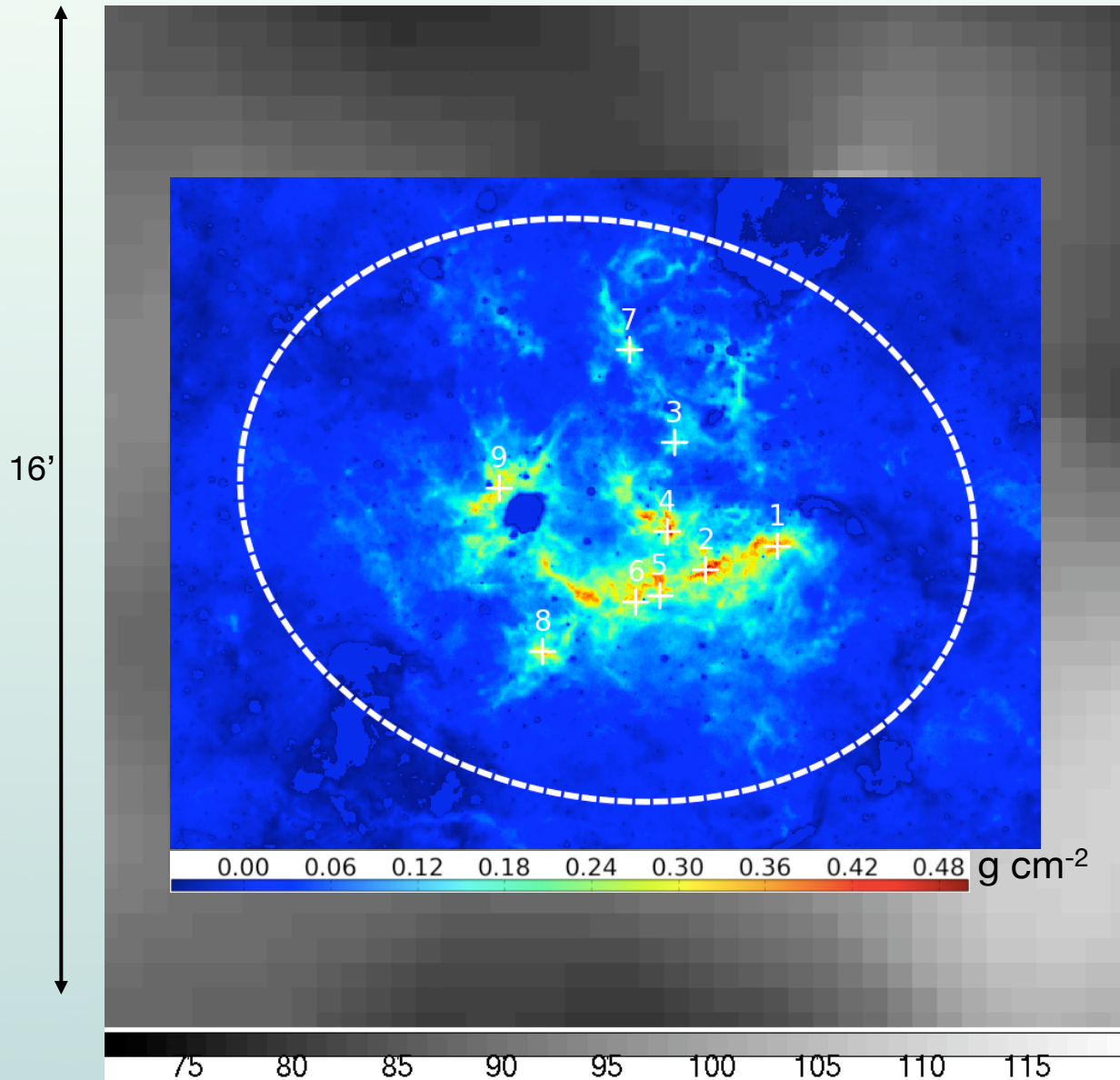




# Mid-IR Extinction Mapping of Infrared Dark Clouds

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**G28.37+00.07**



Spitzer - IRAC  $8\mu\text{m}$  (GLIMPSE)

Median filter for background around IRDC; interpolate for region behind the IRDC

Correct for foreground -> Choose nearby clouds.  
BT09: analytic model  
BT11: observed saturation in independent cores

**~Arcsecond scale maps of regions up to  $\Sigma \sim 0.5 \text{ g cm}^{-2}$ ; independent of dust temp.**

Distance from molecular line velocities ->  $M(\Sigma)$

# Massive starless cores

Cores with power law density structure

$$\rho_c(r) = \rho_{s,c} \left( \frac{r}{R_c} \right)^{-k_{\rho,c}} \quad k_{\rho,c} = 1.8$$

They contain many thermal Jeans masses.

B-fields may be suppressing fragmentation within the core.

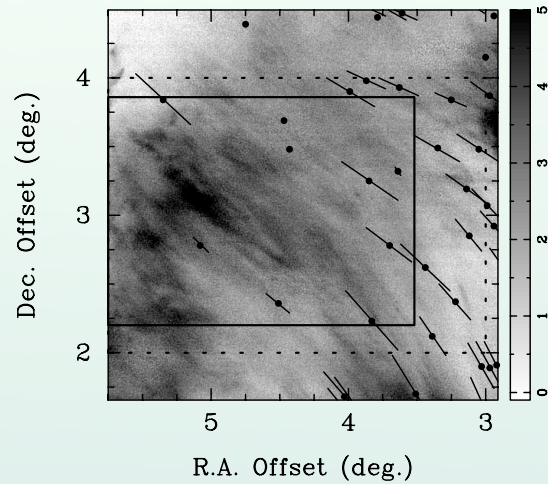
$$M_{BE} = 1.182 \frac{c_{th}^4}{(G^3 P_{s,core})^{1/2}} \rightarrow 0.0504 \left( \frac{T}{20 \text{ K}} \right)^2 \frac{1}{\Sigma_{cl}} M_{\odot}$$

Magnetic Critical Mass (Bertoldi & McKee 1992)

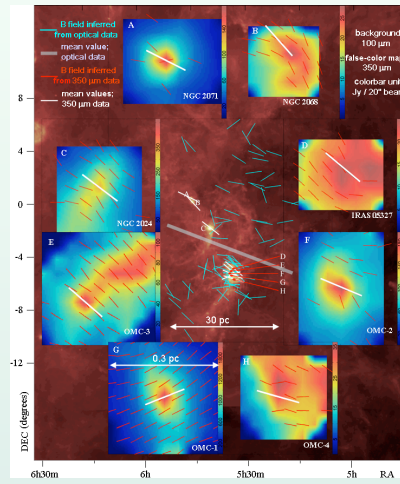
$$M_B = 79 c_{\Phi}^3 \left( \frac{R}{Z} \right)^2 \frac{\bar{v}_A^3}{(G^3 \bar{\rho})^{1/2}} = 1020 \left( \frac{R}{Z} \right)^2 \left( \frac{\bar{B}}{30 \mu\text{G}} \right)^3 \left( \frac{10^3 \text{ cm}^{-3}}{\bar{n}_H} \right)^2 M_{\odot}$$

$$n_H \sim 10^5 \text{ cm}^{-3}, B \sim 200 \mu\text{G} \rightarrow M_B \sim 100 M_{\odot}$$

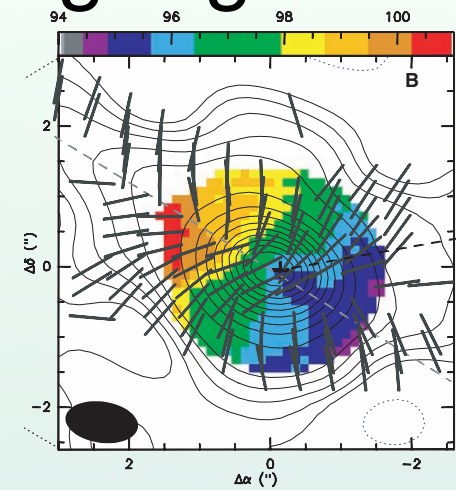
# Strong magnetic fields in star-forming regions



Taurus (Heyer et al. 2008)

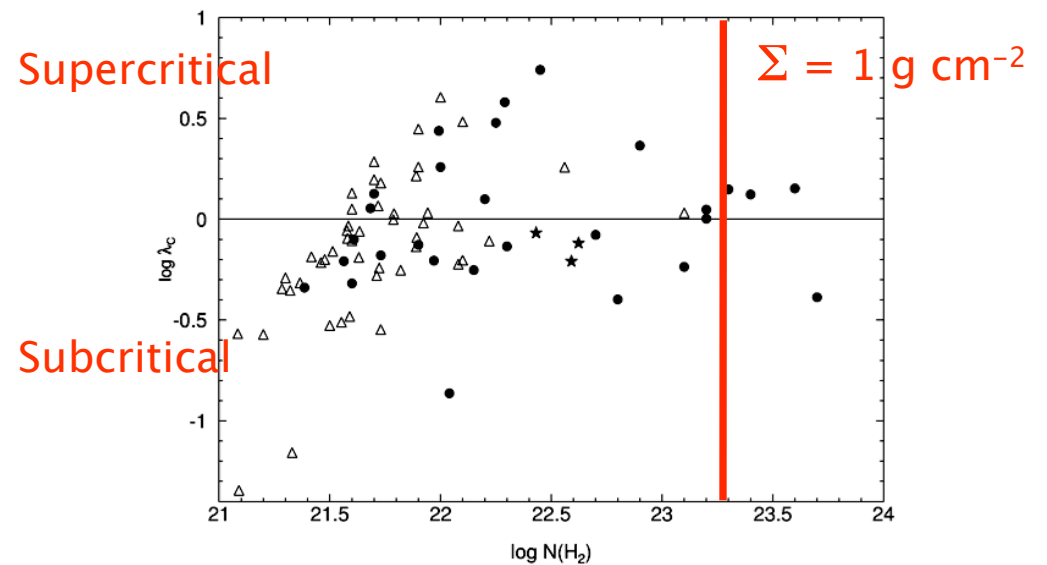


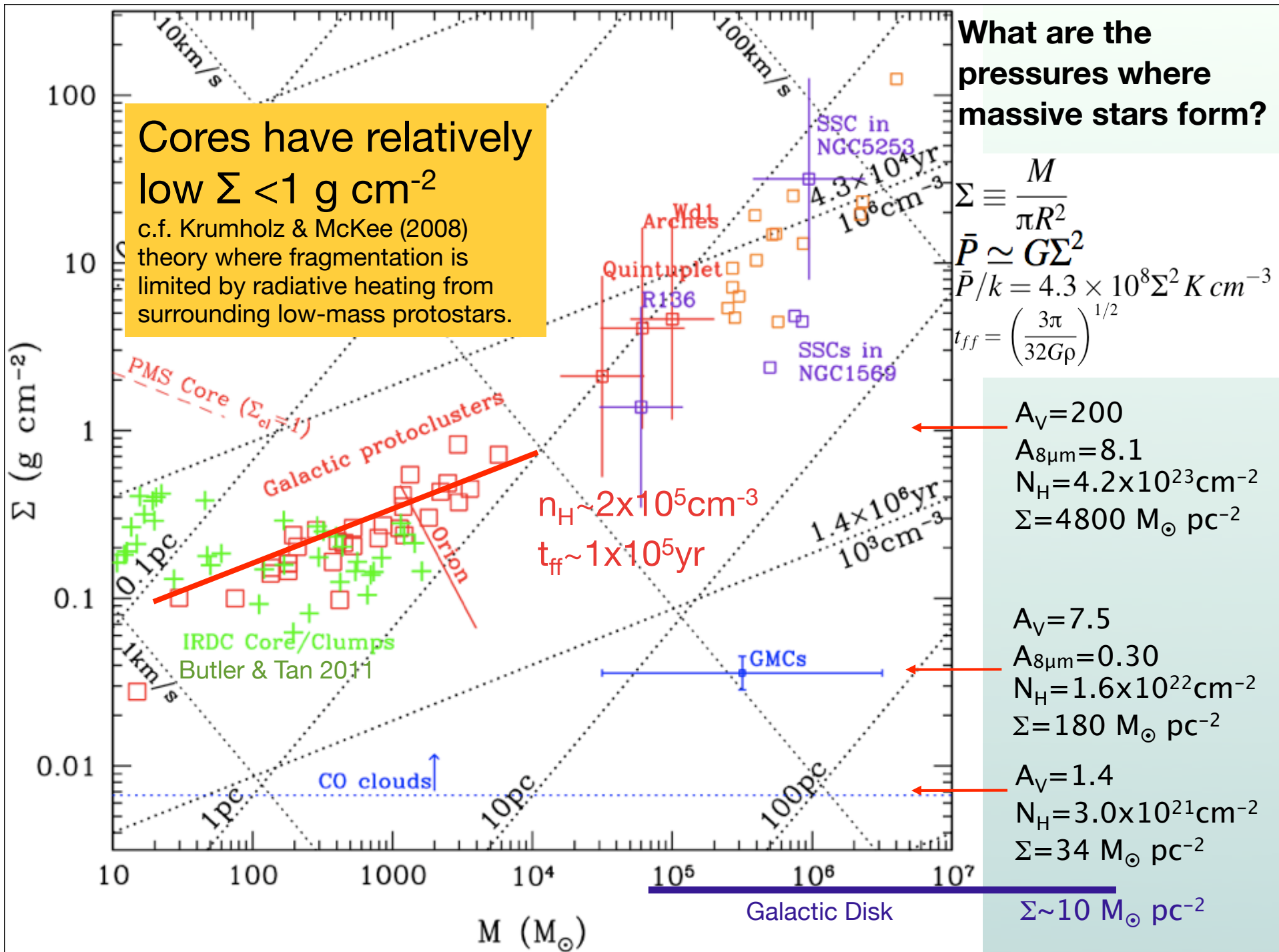
Correlation of field orientations from  $\sim 100$ pc to  $< 1$ pc scales (Hua-bai Li et al. 2009)



Girart et al. (2009)

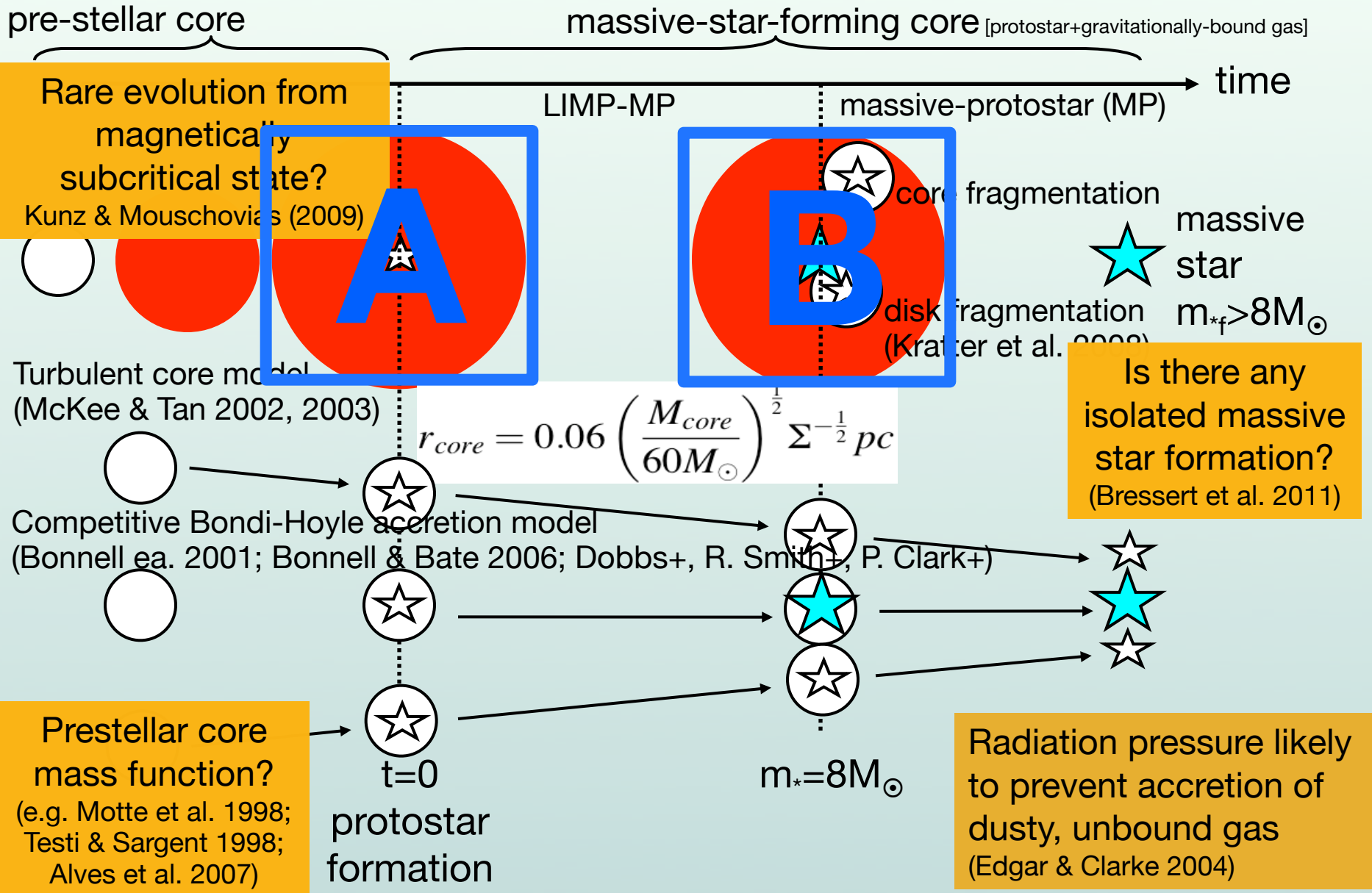
Strength of B-field vs.  $\Sigma$   
(Crutcher 2005;  
Falgarone et al. 2008)





# Schematic Differences Between Massive Star Formation Theories

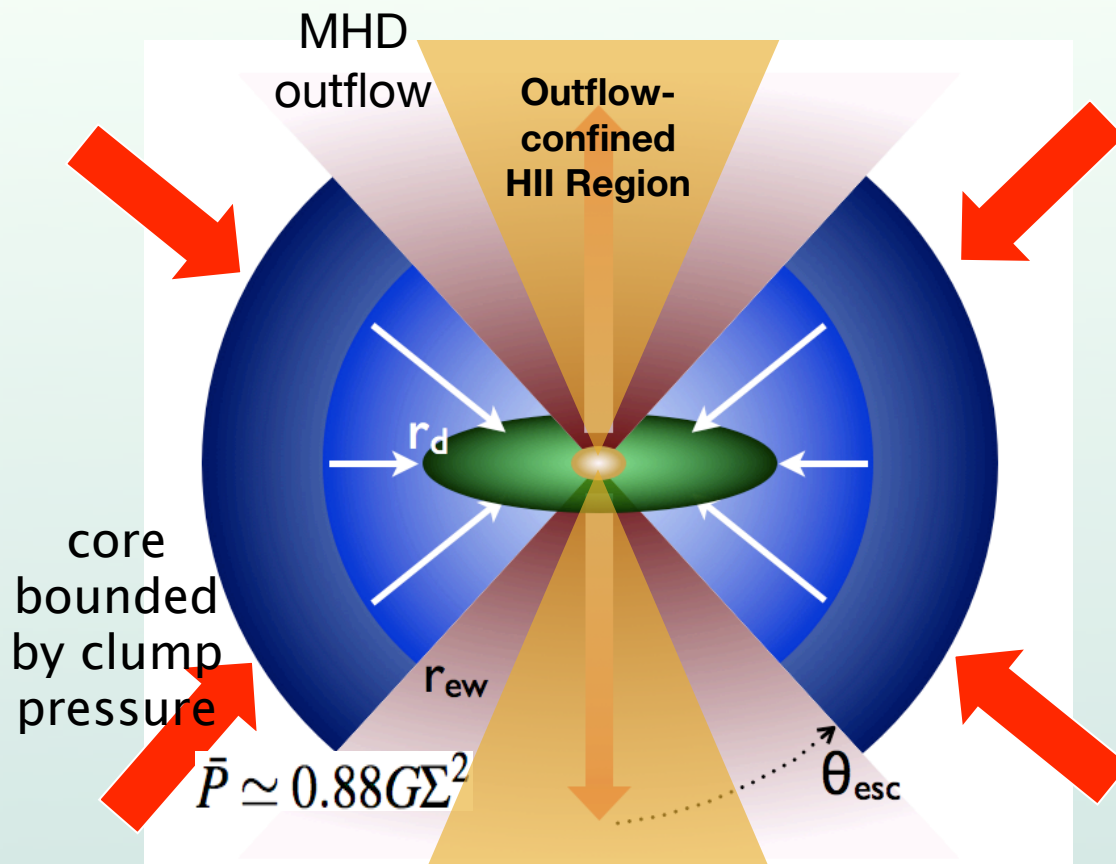
Beuther, Churchwell, McKee, Tan (2007);  
Tan (2008)





# Massive Protostars

MT03, TM, in prep.

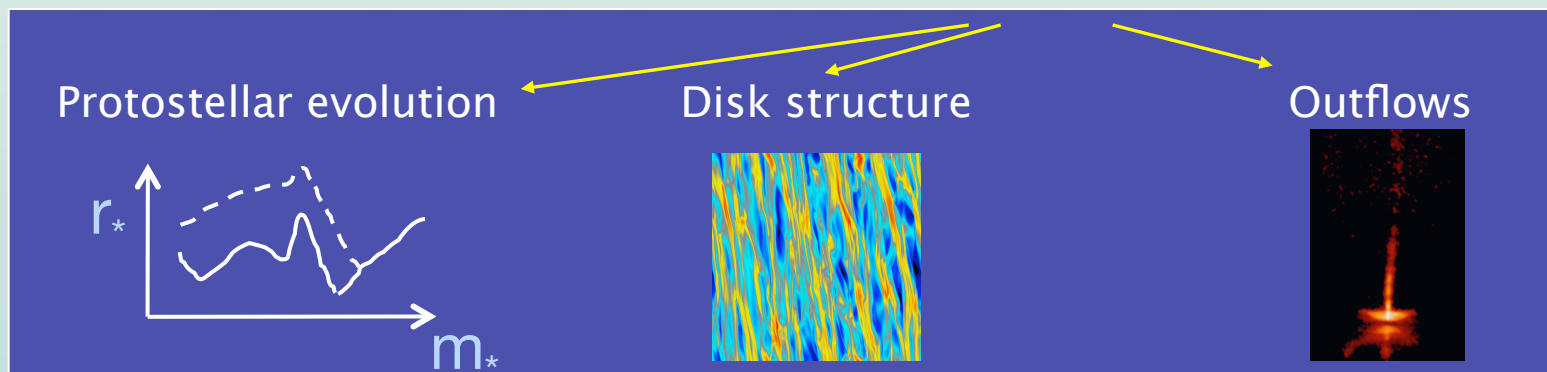


$$r_{core} = 0.06 \left( \frac{M_{core}}{60 M_{\odot}} \right)^{\frac{1}{2}} \Sigma^{-\frac{1}{2}} pc$$

$$r_{disk} = 1200 \frac{\beta}{0.02} \left( \frac{M_{core}}{60 M_{\odot}} \right)^{\frac{1}{2}} \Sigma^{-\frac{1}{2}} AU$$

$$t_{*f} = 1.3 \times 10^5 \left( \frac{M_{core}}{60 M_{\odot}} \right)^{\frac{1}{4}} \Sigma^{-\frac{3}{4}} yr$$

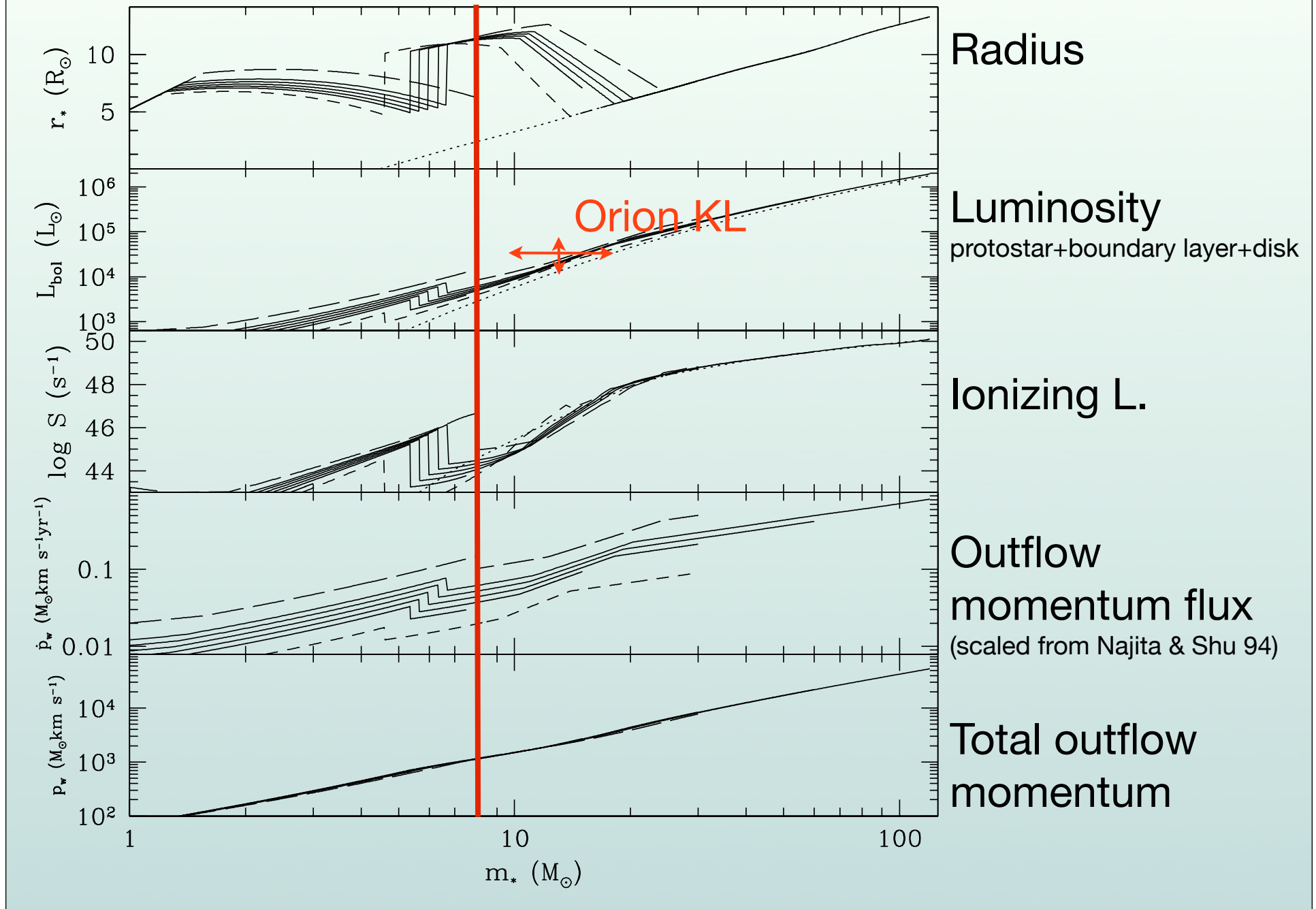
$$\dot{m}_* = 4.6 \times 10^{-4} \left( \frac{M_{core}}{60 M_{\odot}} \right)^{\frac{3}{4}} \Sigma^{\frac{3}{4}} M_{\odot} yr^{-1}$$





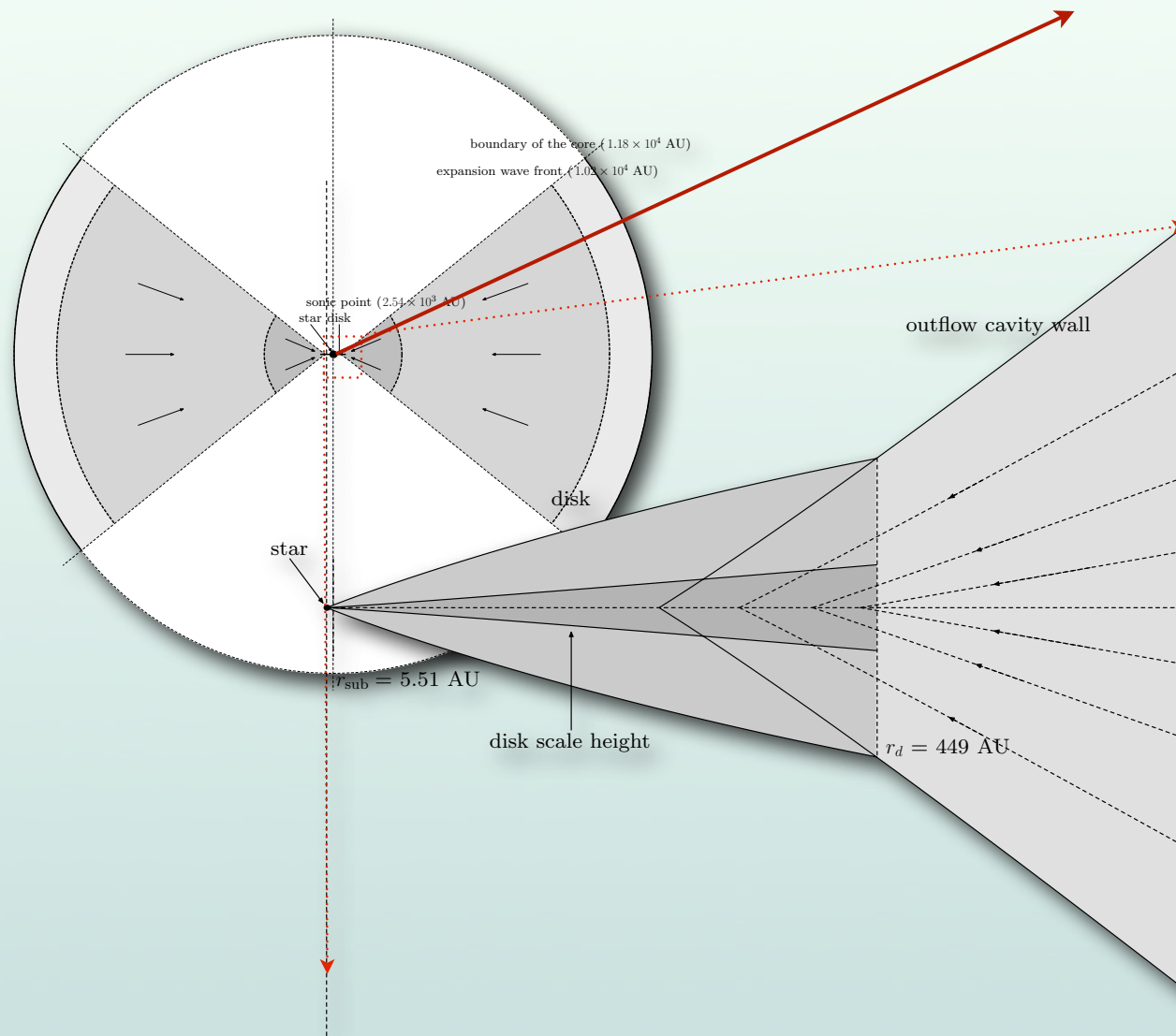
# Protostellar Evolution

[Tan & McKee 2003; see also Hosokawa & Omukai 2009]



# Continuum Radiative Transfer Modeling

Zhang & Tan (2011)



# Radiative Transfer Models

Zhang & Tan (2011)

see also:

Robitaille et al. 2006;

Molinari et al. 2008;

Johnston et al. 2011

Rotation and outflow axis inclined at  $60^\circ$  to line of sight.

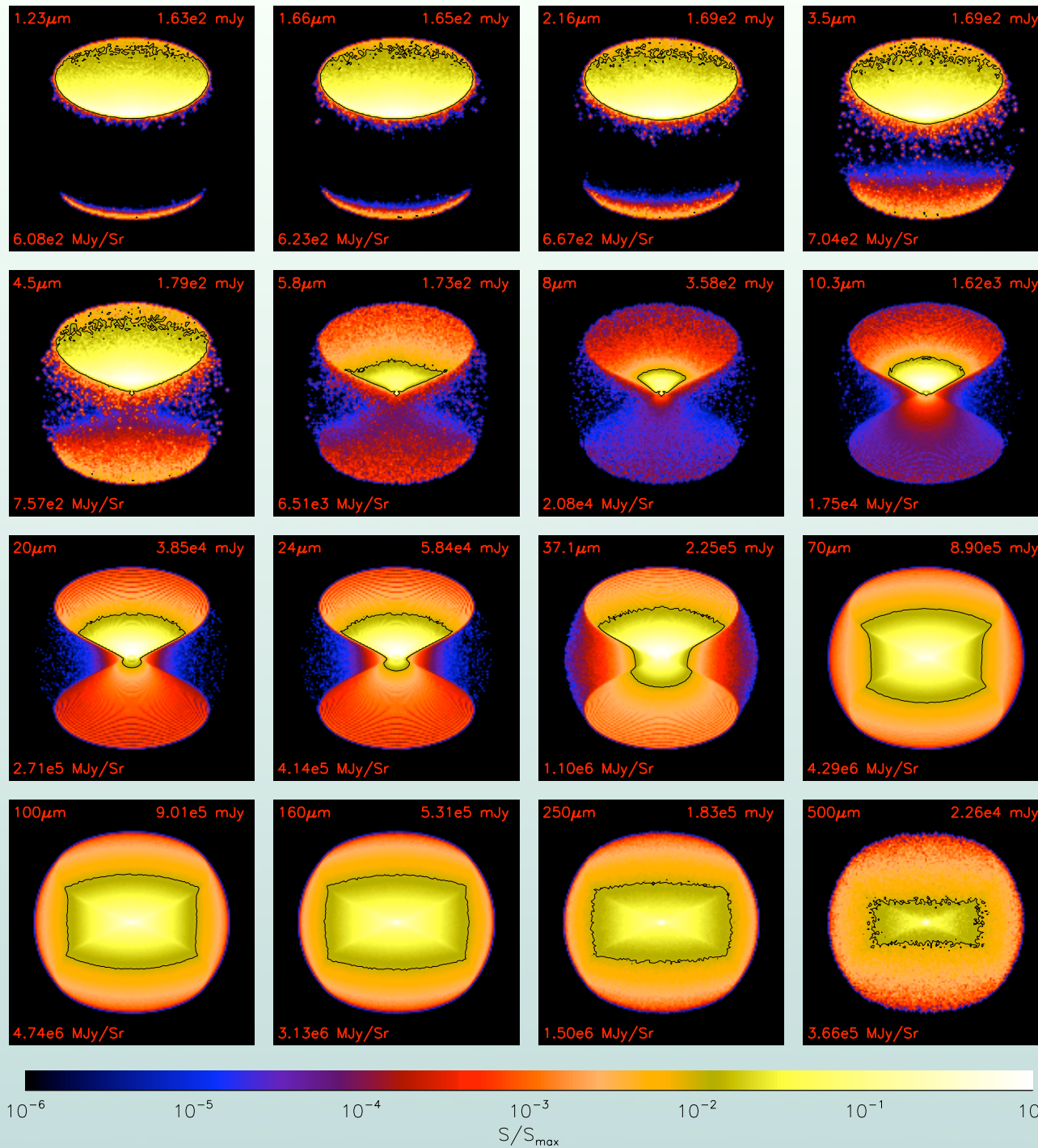
$$\Sigma = 1 \text{ g cm}^{-2}$$

$$M_{\text{core}} = 60 M_{\odot}$$

$$m^* = 8 M_{\odot}$$

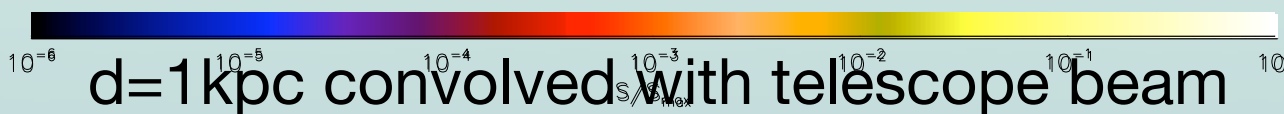
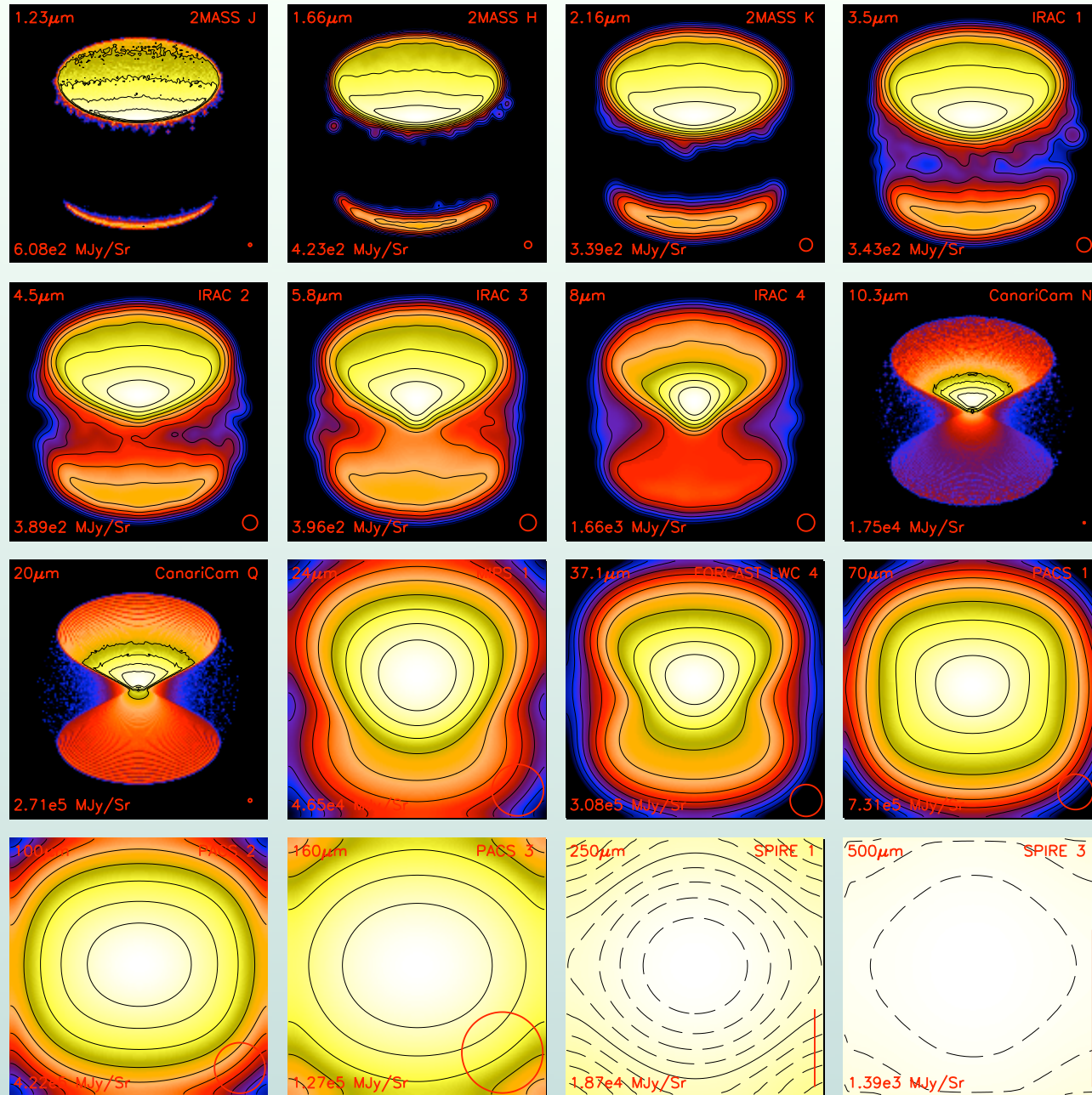
$$m_{\text{disk}} = m^*/3$$

$$L_{\text{bol}} = 6 \times 10^3 L_{\odot}$$



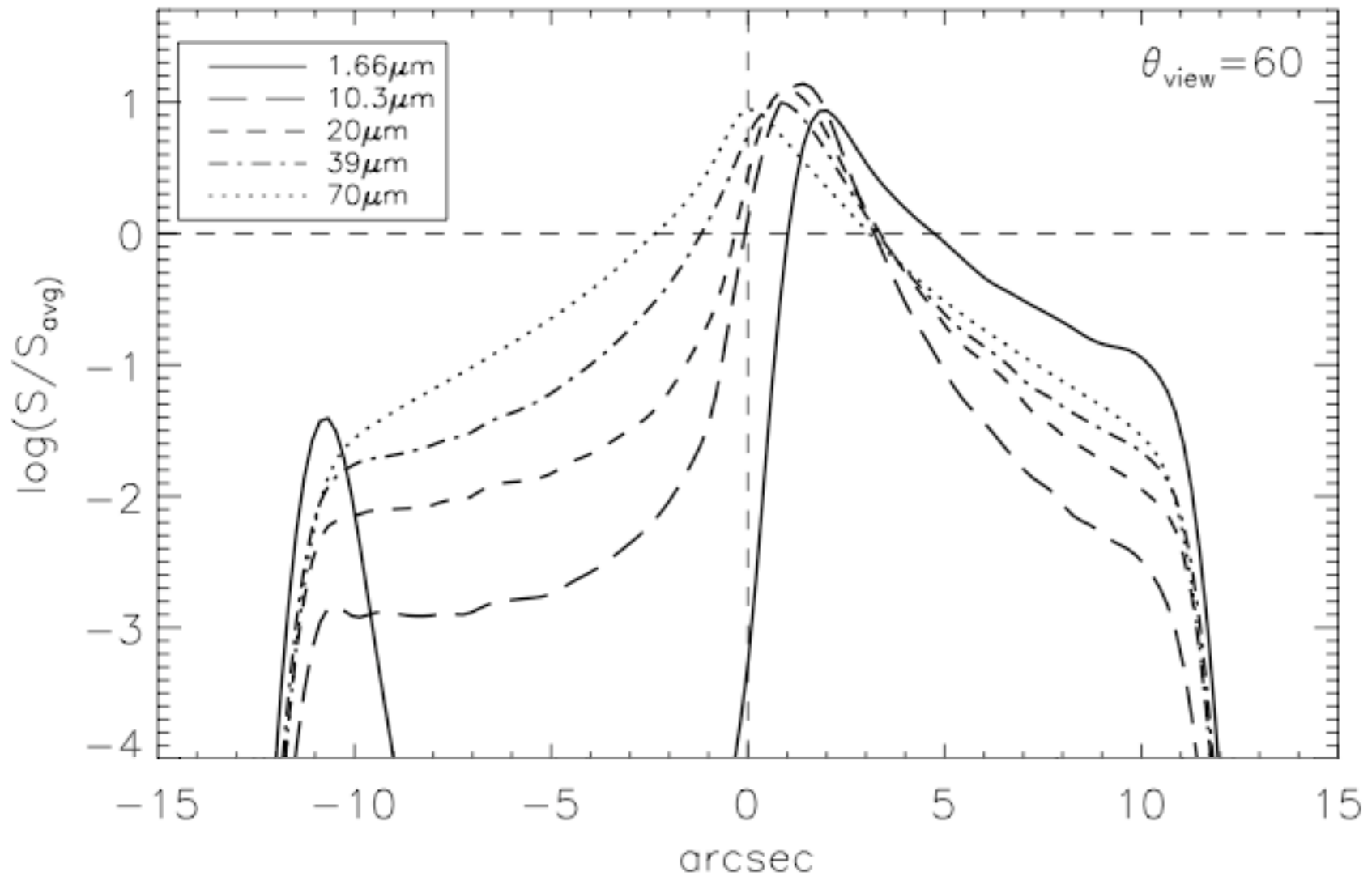
# Radiative Transfer Models

Zhang & Tan (2011)  
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 Rotation and  
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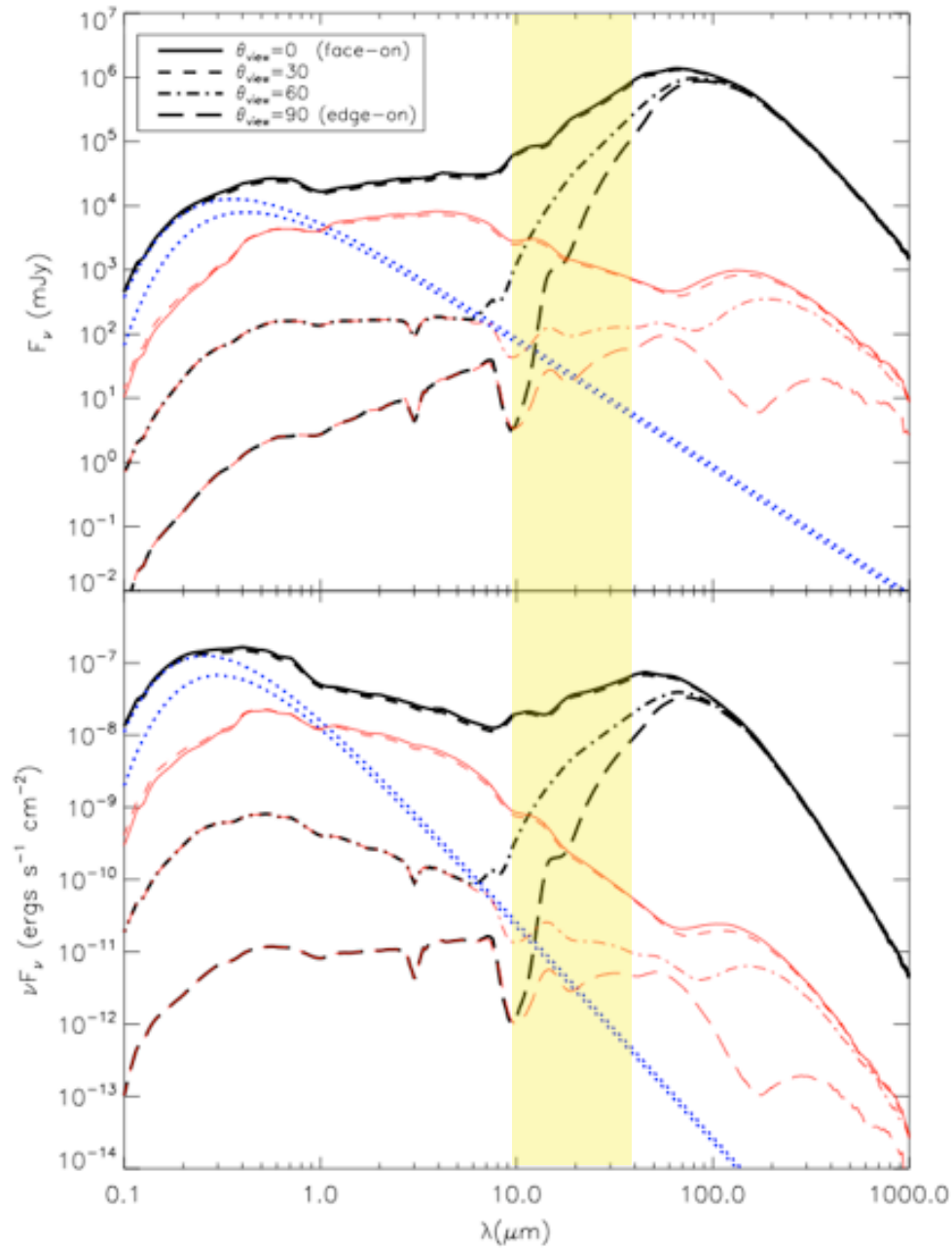
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 $L_{\text{bol}} = 6 \times 10^3 L_{\odot}$

# Flux Profile along Outflow Axis





## SEDs



**Figure 7.** SEDs of the fiducial model (Model 8) at four inclinations, assuming a distance of 1 kpc ( $F_\nu$  in upper panel and  $\nu F_\nu$  in lower panel). The thicker lines are total fluxes and the thinner red lines are only the scattered light. The lower dotted blue line is the input stellar spectrum (blackbody). The upper dotted line is the total luminosity of the star and hot spot emitted as a blackbody.



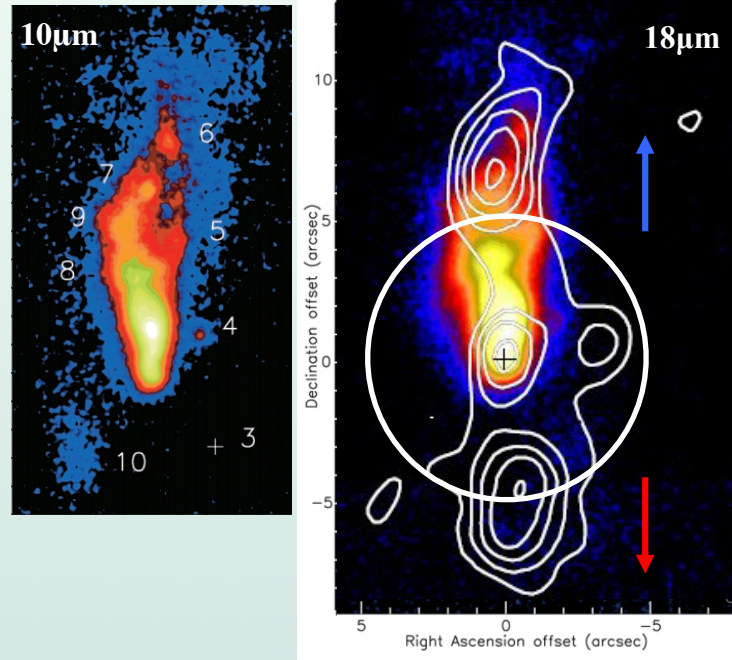
# Mid IR Emission from the Outflow Cavity

G35.2N

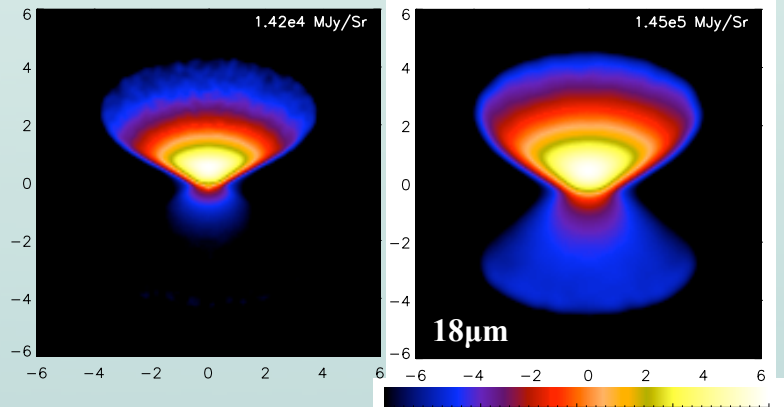
d=2.3 kpc

$L_{\text{MIR}} \sim 1.6 \times 10^3 L_{\odot}$

(De Buizer 2006)



Rotation and outflow axis inclined at  $60^\circ$  to line of sight.  
 $m^* = 8 M_{\odot}$   
 $L_{\text{bol}} = 6 \times 10^3 L_{\odot}$



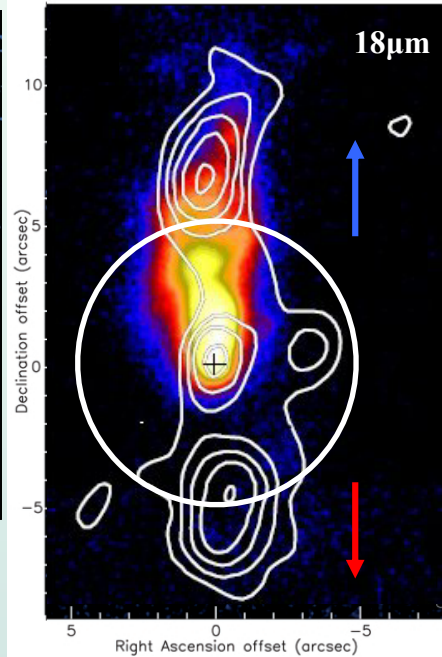
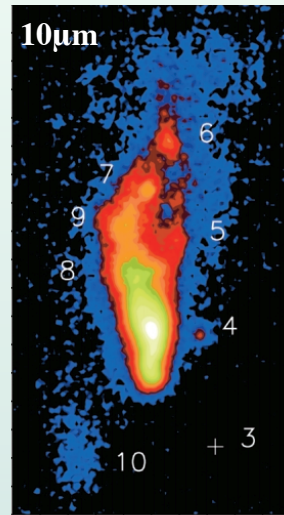
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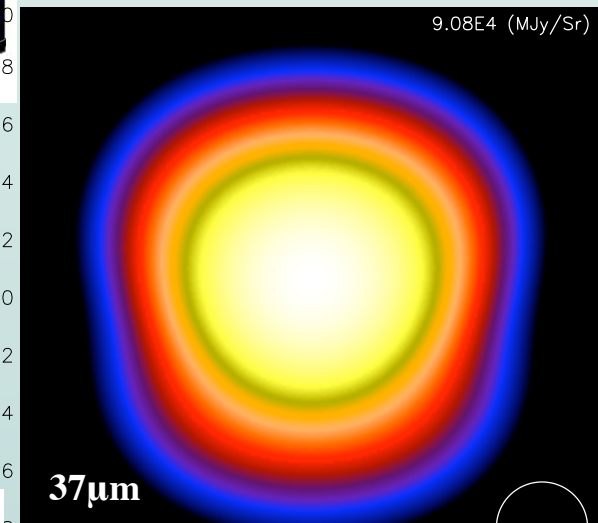
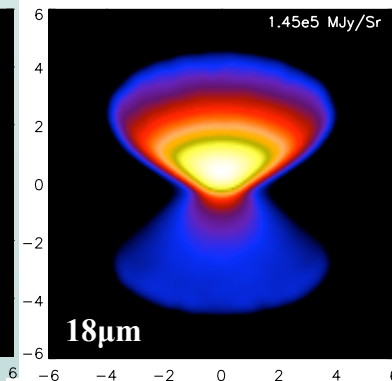
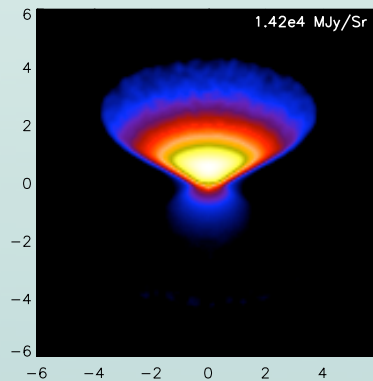
$d=2.3$  kpc

$L_{MIR} \sim 1.6 \times 10^3 L_{\odot}$

(De Buizer 2006)



Rotation and outflow axis inclined at  $60^\circ$  to line of sight.  
 $m^* = 8 M_{\odot}$   
 $L_{bol} = 6 \times 10^3 L_{\odot}$



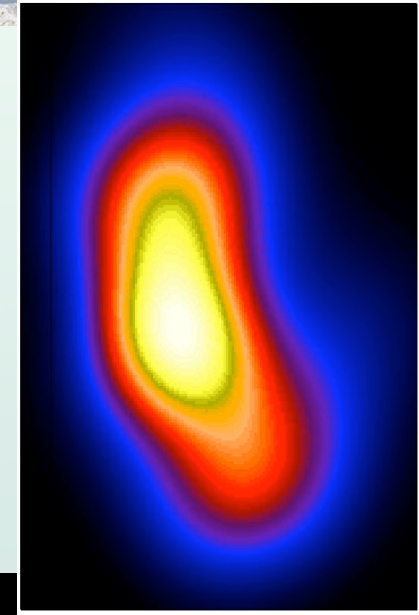
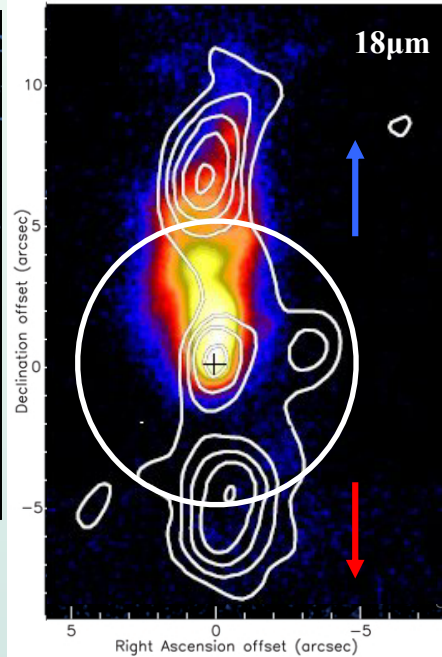
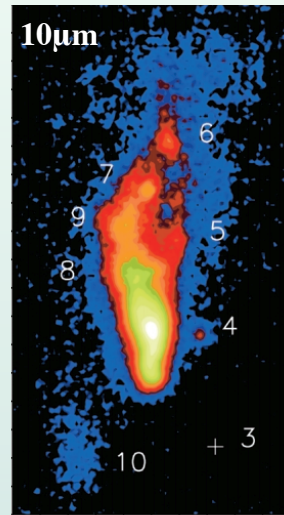
# Mid IR Emission from the Outflow Cavity

G35.2N

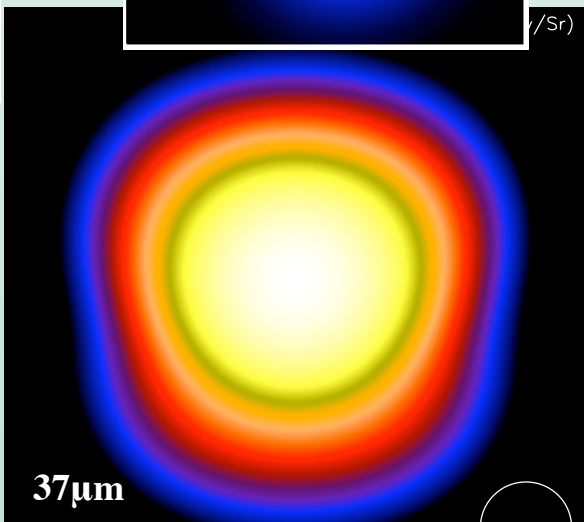
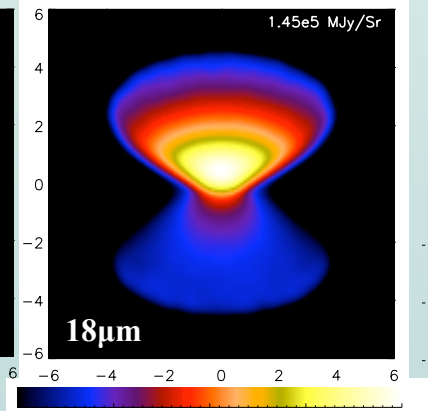
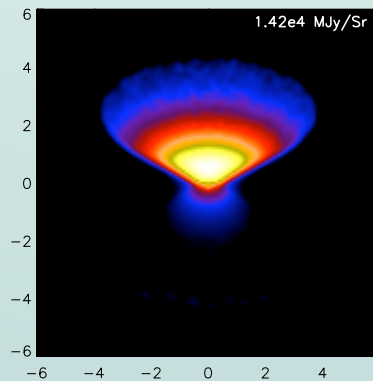
$d=2.3$  kpc

$L_{MIR} \sim 1.6 \times 10^3 L_{\odot}$

(De Buizer 2006)

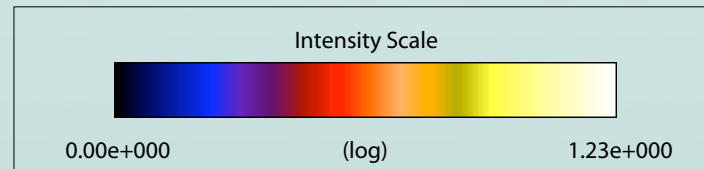
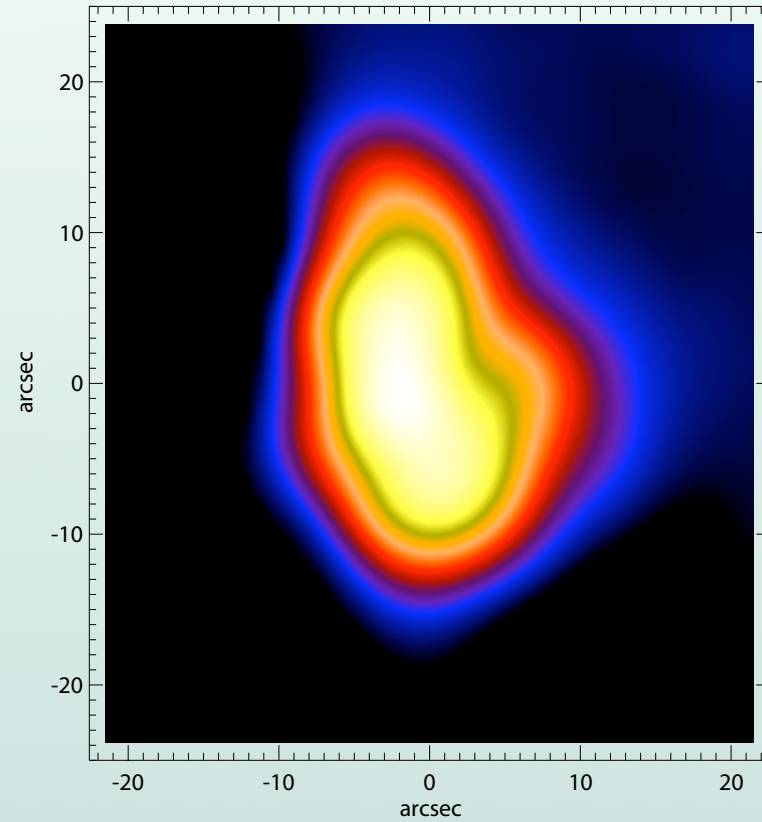
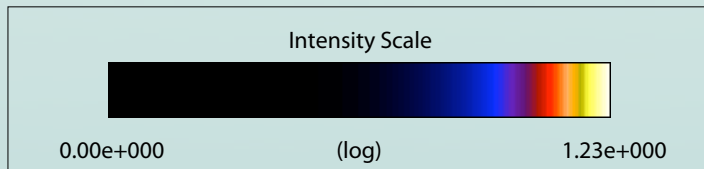
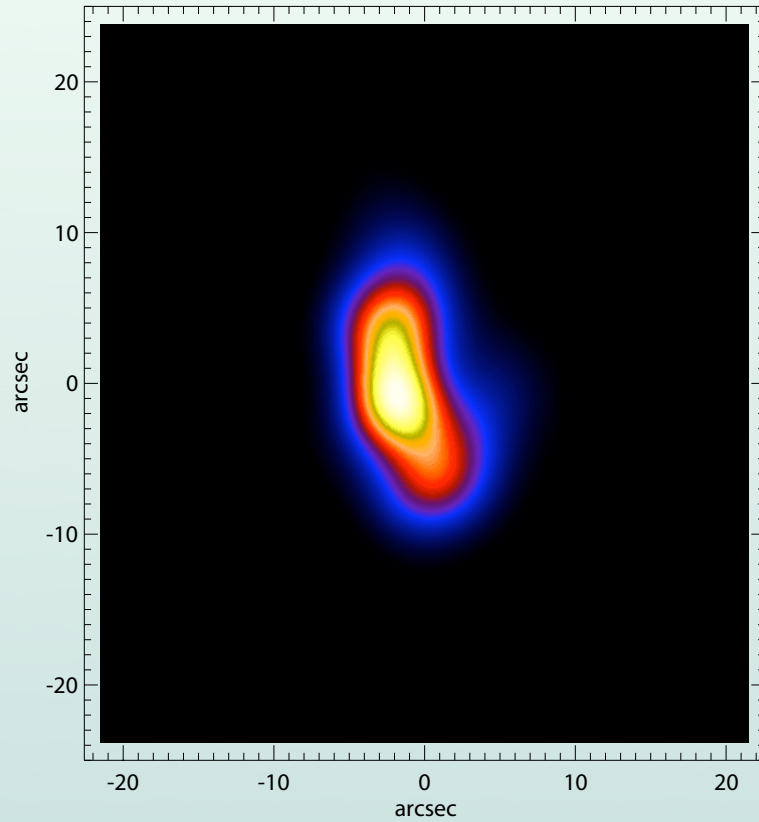


Rotation and outflow axis inclined at  $60^\circ$  to line of sight.  
 $m^* = 8 M_{\odot}$   
 $L_{bol} = 6 \times 10^3 L_{\odot}$

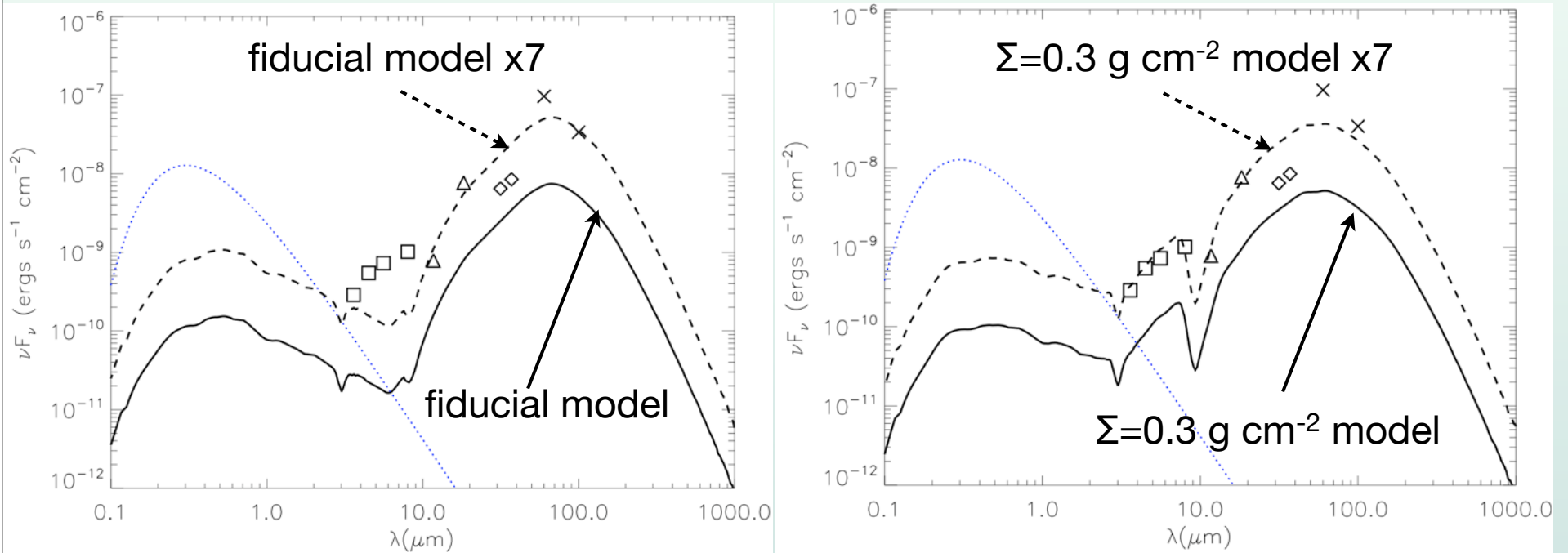


# Extended Faint Emission

G35 (SOFIA 37 micron)



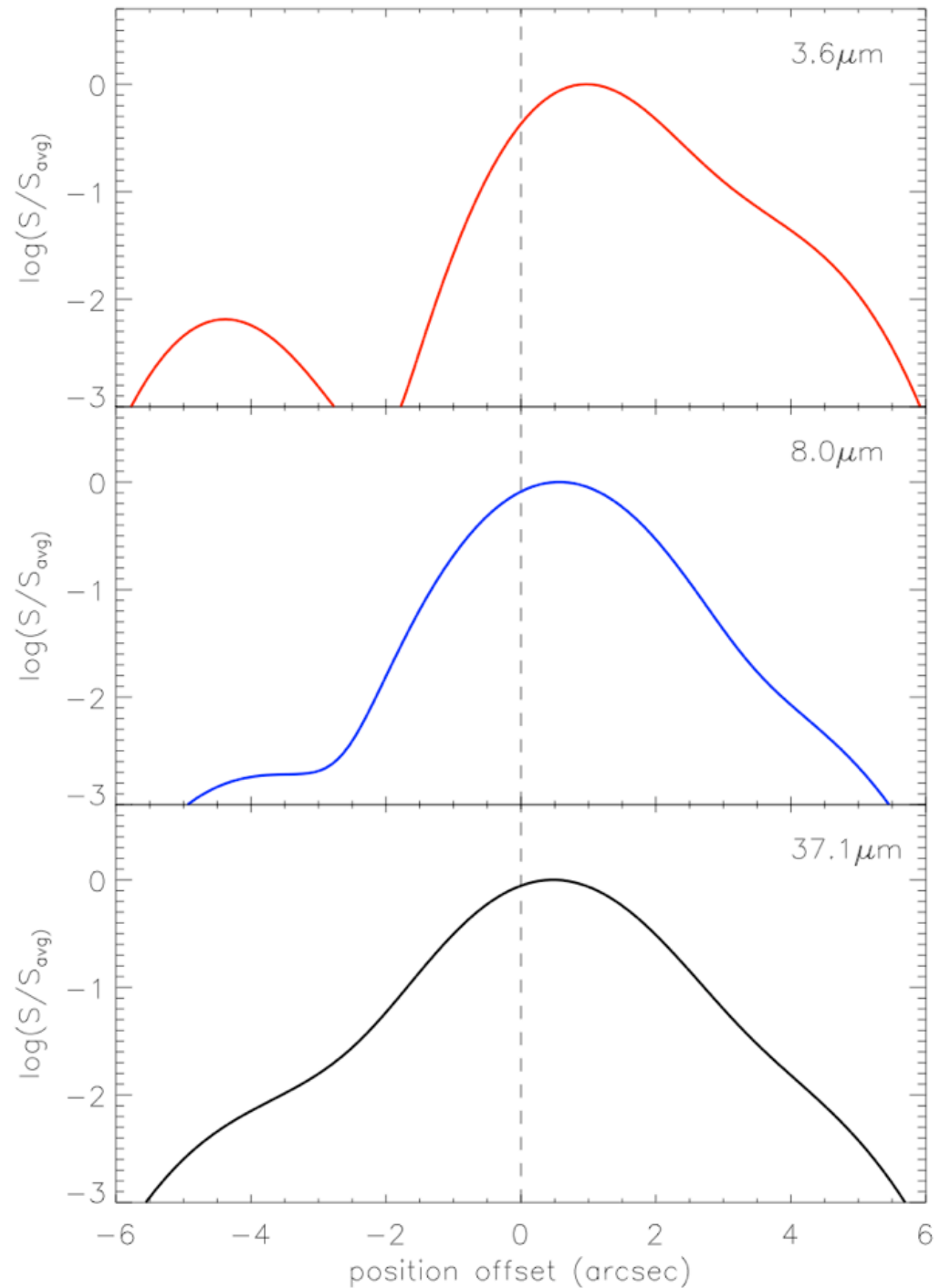
# SED of G35.2N





# Flux Profile Along Outflow Axis

For fiducial model (solid lines) at 2.3kpc, convolved with IRAC and SOFIA beams.

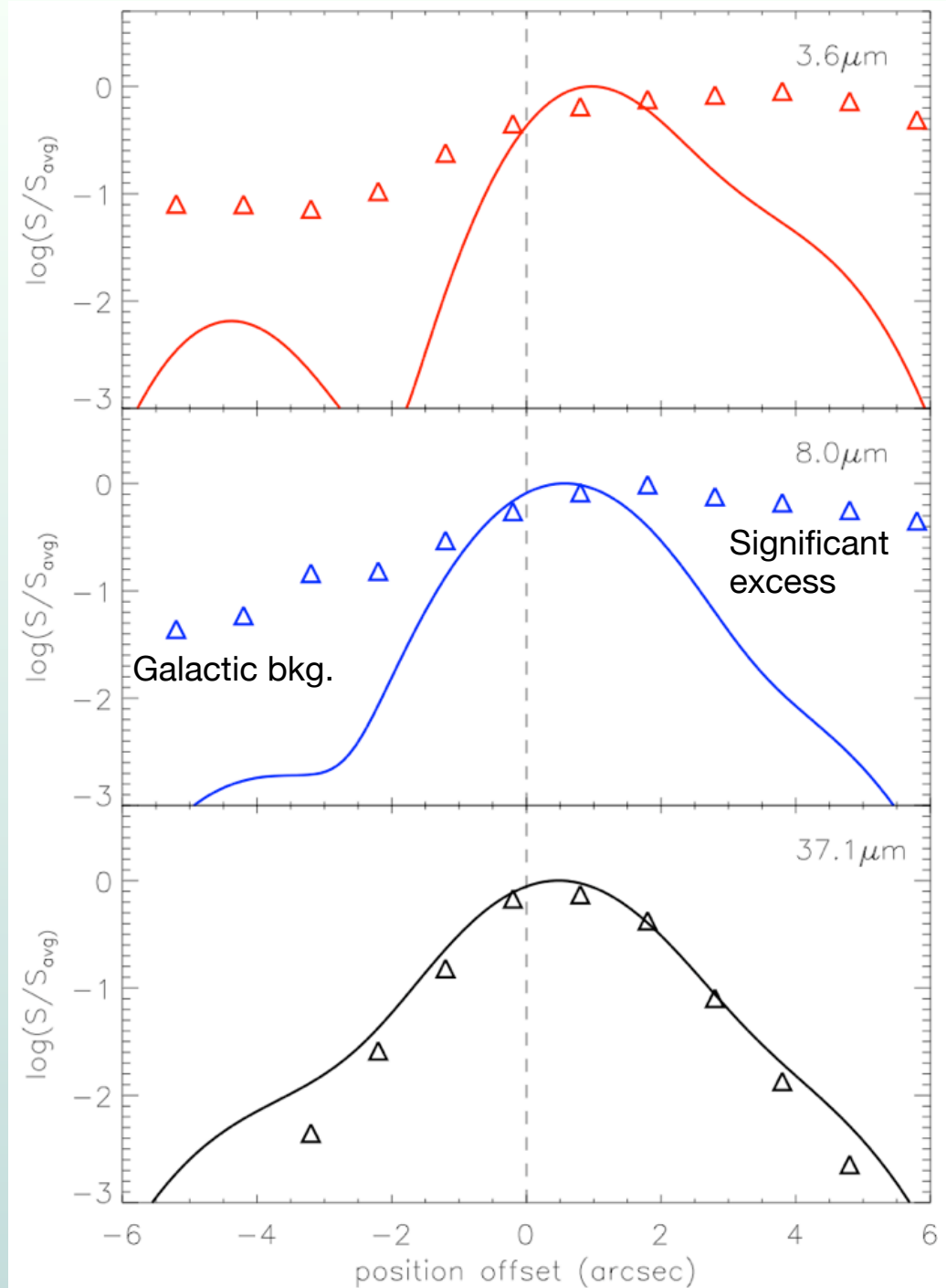




# Flux Profile Along Outflow Axis

For fiducial model (solid lines) at 2.3kpc, convolved with IRAC and SOFIA beams.

Triangles show the observed profiles.

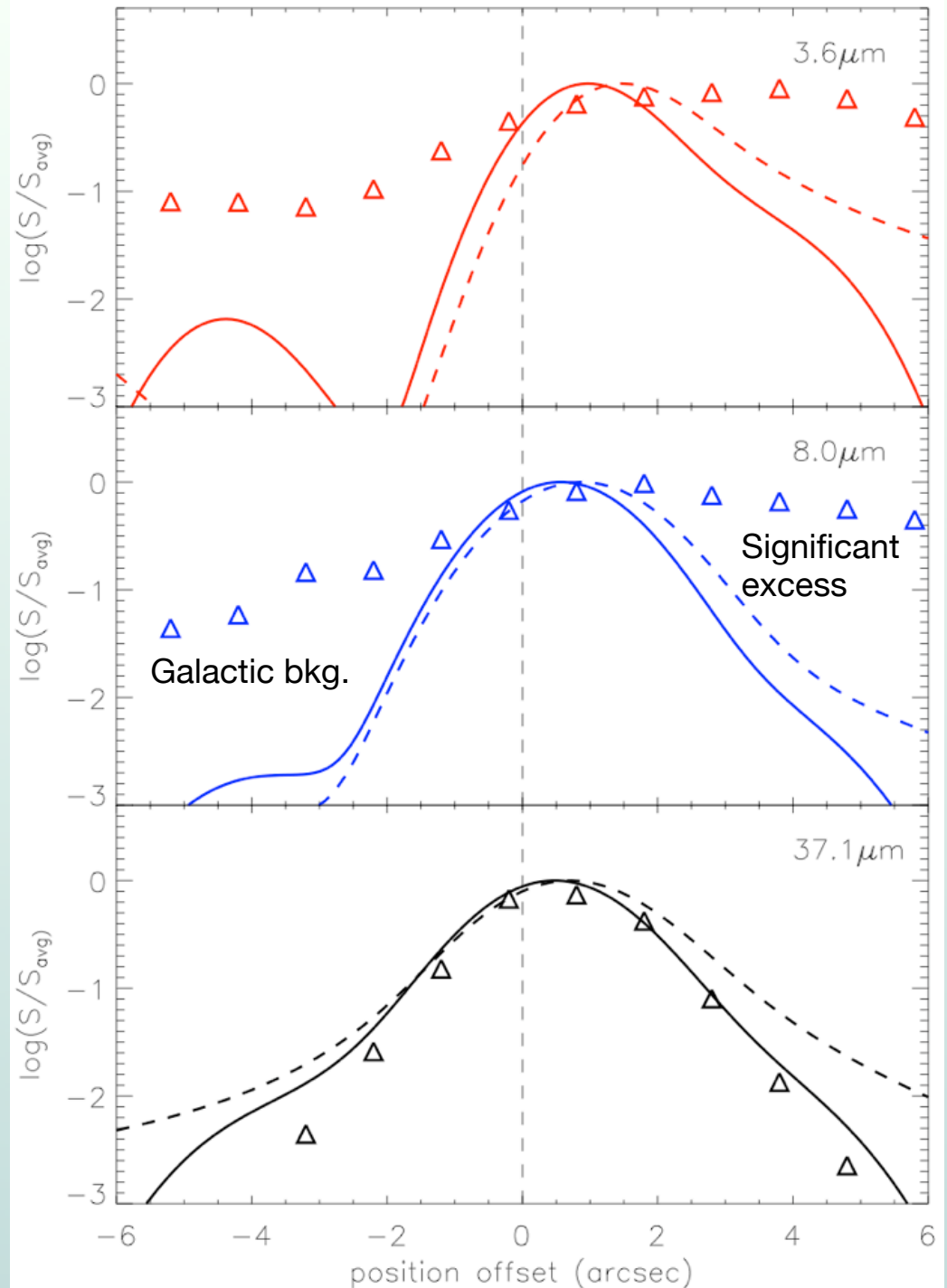


# Flux Profile Along Outflow Axis

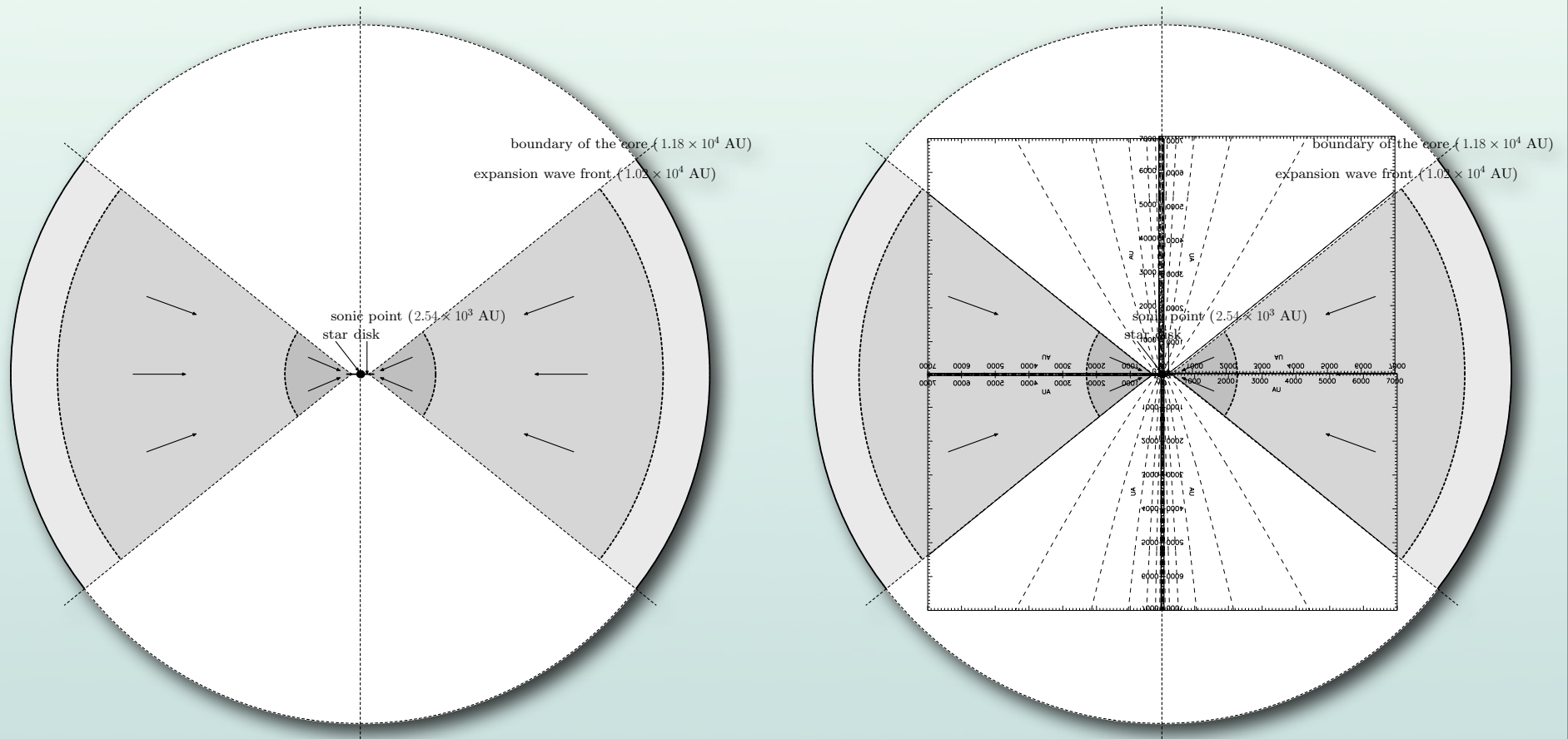
For fiducial model (solid lines) at 2.3kpc, convolved with IRAC and SOFIA beams.

Triangles show the observed profiles.

Dashed lines show  $\Sigma = 0.3 \text{ g cm}^{-2}$  model.

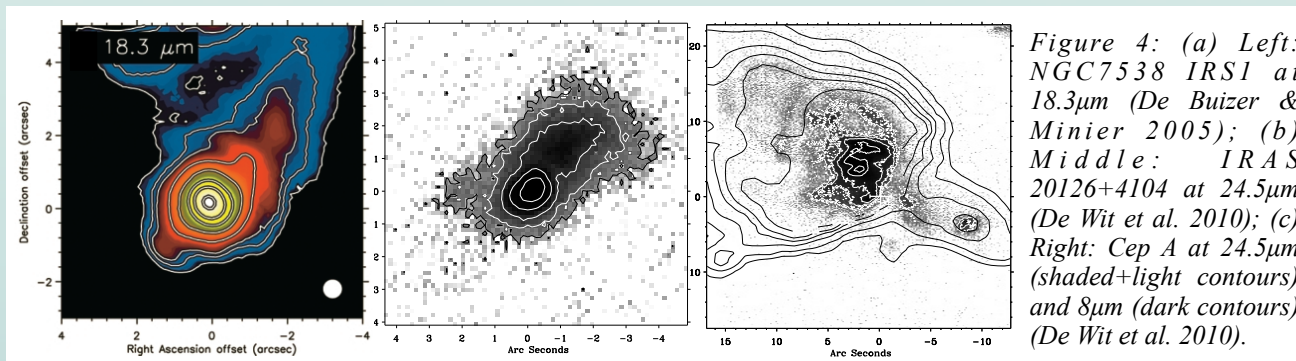
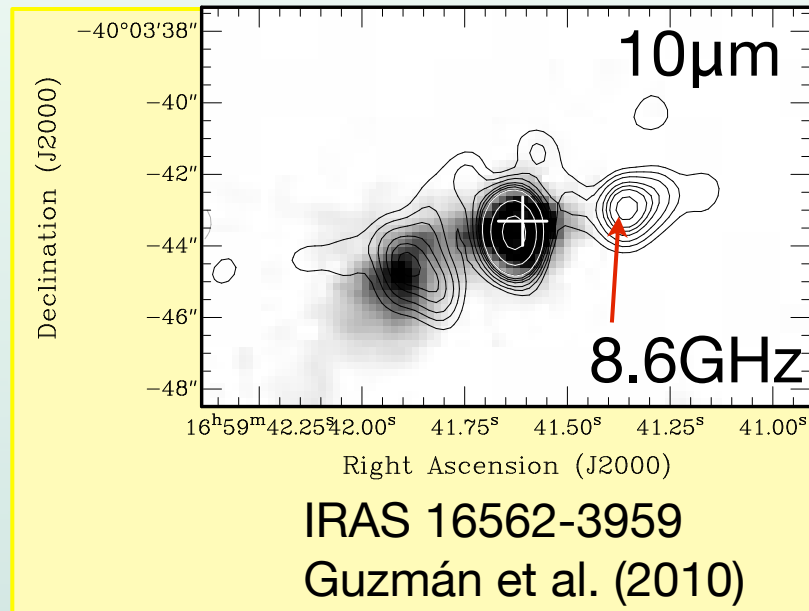


# Including Dusty Outflow May Help Resolve the Discrepancy



# G35.2N is not alone

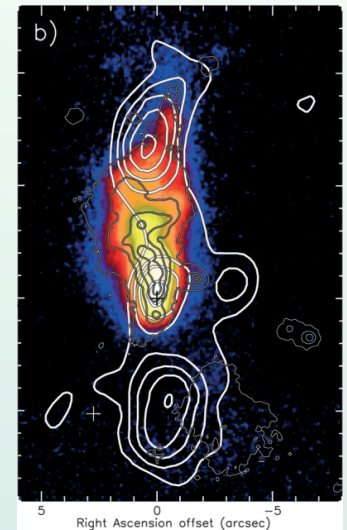
A number of other ionized HCHIIIs seen in nearby sources  
(e.g. van der Tak & Menten 2005; Guzmán et al. 2010)



# Conclusions: The Formation of Massive Stars

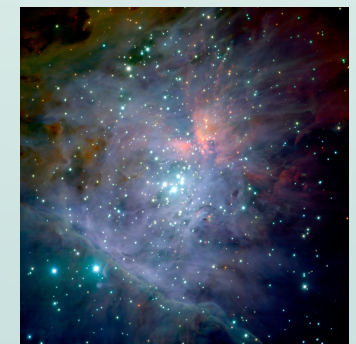
## Massive star formation:

- Sometimes appears to be a scaled-up version of low-mass star formation: massive starless cores, ordered B-fields, rotating toroids, some claims of disks, collimated outflows & outflow-confined HII regions.
- “Core Accretion Model” invokes turbulent/magnetic support of rare, massive cores & ~pressure equilibrium w/ surrounding clump.
- Fragmentation inside core may be limited by B-fields, rather than radiative heating.
- Detailed Radiative Transfer modeling and comparison to MIR/FIR observations, including SOFIA-FORCAST, can provide powerful constraints on the Core Accretion model.



## Star Cluster Formation:

- SFE per free-fall time from turbulent gas is low, ~few%.
- Formation of “rich”, bound (total SFE >20%) star clusters should take many free-fall times, i.e. >~1 Myr.
- Clump reaches ~equilibrium; turbulence maintained by outflows.



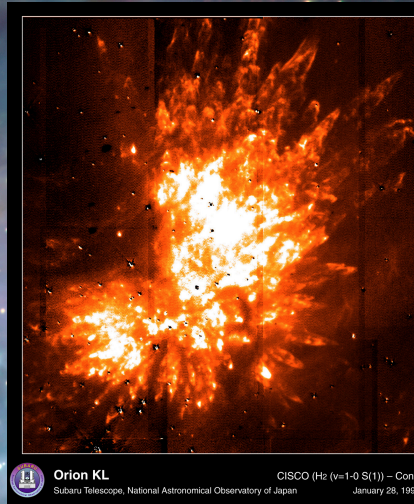


# Orion



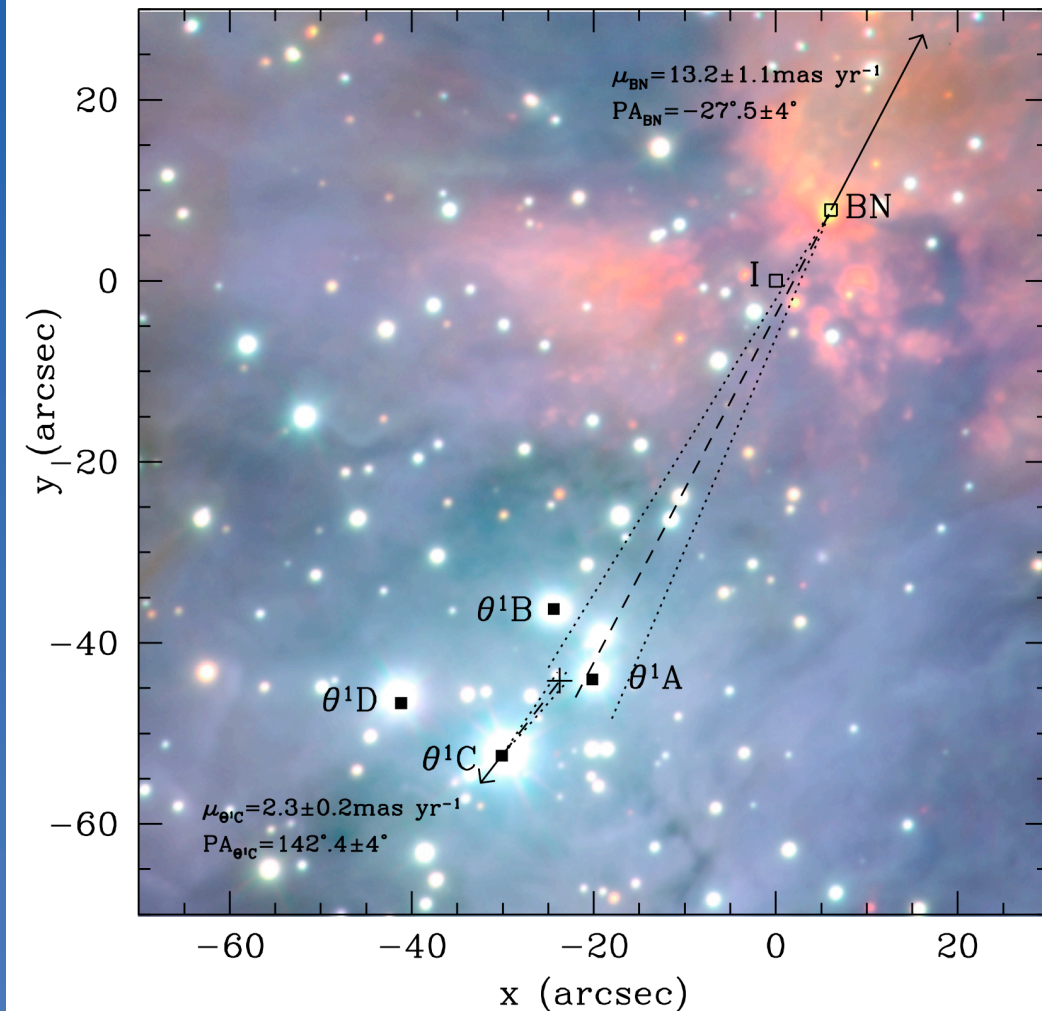


# Orion



Close passage of BN within  $\sim 300$  AU of the KL protostar (source I), 500 years ago, would have tidally perturbed the accretion disk, enhanced accretion, and thus enhanced the outflow (Tan 2004, 2008).

cf. Bally & Zinnecker (2005), who argue for ejection of BN from source I.



Cluster frame proper motions of BN (Gomez et al. 2008; Tan 2008) and  $\theta^1C$  (van Altena et al. 1988), tracing back to a common origin about 4500 years ago, shown by the cross.



# Constraining source I from reflected NIR spectra

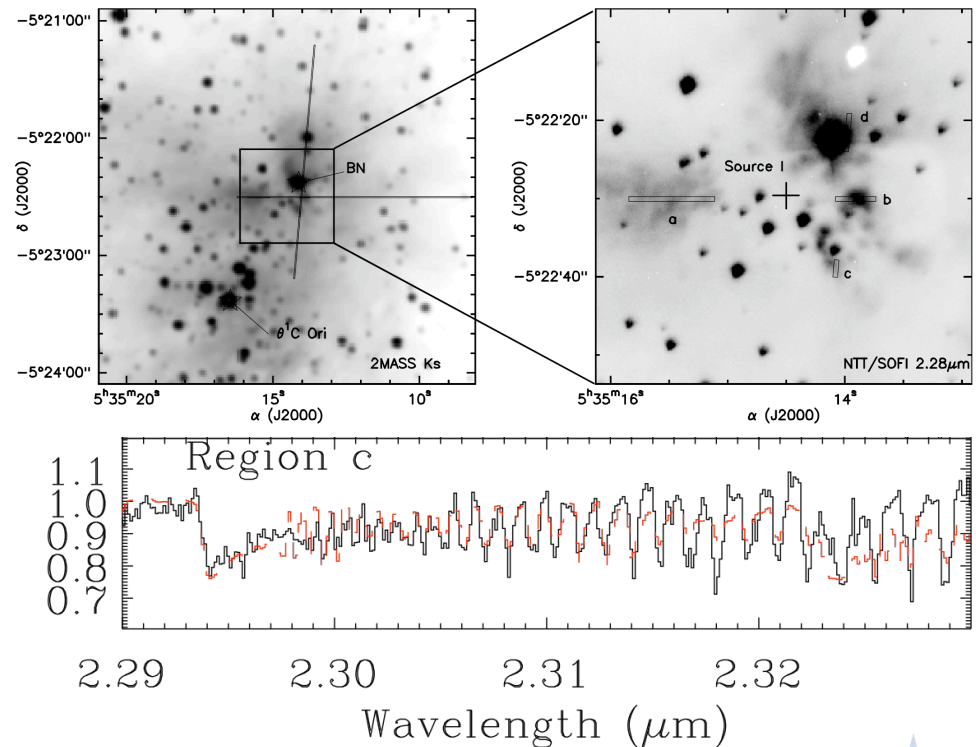
Testi, Tan, Palla 2010

NIR light from source I escapes along the outflow cavity aligned NW-SE (Minchin et al. 1991). Some of this light is scattered towards us by dust, and the spectrum can be used to constrain properties of the protostar.

Morino et al. (1998) found  $T_{\text{phot}} \sim 4000\text{K}$

If from a massive protostar, requires large size ( $\sim 300 R_{\text{sun}}$ ). Such sizes can be achieved if the average accretion rate is  $> \sim 4 \times 10^{-3} M_{\text{sun}}/\text{yr}$  (Hosokawa & Omukai 2009),  $\gg$  than expected from the Core Accretion Model (McKee & Tan 2003).

Or, the large size may indicate the star is rotating close to break-up, which would predict large rotational broadening of the spectral features.



We observed the NIR reflection spectrum with the VLT at higher spectral resolution than Morino et al., finding ID velocity dispersions of  $\sim < 30$  km/s.

We expect most NIR emission comes from the disk.

The narrow profiles constrain protostar + disk models: a  $10 M_{\text{sun}}$  protostar with a large current accretion rate of  $\sim \text{few} \times 10^{-3} M_{\text{sun}}/\text{yr}$  is favored. Such a large accretion rate may be the result of the recent perturbation of the disk by close passage of BN.