



thanks to ...



... people in the star formation group at Heidelberg University:

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... many collaborators abroad!







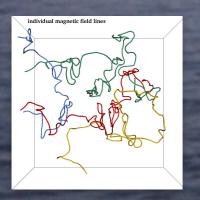






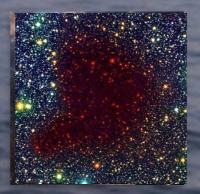
European Research Council

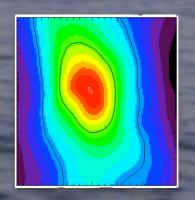
theoretical perspectives

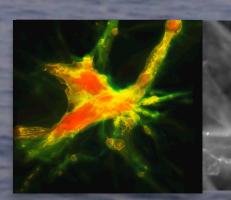












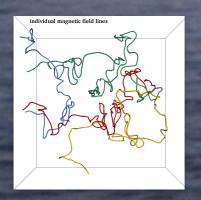


Ralf Klessen

Zentrum für Astronomie der Universität Heidelberg tut für Theoretische Astrophysik

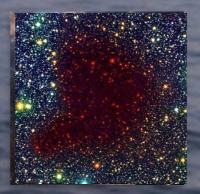


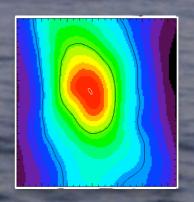
ISM amics & star formation

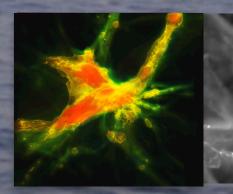














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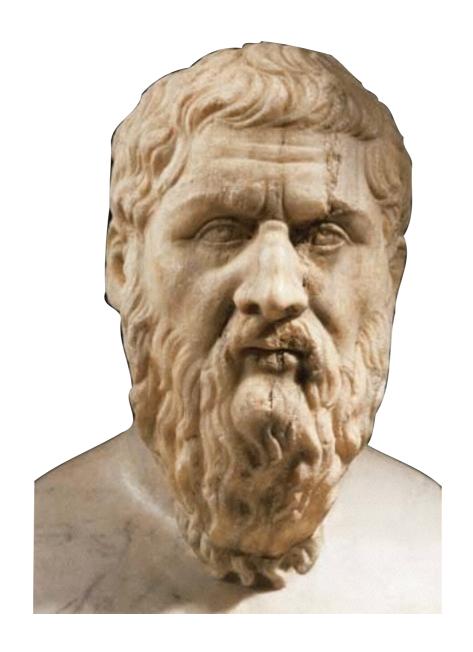






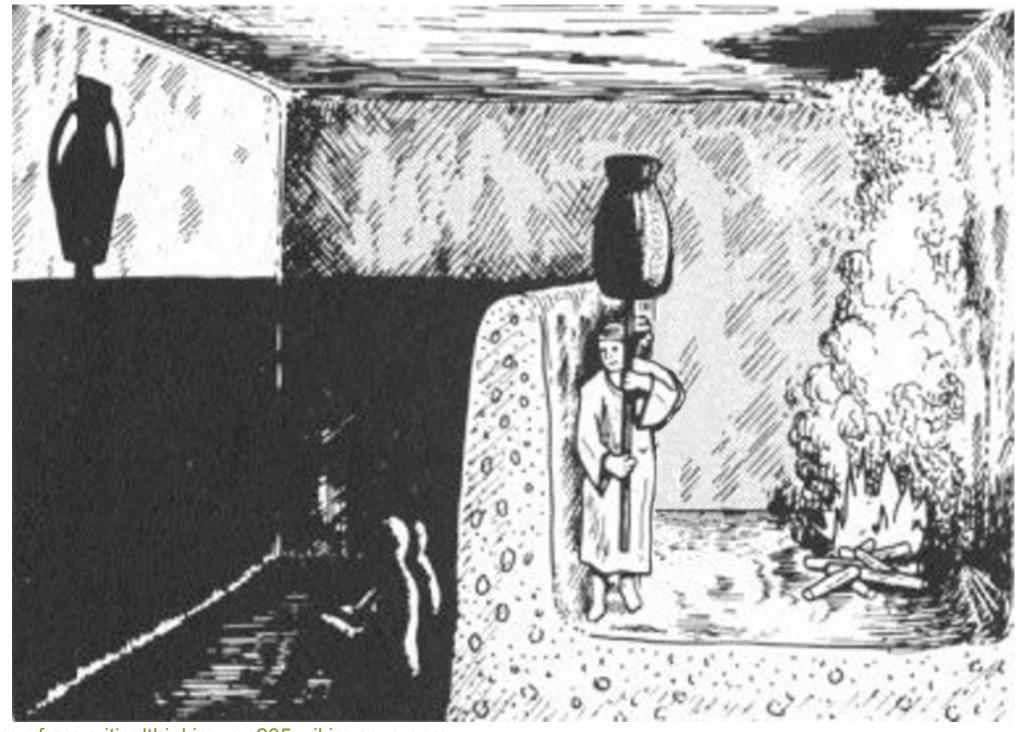






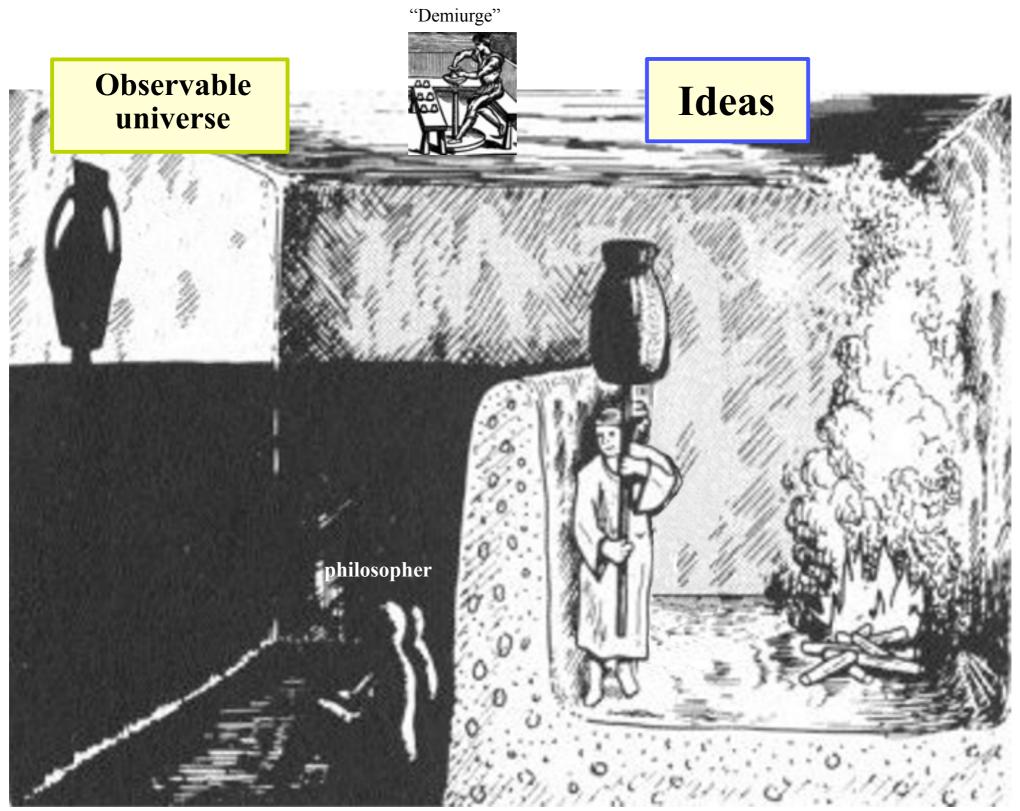
Platon 428/427–348/347 BC

Plato's allegory of the cave*

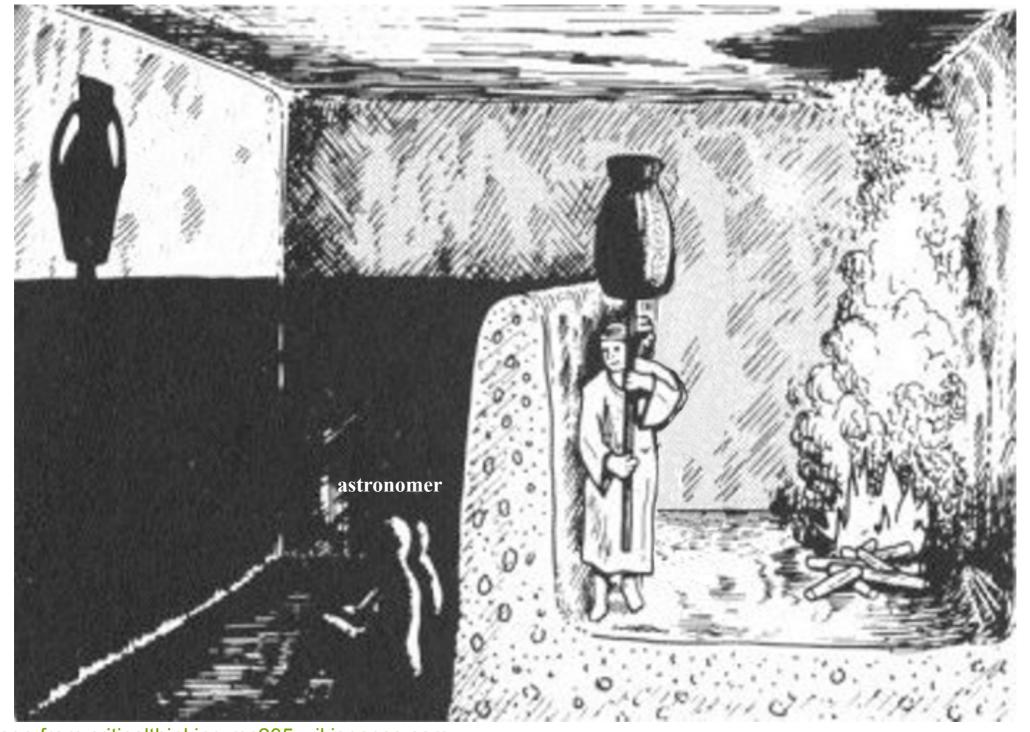


^{*} The Republic (514a-520a)

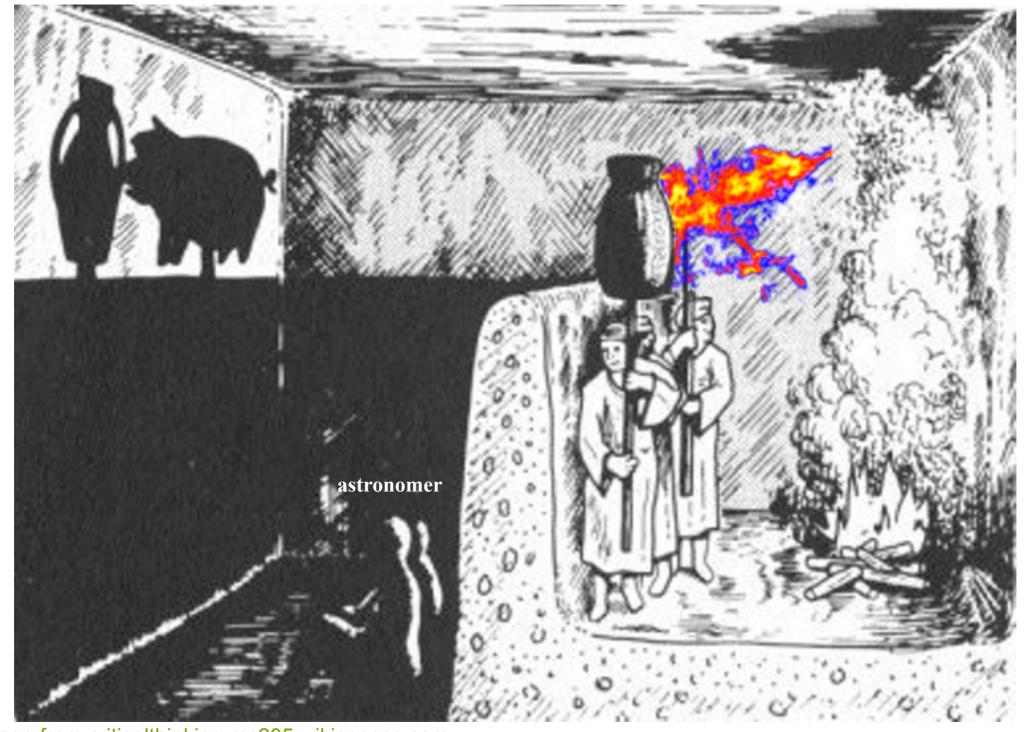
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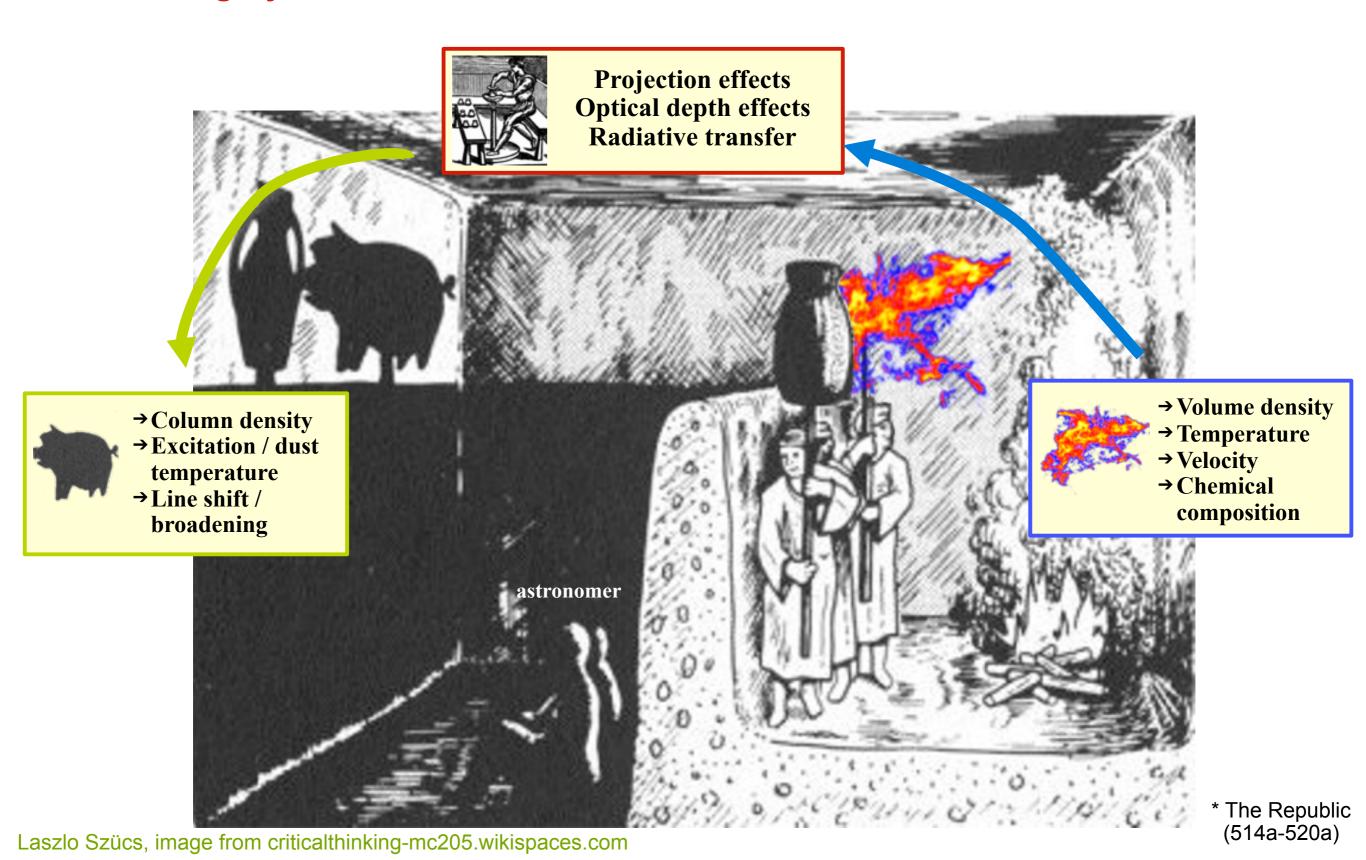
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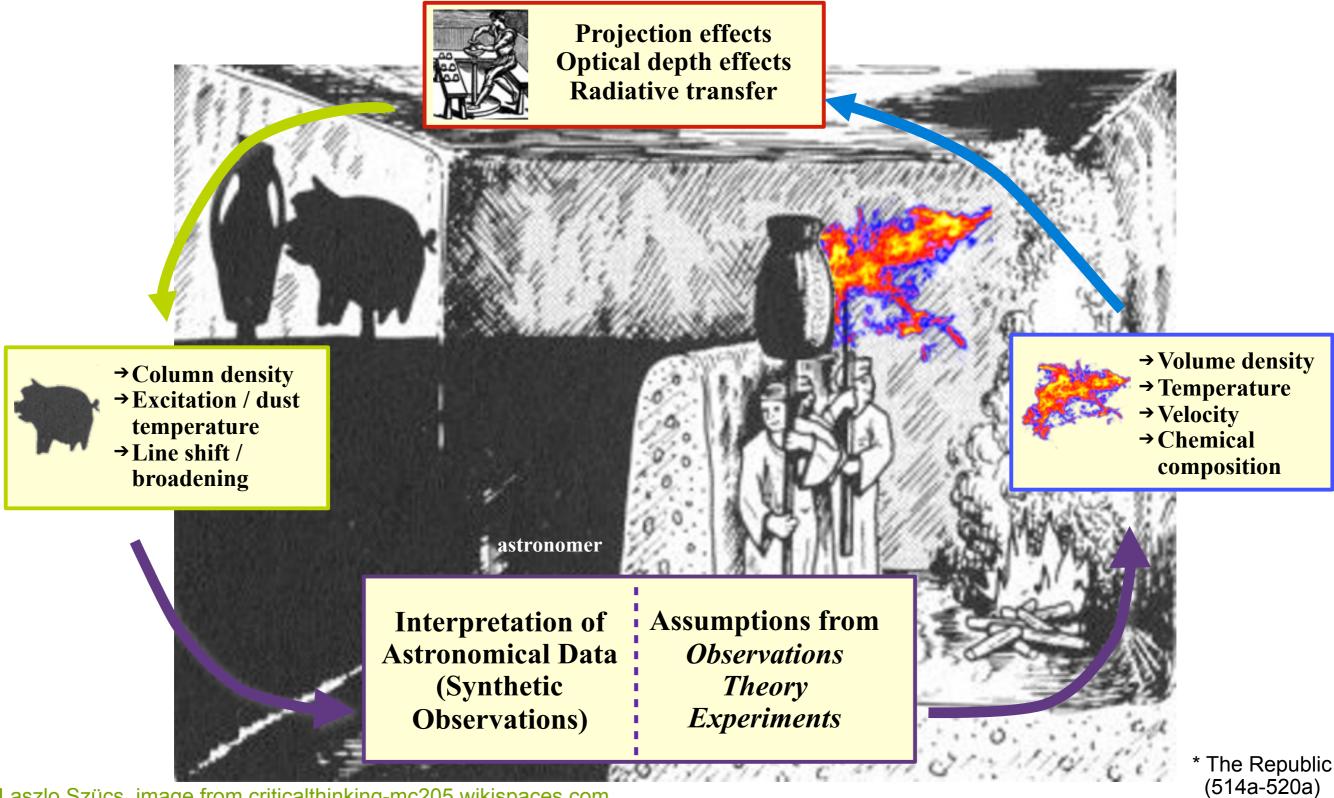


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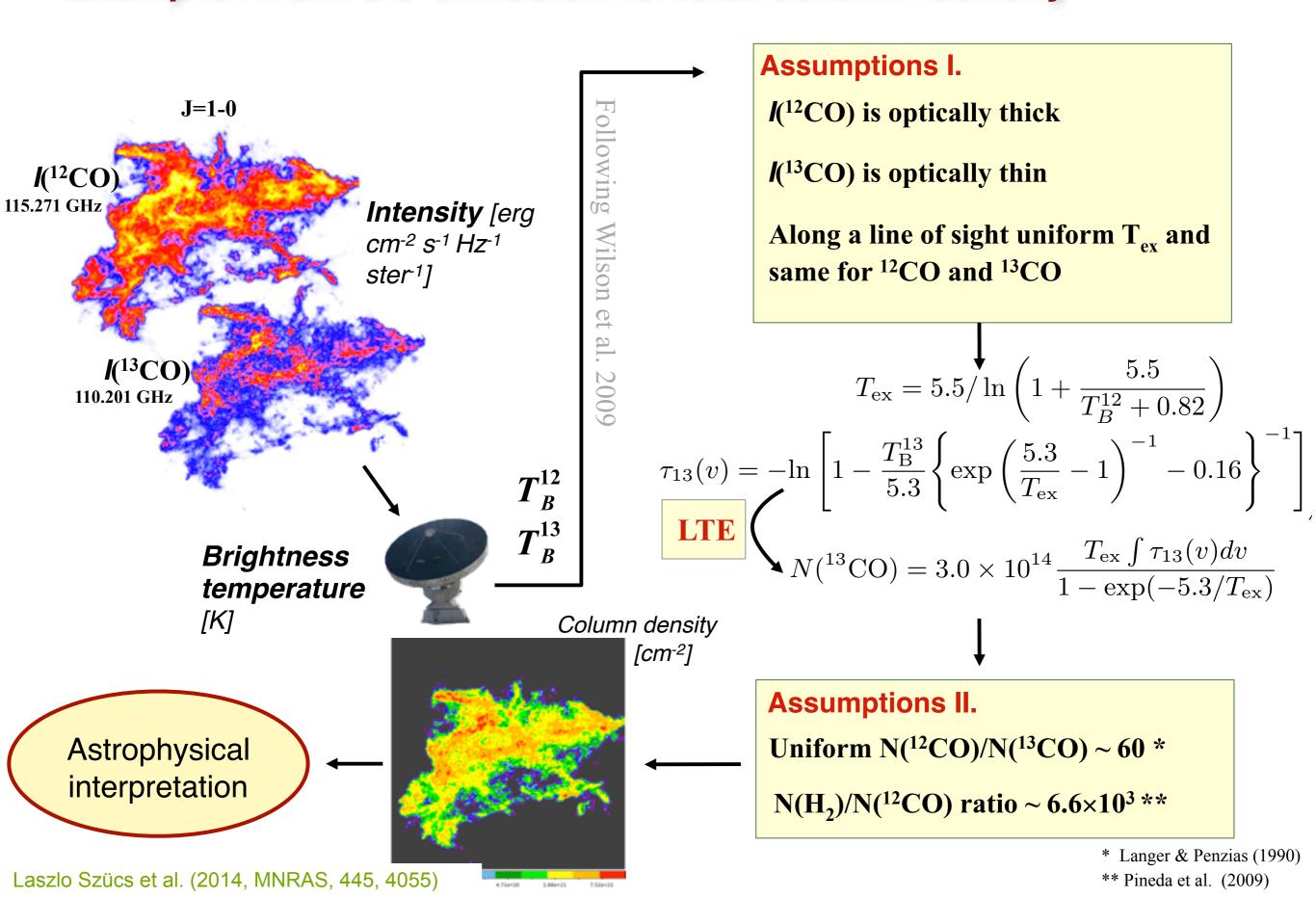


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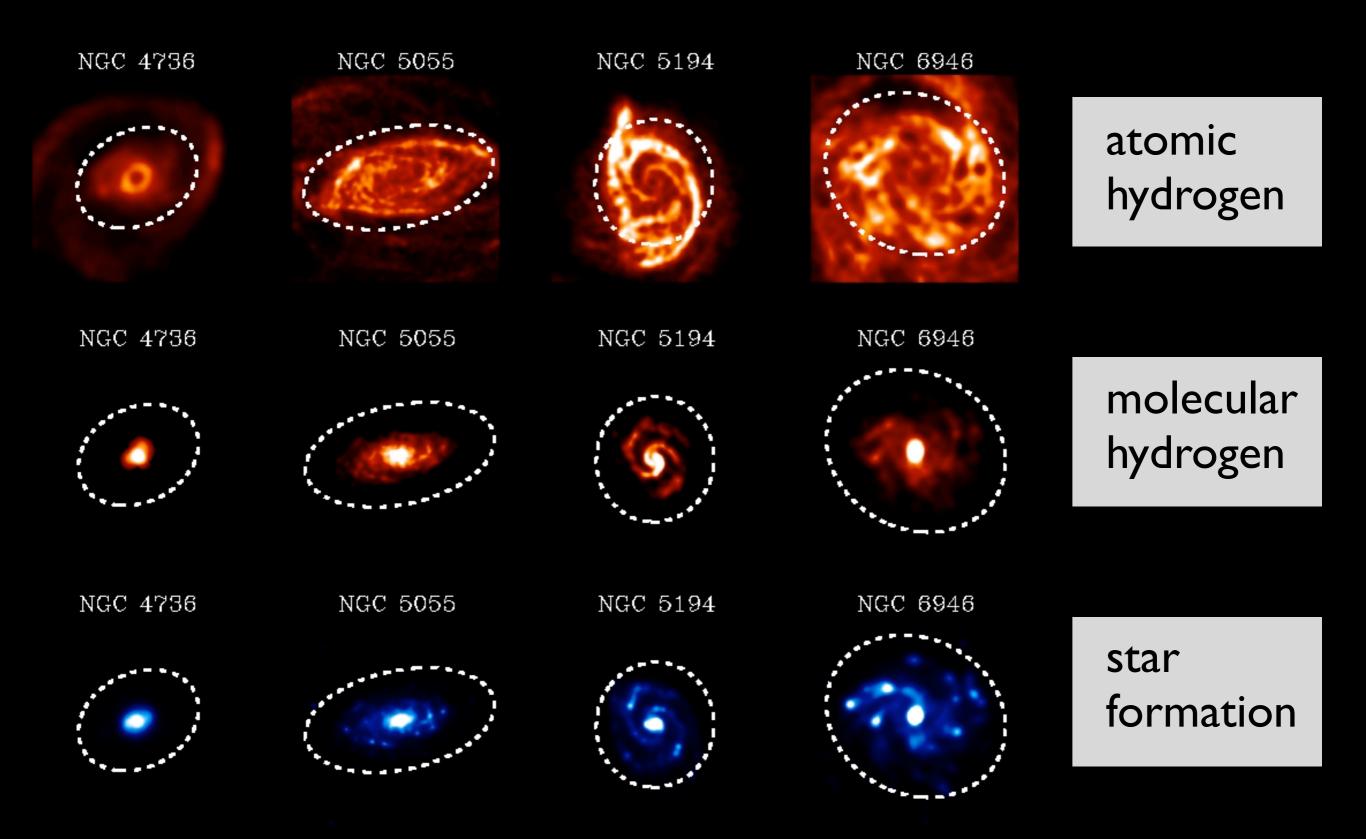


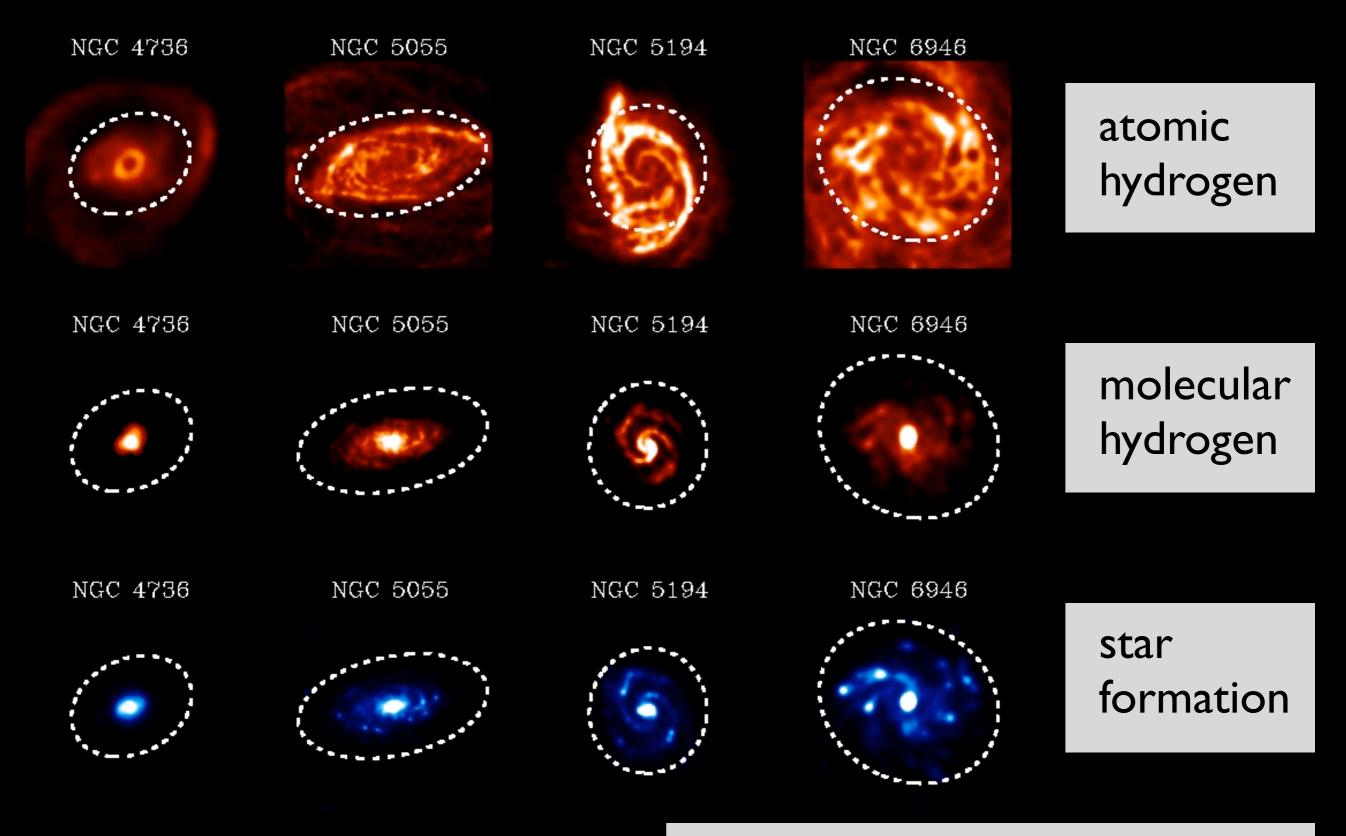


Example: from CO emission to total column density

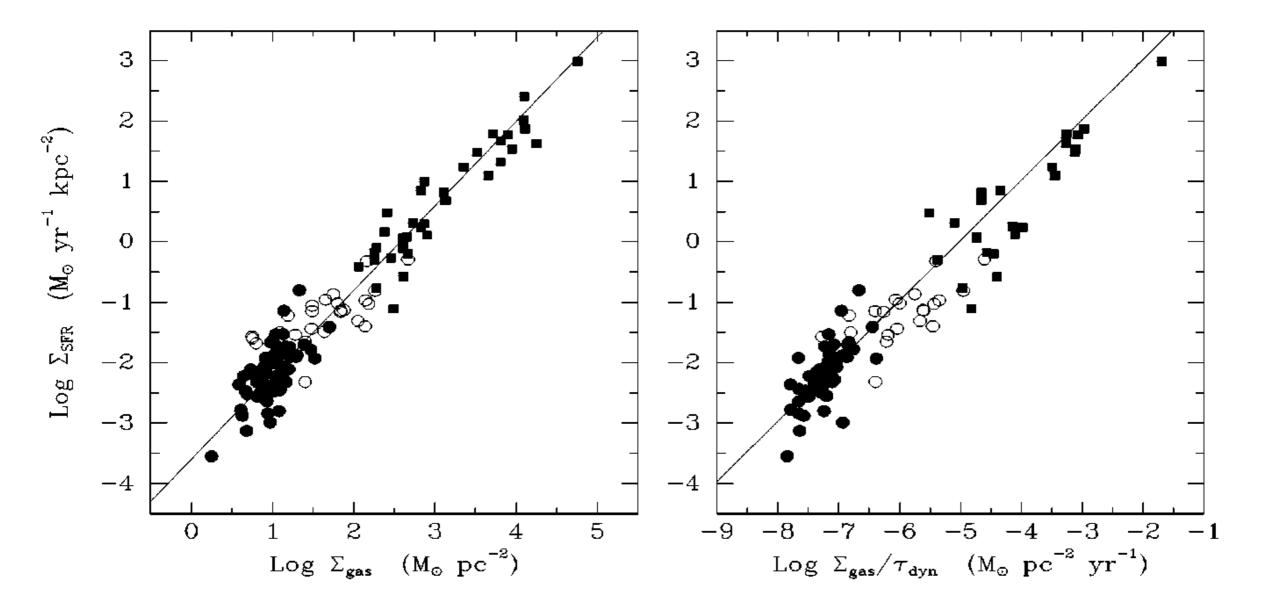








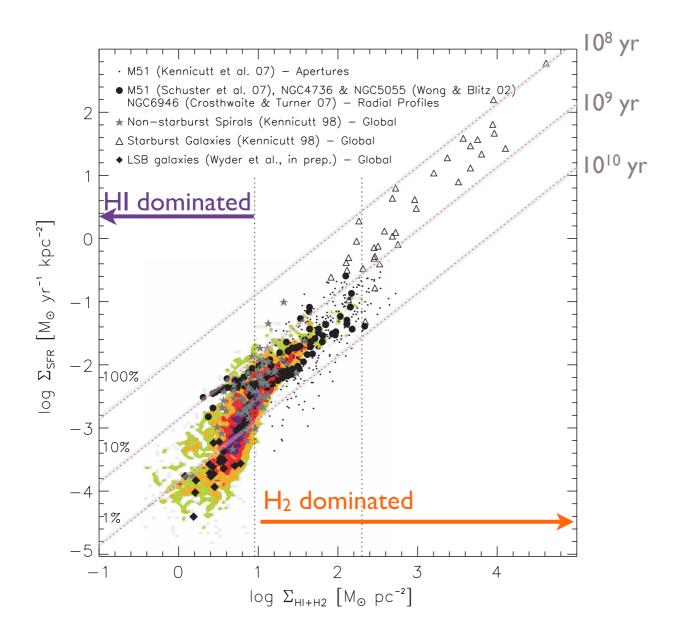
- HI gas more extended
- H2 and SF well correlated



Kennicutt (1998, ARAA, 36, 189)

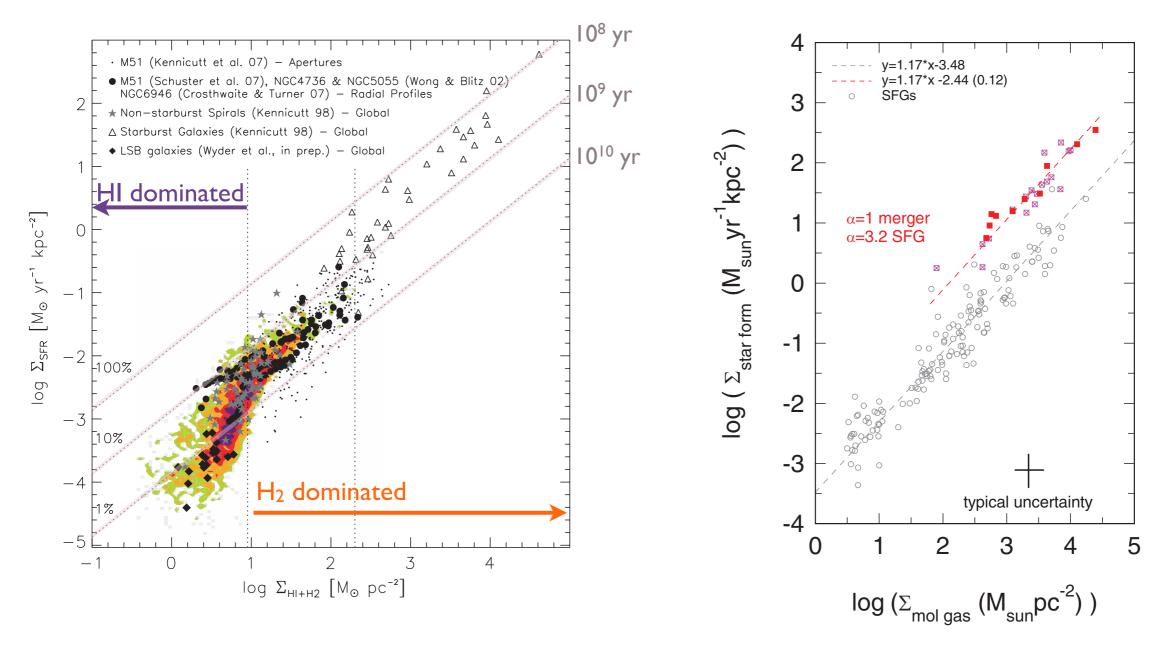
• when considering galaxies as a whole, there seems to be a super-linear relation between total gas (H_2+H) and the star formation rate (SFR) with slope ~ 1.4 :

$$\Sigma_{SFR} = (2.5 \pm 0.7) \times 10^{-4} \left(\frac{\Sigma_{gas}}{1 \ M_{\odot} \ \text{pc}^{-2}} \right)^{1.4 \pm 0.15} M_{\odot} \ \text{year}^{-1} \, \text{kpc}^{-2}$$



Bigiel et al. (2008, AJ, 136, 2846)

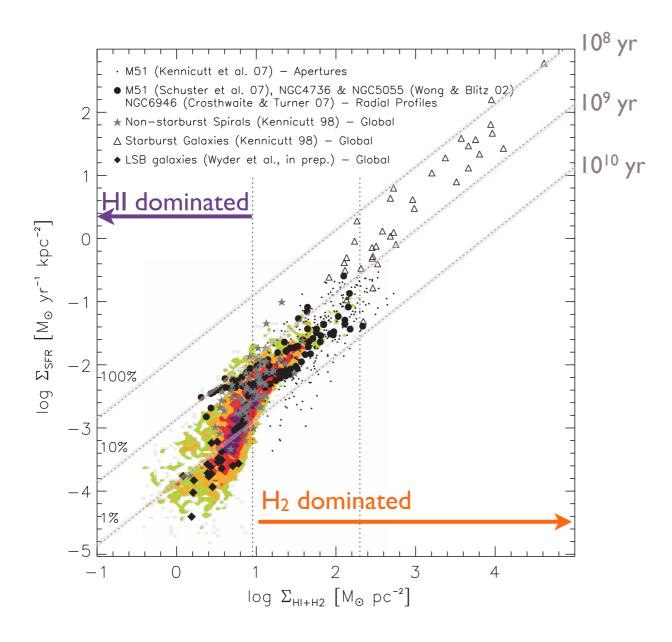
- for "resolved" galaxies on scales of 0.5-1 kpc, there seems to be a linear relation between H_2 and SFR
- implying a roughly constant depletion time of a few x 109 yr



Bigiel et al. (2008, AJ, 136, 2846)

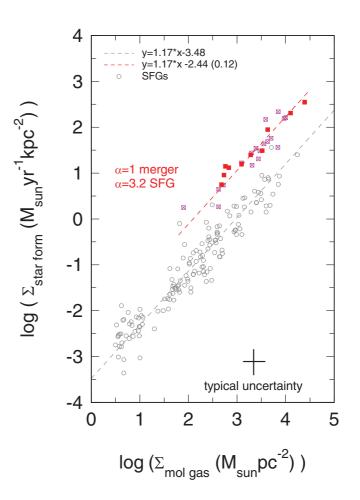
Genzel et al. (2010, MNRAS, AJ, 407, 2091)

- for "resolved" galaxies on scales of 0.5-1 kpc, there seems to be a linear relation between H_2 and SFR
- implying a roughly constant depletion time of a few $\times 10^9$ yr
- but with different normalization for starburst galaxies compared to normal ones

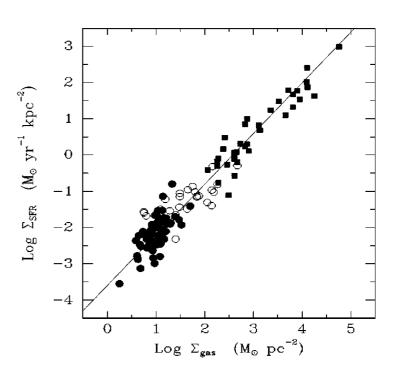


Bigiel et al. (2008, AJ, 136, 2846)

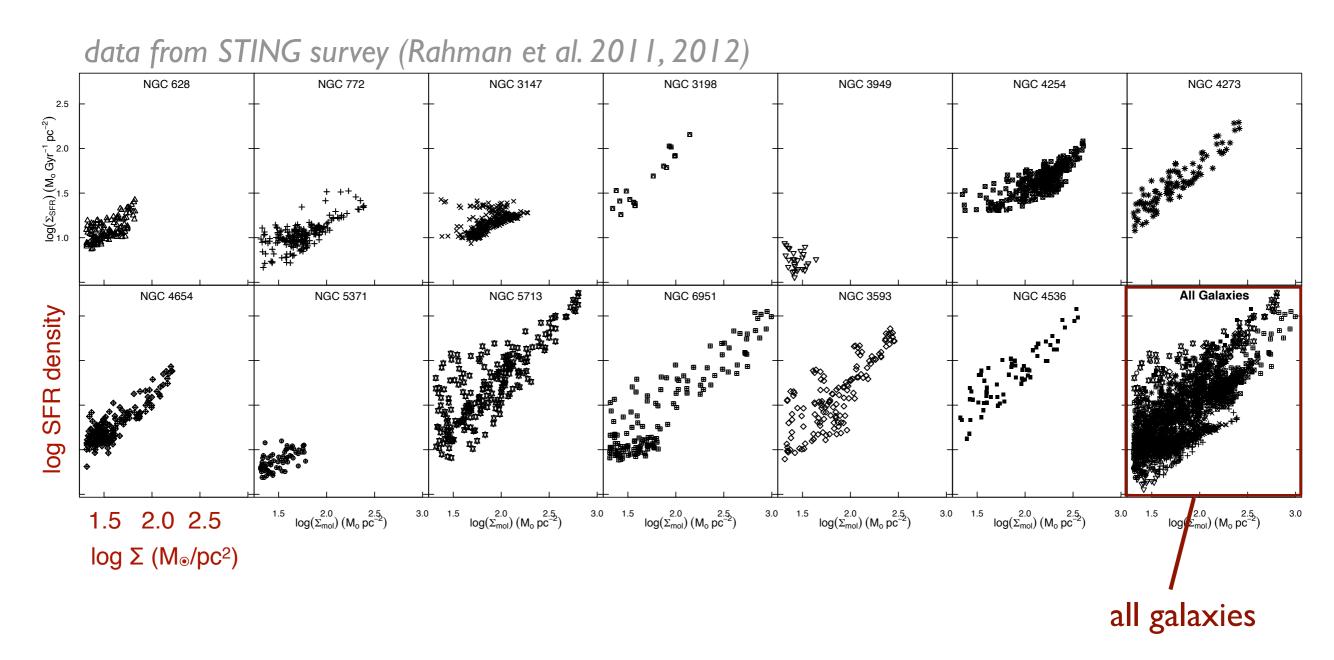
true physical behavior may be (much) more complicated than simple models assume!



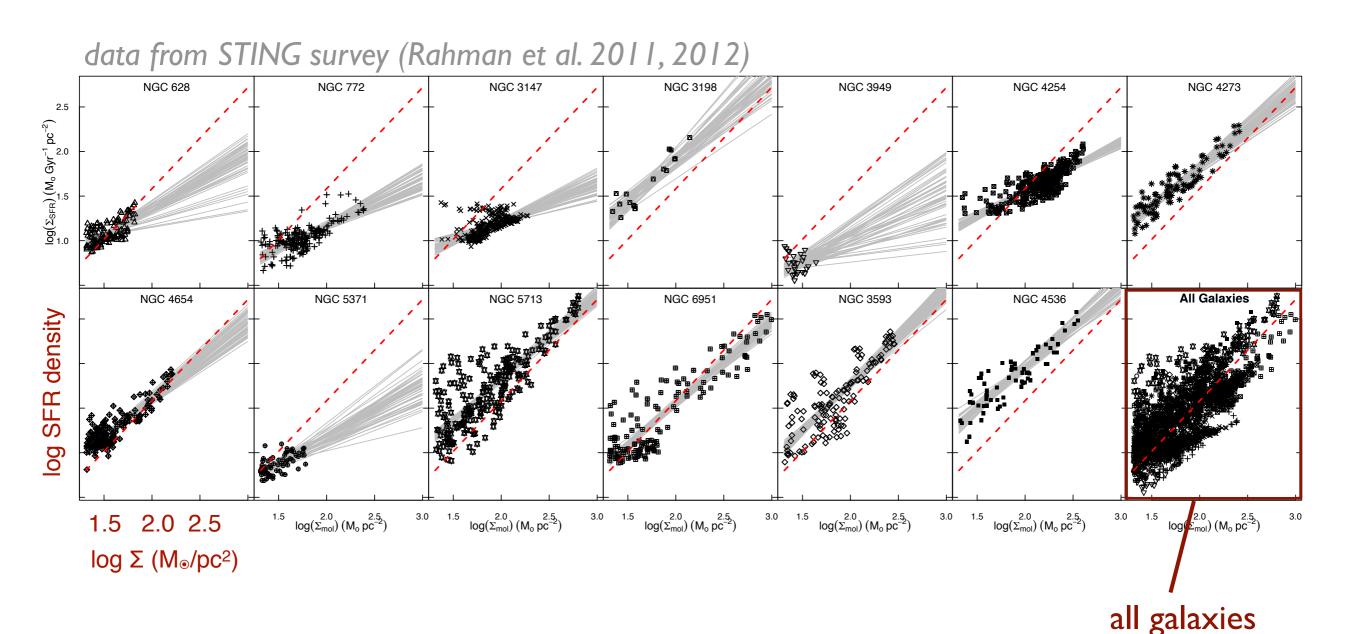
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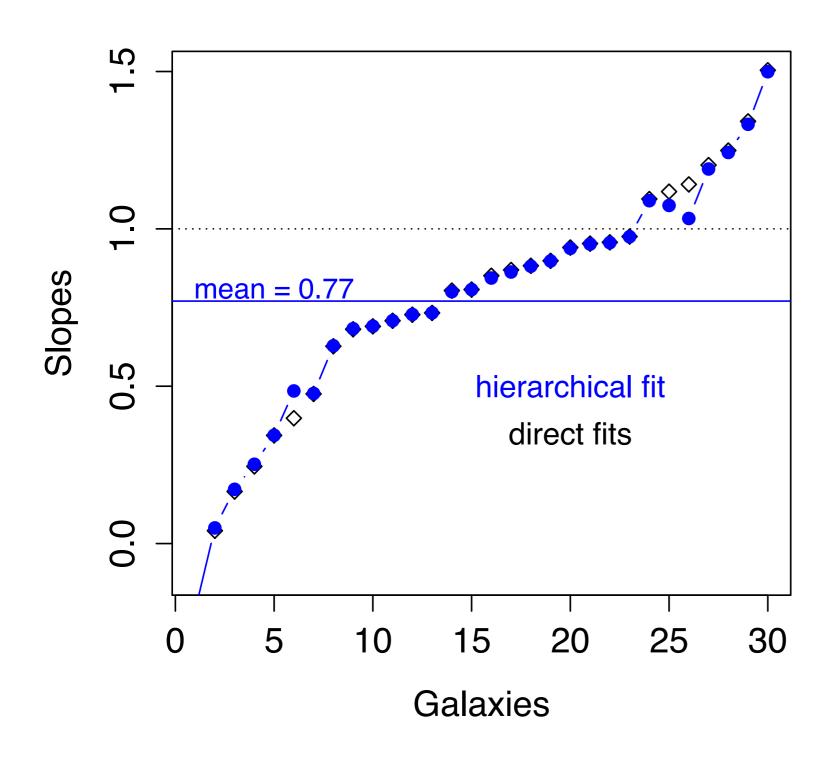
Kennicutt (1998, ARAA, 36, 189)



• QUIZ: do you see a universal Σ_{H2} - Σ_{SFR} relation?

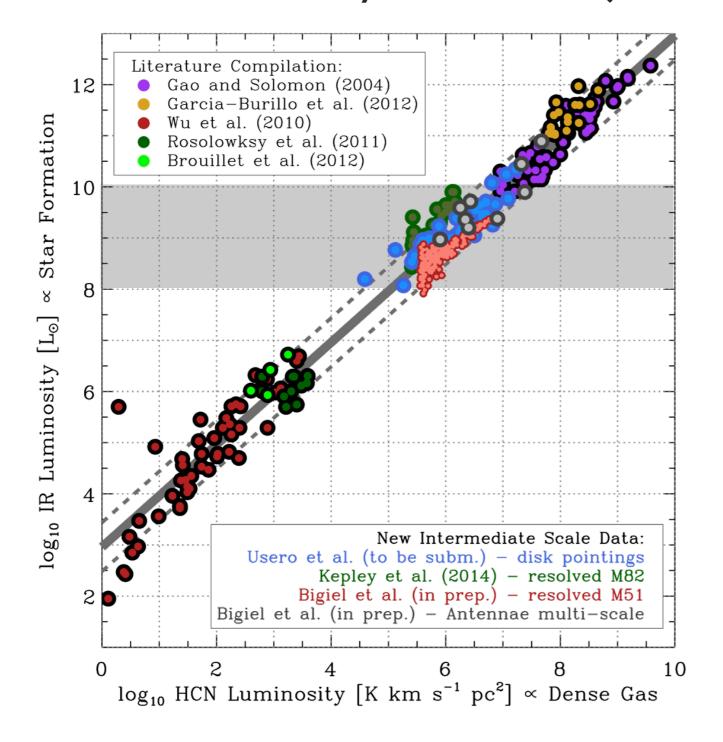


- QUIZ: do we really see a universal Σ_{H2} Σ_{SFR} relation?
- ANSWER: large galaxy-to-galaxy variations
 - relation is often sublinear



- analysis of THINGS/ HERACLES data
- many galaxies show sublinear KS-type relation

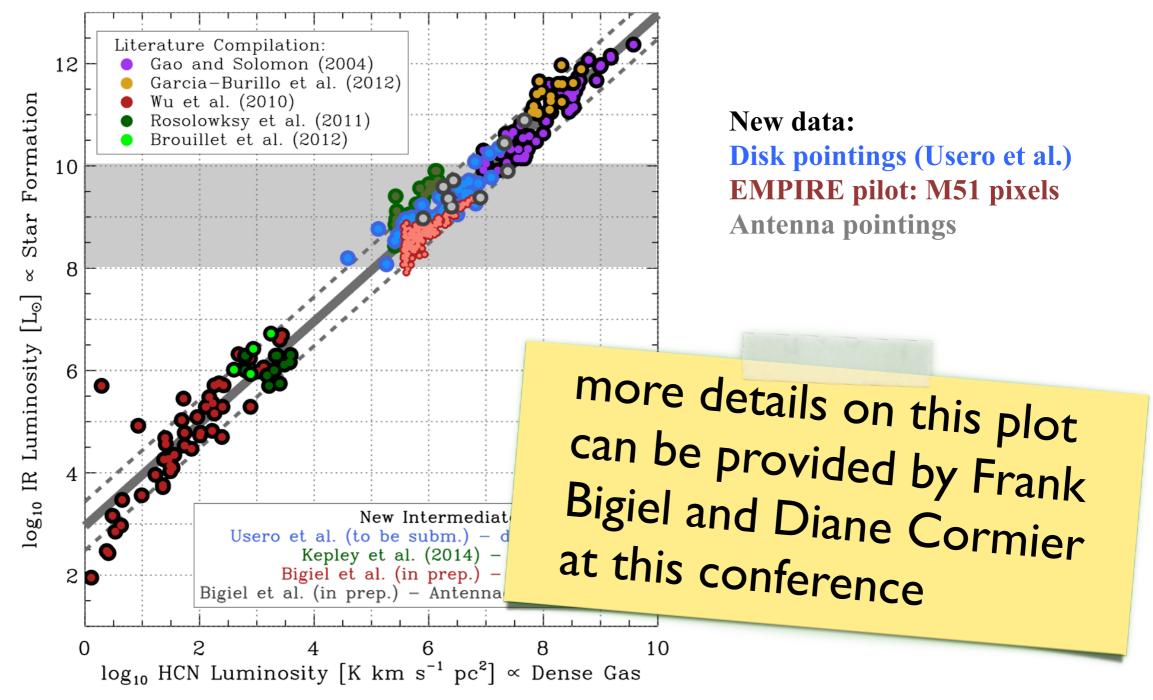
- HOWEVER: there seems to be a relation between SFR tracers and dense gas tracers that extends over many orders of magnitude!!
- this includes many different objects

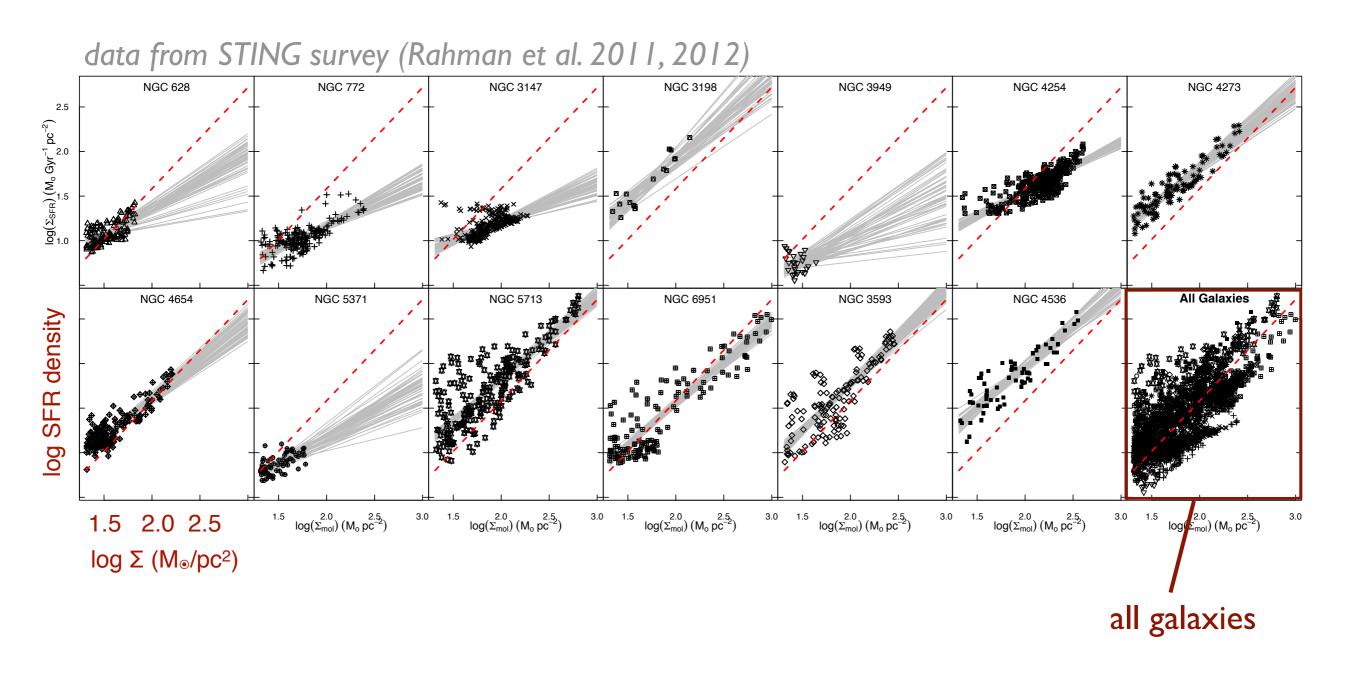


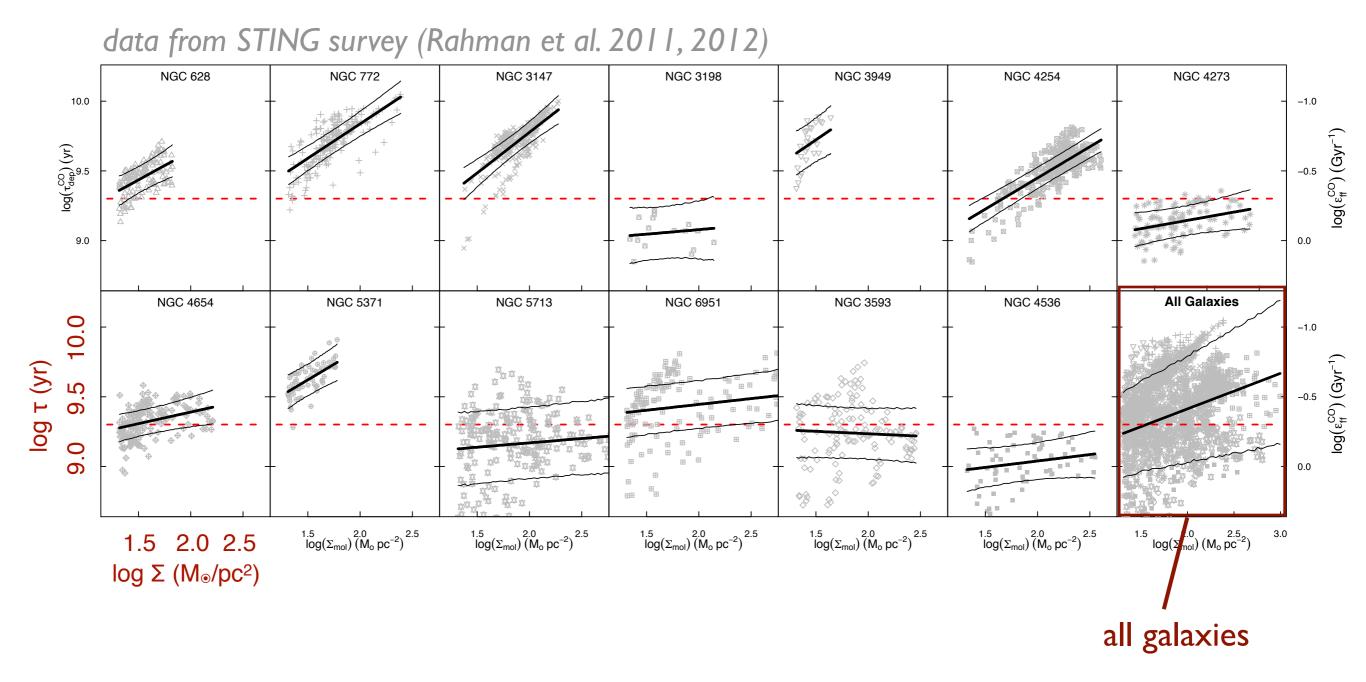
New data:

Disk pointings (Usero et al.) EMPIRE pilot: M51 pixels Antenna pointings

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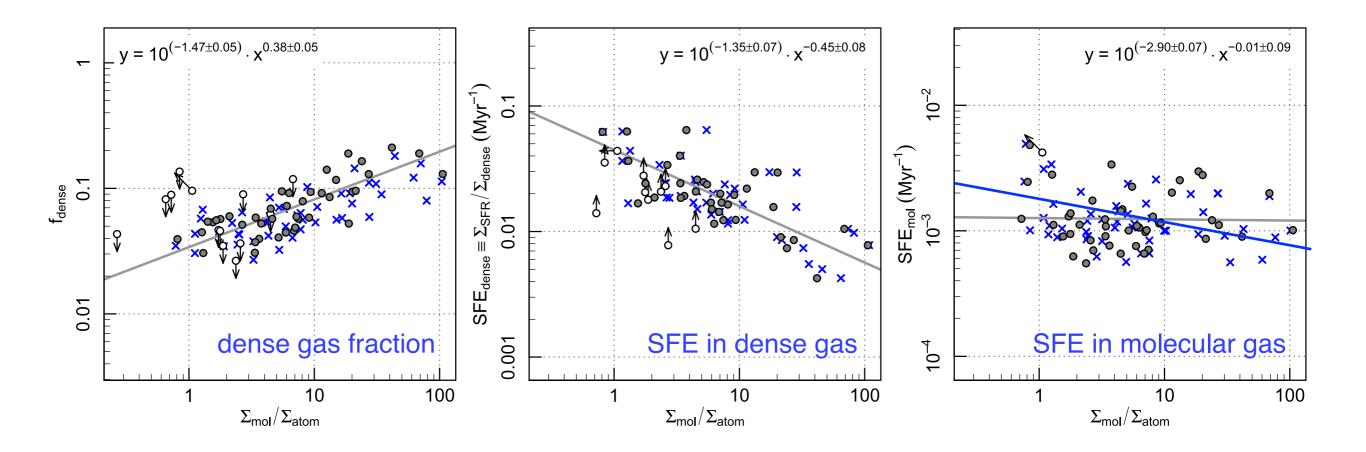






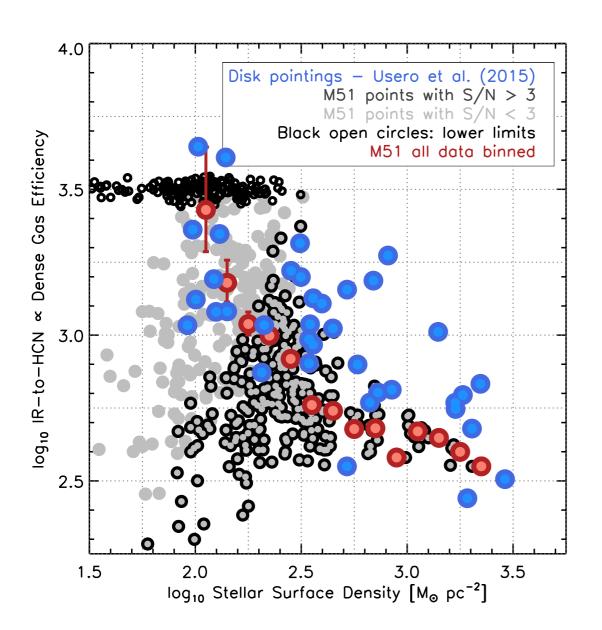
Hierarchical Bayesian model for STING galaxies indicate varying depleting times. Depletion time increases with increasing density. Why ??

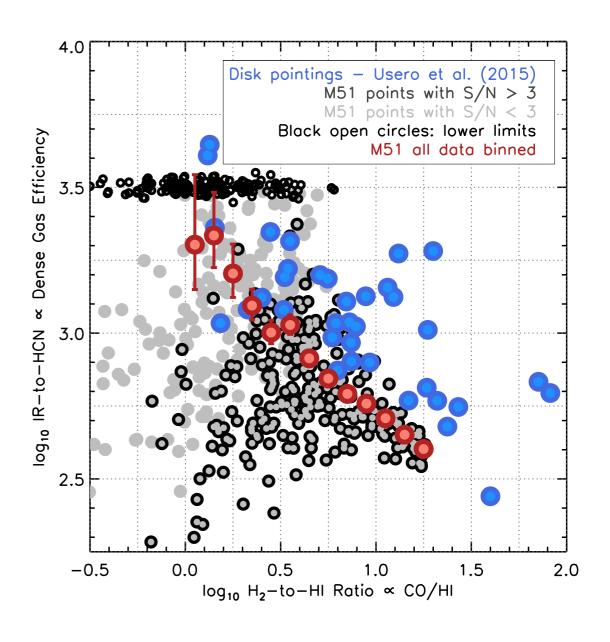
- EMPIRE Survey (PI Frank Bigiel):
- IR-to-HCN ratio varies systematically as function of local disk structure (here stellar surface density)
- dense gas is less good in forming stars in overall dense regions (longer depletion time)



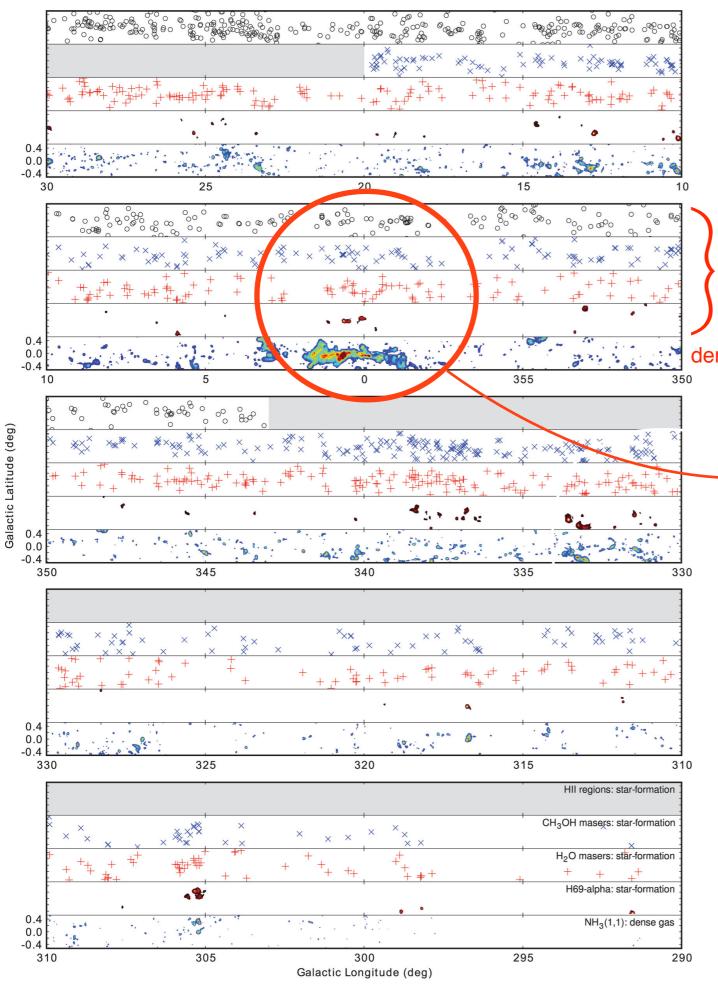
different galaxies in survey

- EMPIRE Survey (PI Frank Bigiel):
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resolved data in M5 I

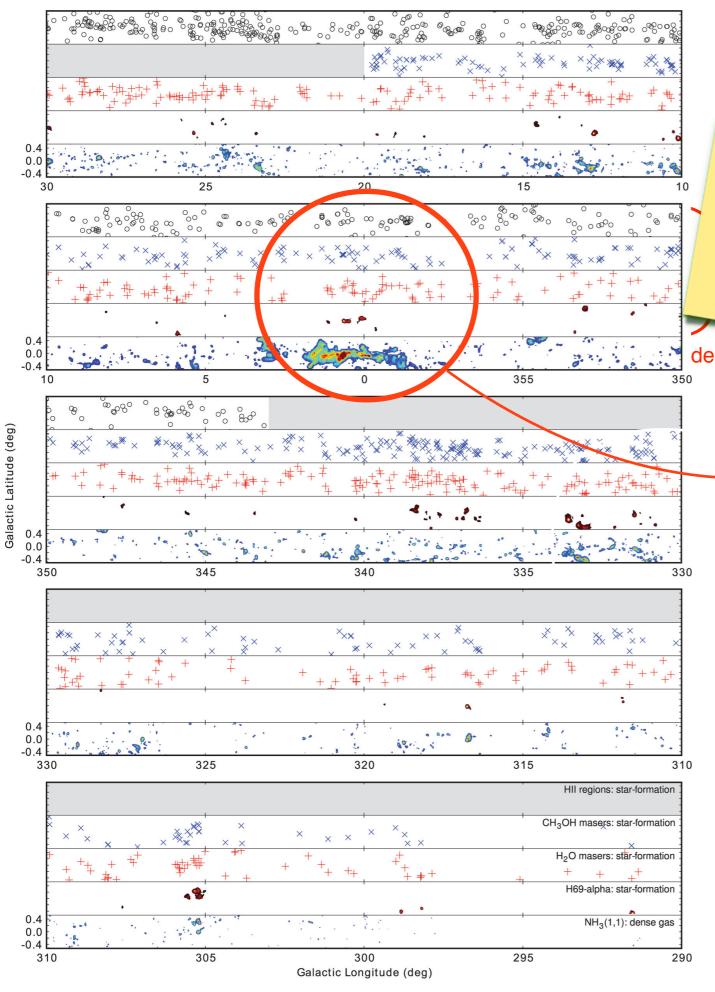


different SF tracers

dense gas tracer: NH₃(1,1)

Galactic Center

- similar holds for Galactic Center:
- dense gas in Central Molecular Zone (CMZ) seems relative inefficient in forming stars
 - for numerical modeling see
 Bertram (2015, MNRAS, 451, 3679),
 Bertram (2016, MNRAS, 455, 3763)



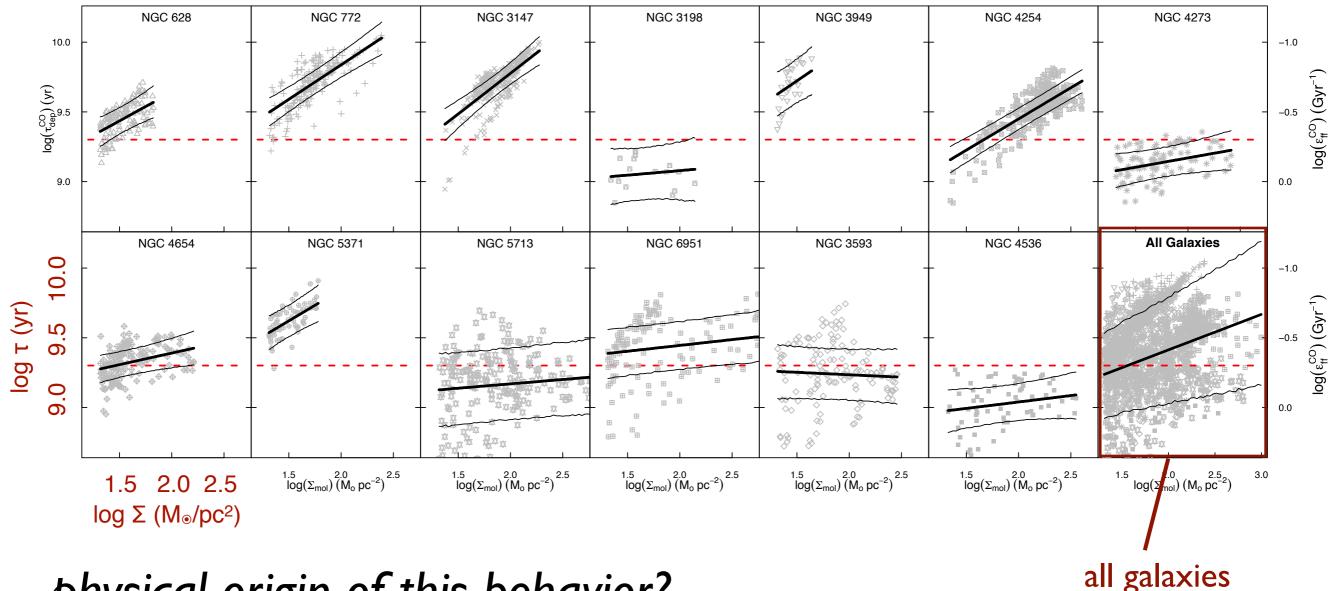
more Galactic Center in review talk by Mark Morris

dense gas tracer: NH₃(1

Galactic Center

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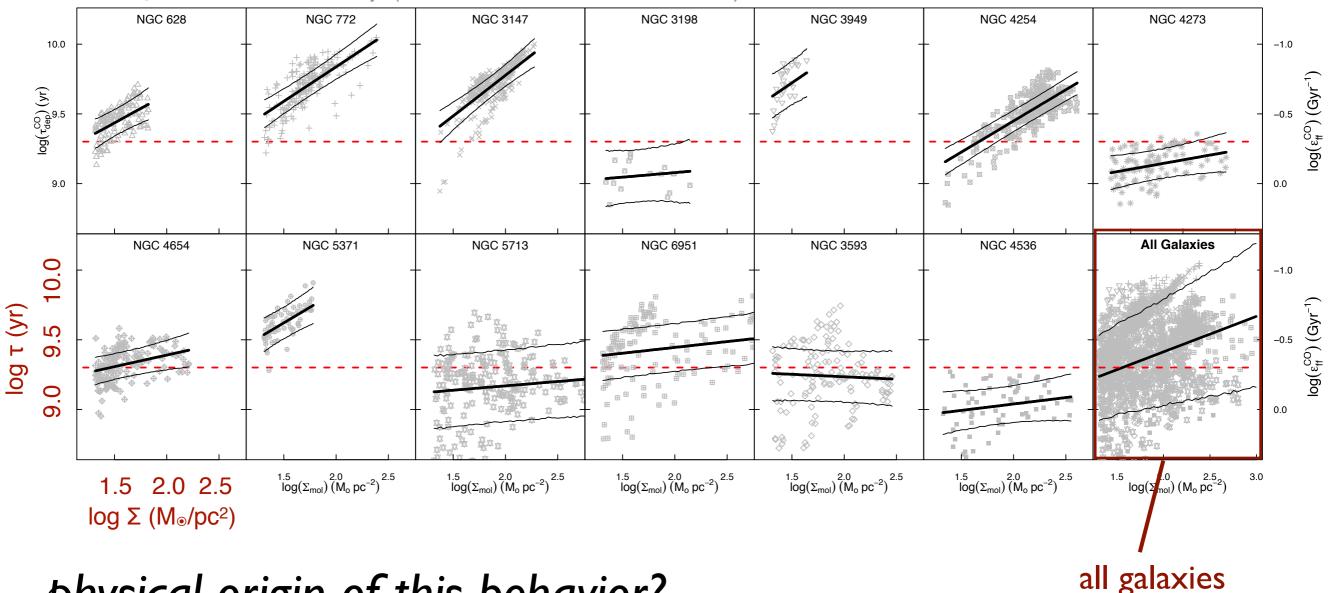




physical origin of this behavior?

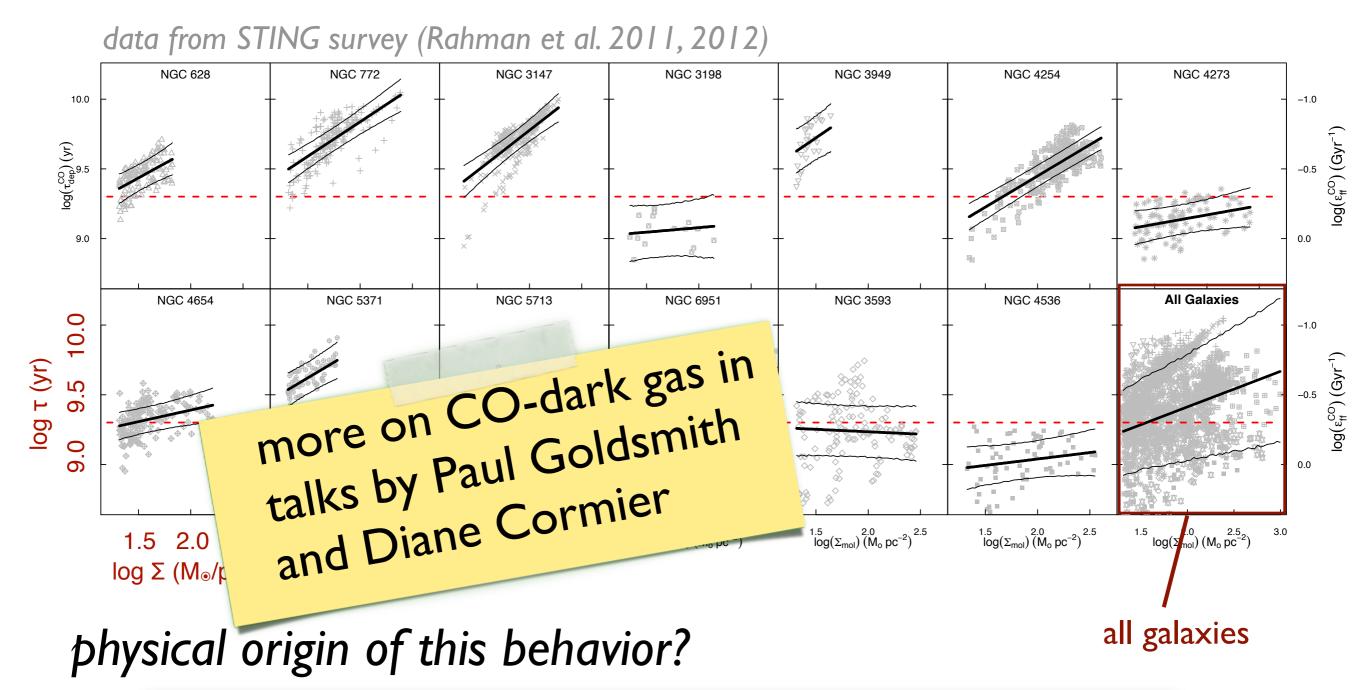
- maybe strong shear in dense arms (example M51, Meidt et al. 2013)...
- maybe non-star forming H_2 gas becomes traced by CO at high column densities (recall H_2 needs $A_v \sim I$, CO needs $A_v \sim 2$,)...





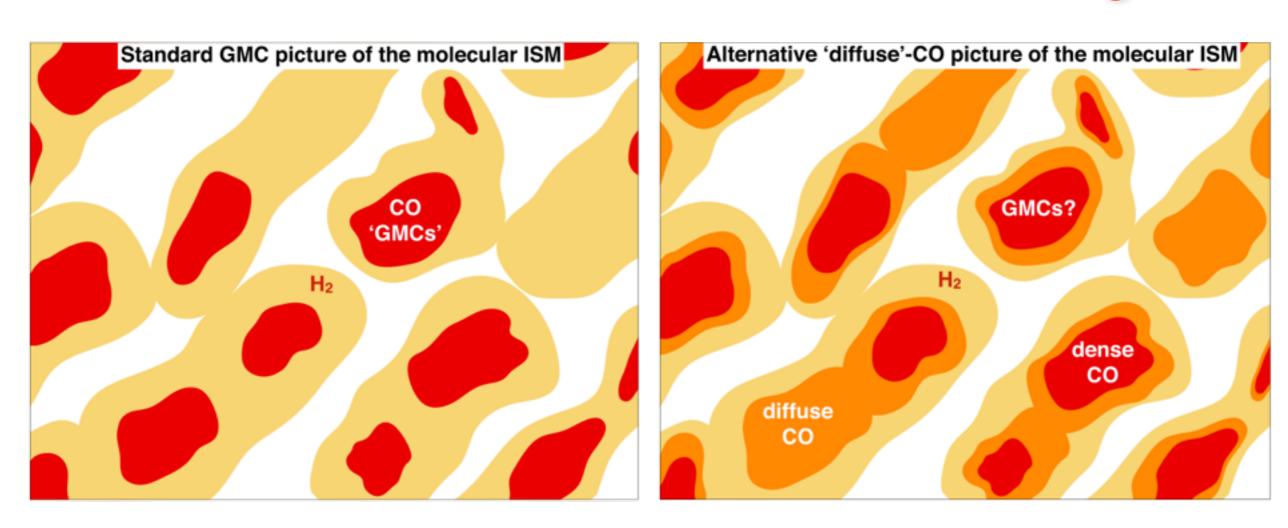
physical origin of this behavior?

SEARCH FOR CO-dark H₂ GAS here SOFIA can provide major input



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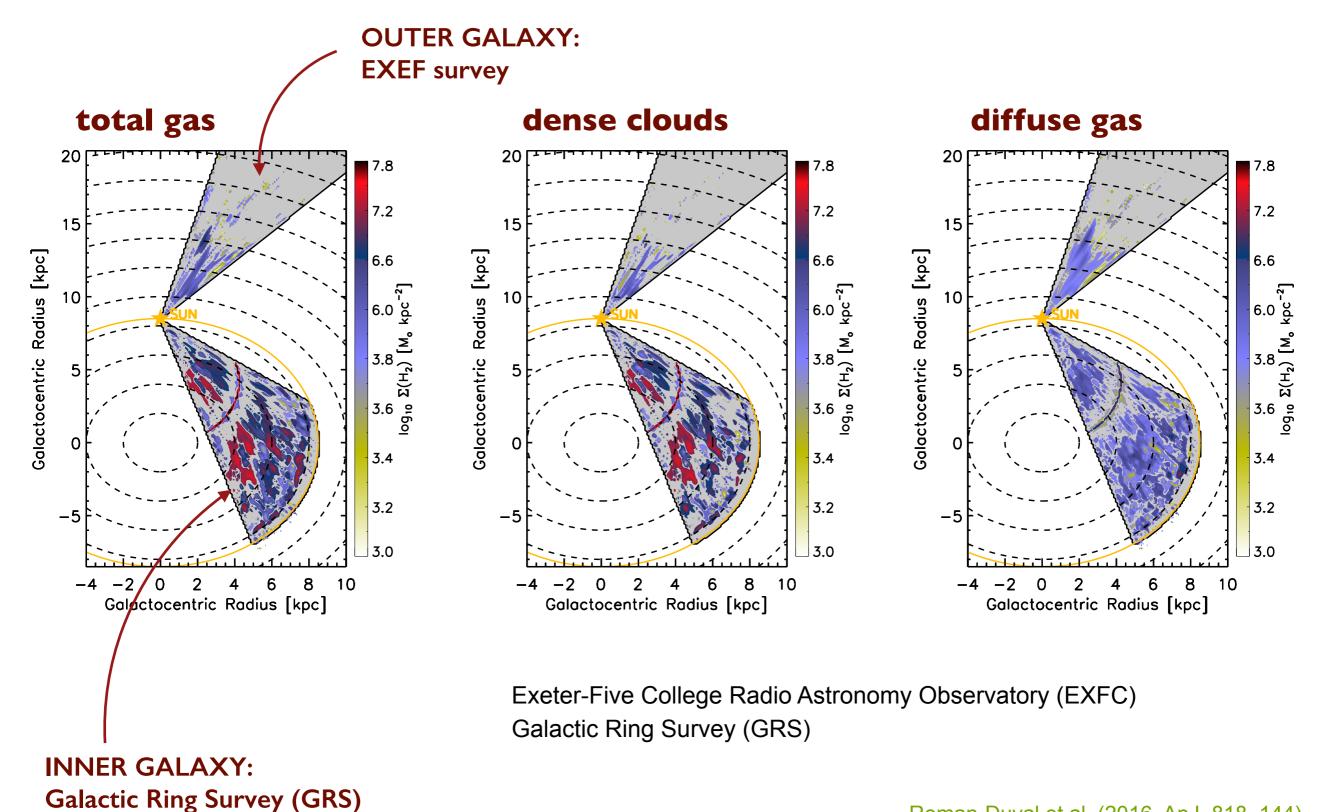
dense vs. diffuse CO-traced H2 gas

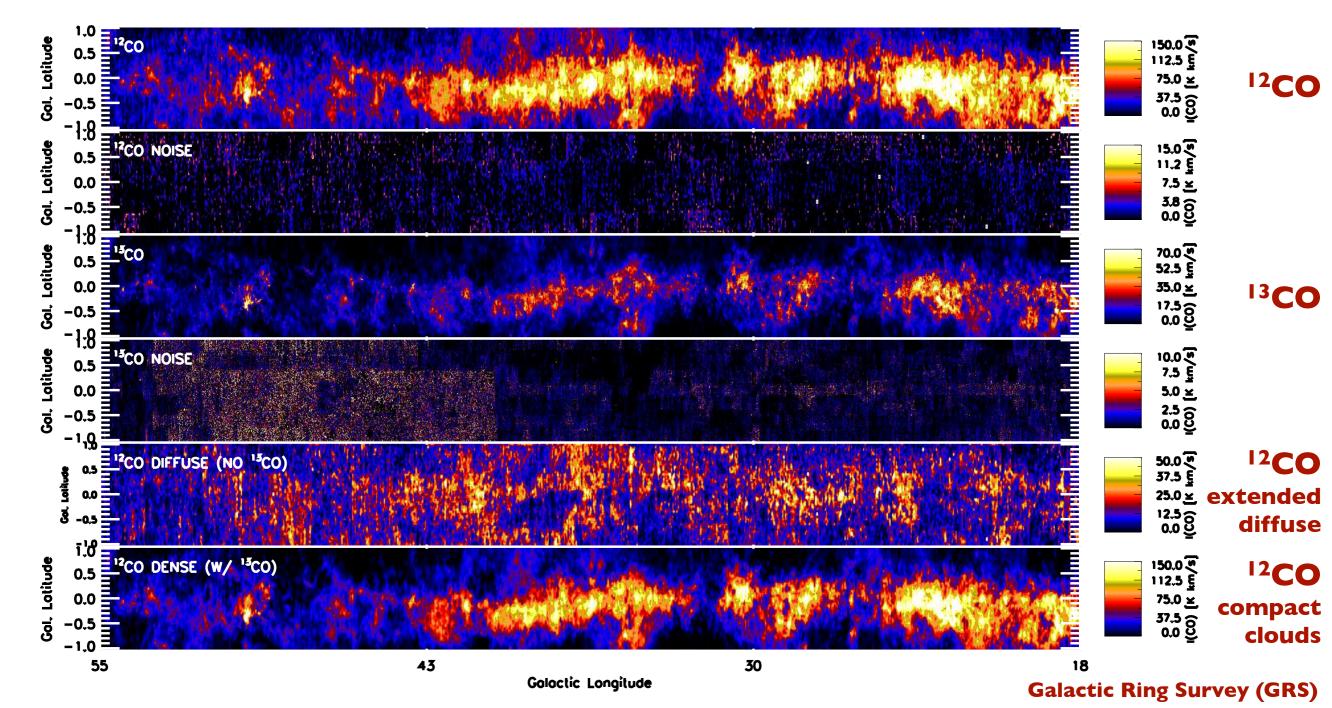


in addition:

 maybe a large fraction of H₂ (even if traced by CO) may not be in dense clouds, but in a diffuse state!

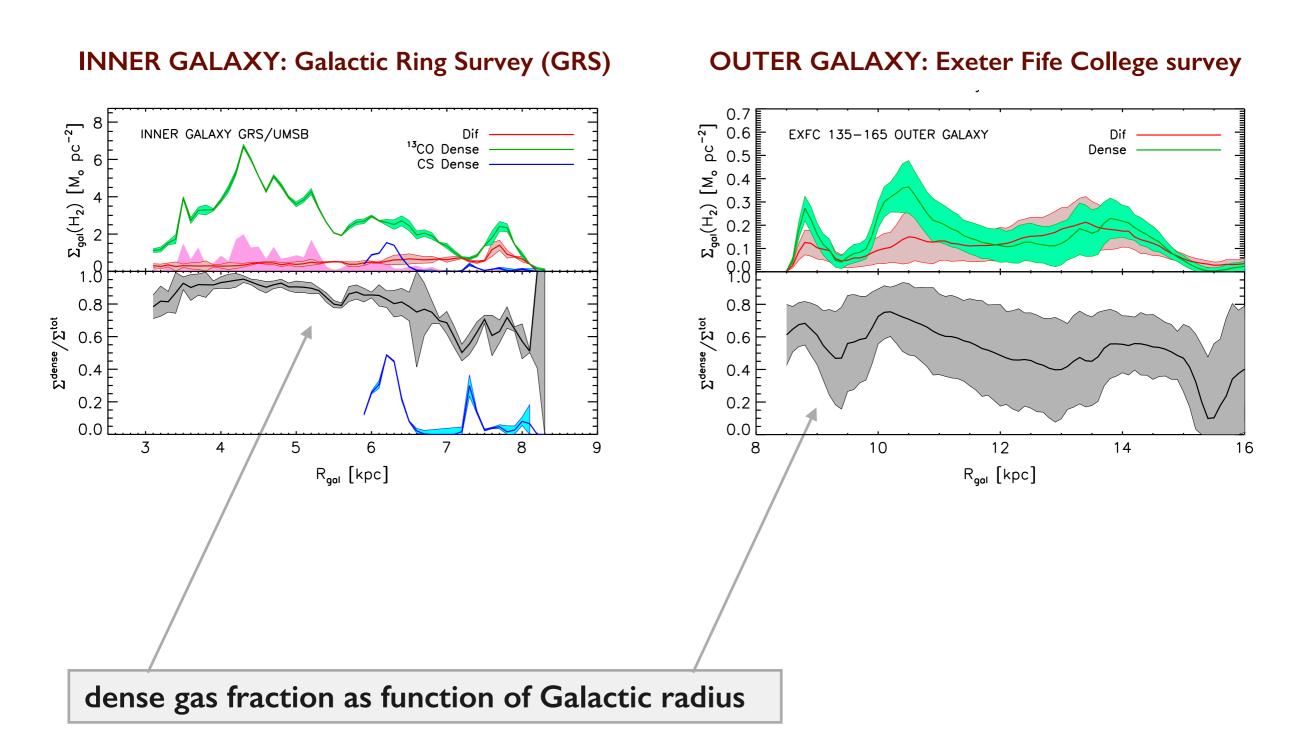
observational approach





observational approach:

comparison of ¹³CO (tracing mostly dense clouds) and ¹²CO tracing all the gas (including the more diffuse component)



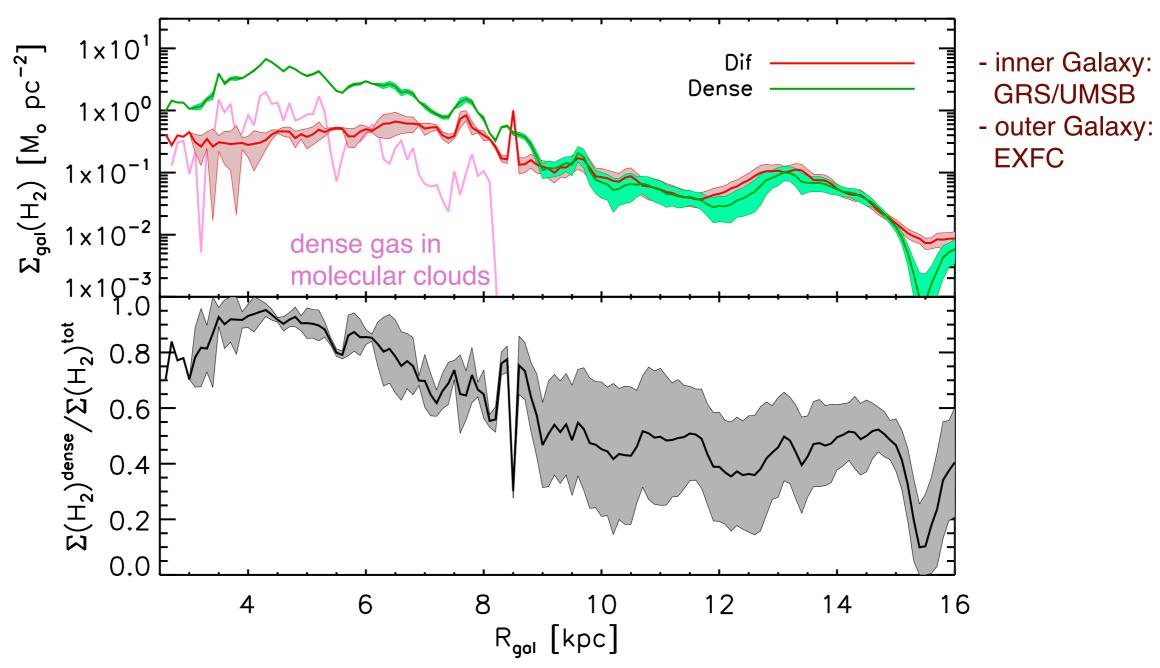
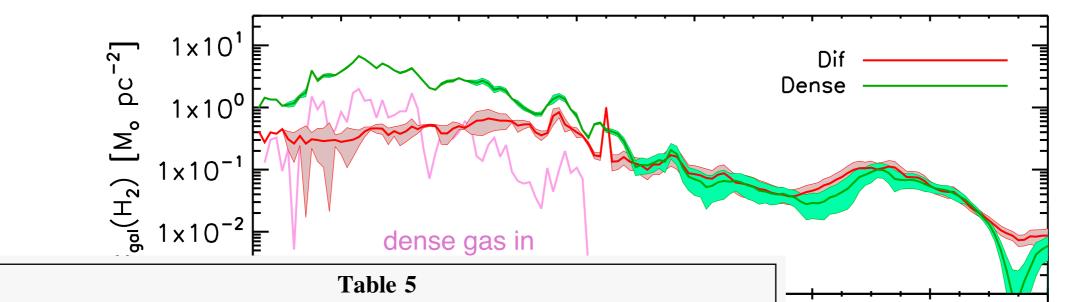


Figure 13. Average Galactic H_2 surface densities of the diffuse (red, detected in ^{12}CO , undetected in ^{13}CO) and dense (green, detected in ^{12}CO and ^{13}CO) components as a function of Galactocentric radius (in bins of width 0.1 kpc), in logarithmic scale, combining all data sets. In the inner Galaxy, the pink line indicates the surface density of H_2 in molecular clouds identified in Roman-Duval et al. (2010).



inner Galaxy:GRS/UMSB

outer Galaxy:EXFC

Table 5

Total Luminosity and Molecular Mass in the Milky Way in the Diffuse and Dense Components Traced by ¹²CO

		Inner	Outer	Total
L(12CO)	Diffuse	2.0×10^{1}	4.0	2.4×10^{1}
	Dense	1.1×10^{2}	3.8	1.1×10^{2}
	Very dense	4.8	•••	4.8
	Total	1.3×10^{2}	7.7	1.4×10^{2}
$M(H_2)$	Diffuse	9.3×10^{7}	6.0×10^{7}	1.5×10^{8}
	Dense	4.6×10^{8}	3.9×10^{7}	4.9×10^{8}
	Very dense	2.9×10^7		2.9×10^{7}
	Total	5.5×10^8	9.9×10^7	6.5×10^{8}

fraction CO-traced H2 gas in Milky Way:

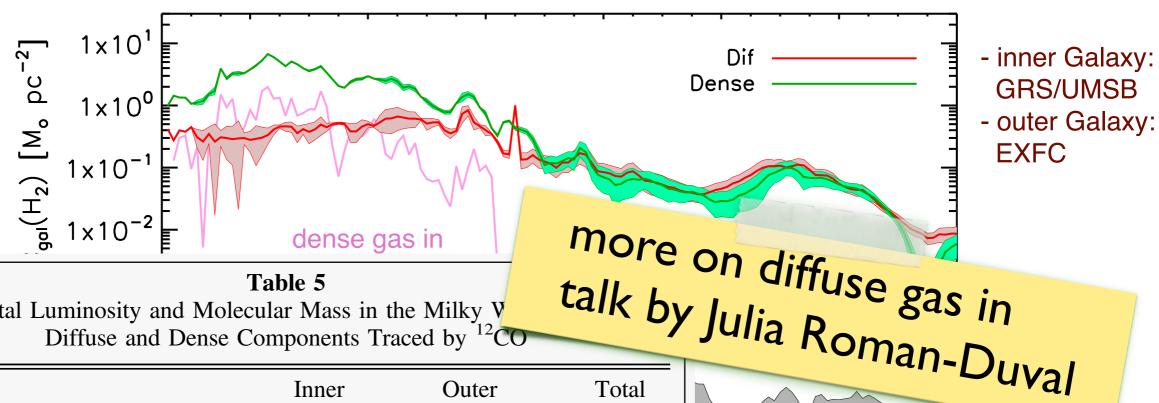
~1/4 diffuse

~3/4 dense

nd ~1/20 in known molecular clouds only !!!

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Inner Oute	er T

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fraction CO-traced H2 gas in Milky Way: ~1/4 diffuse

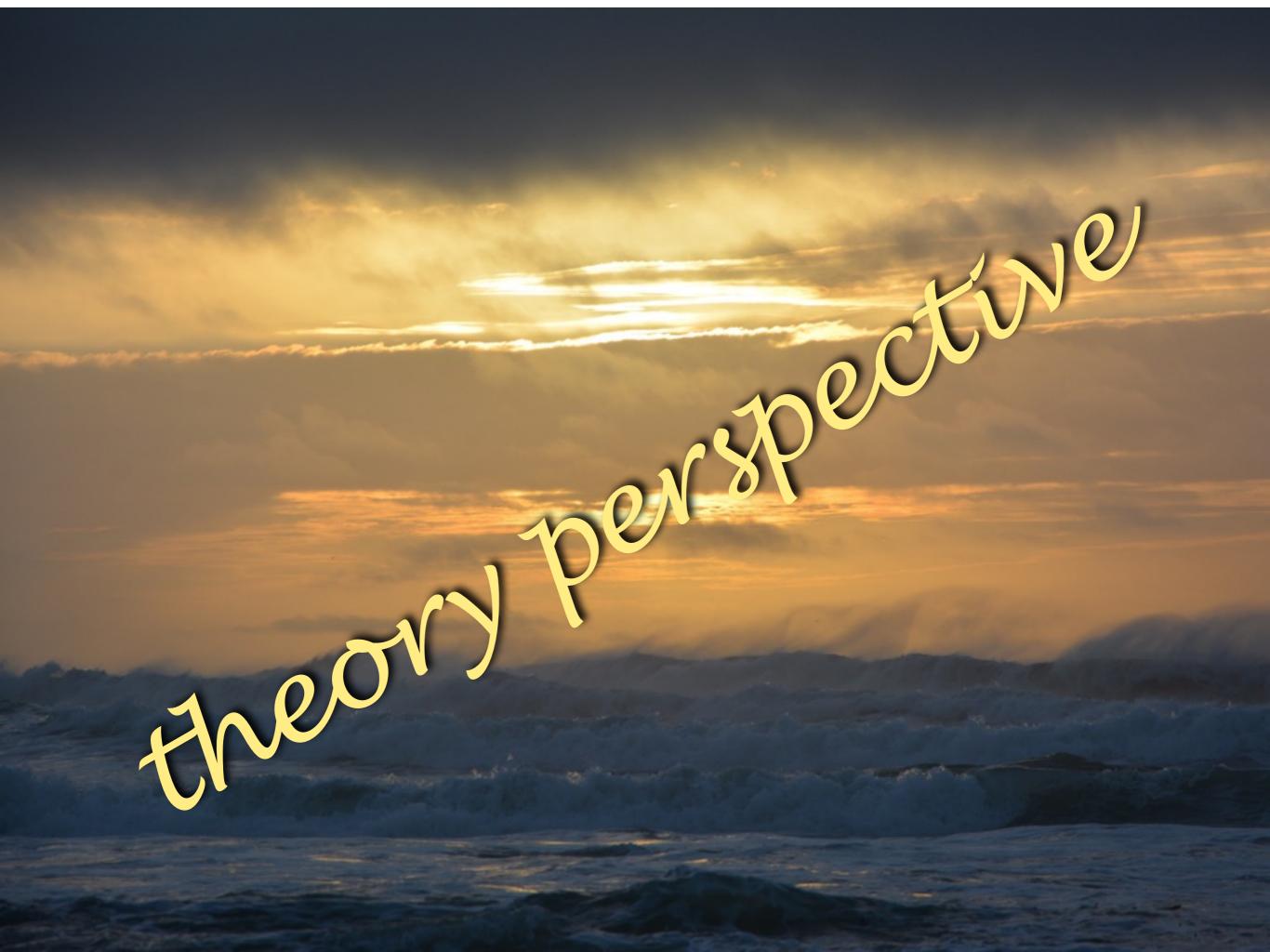
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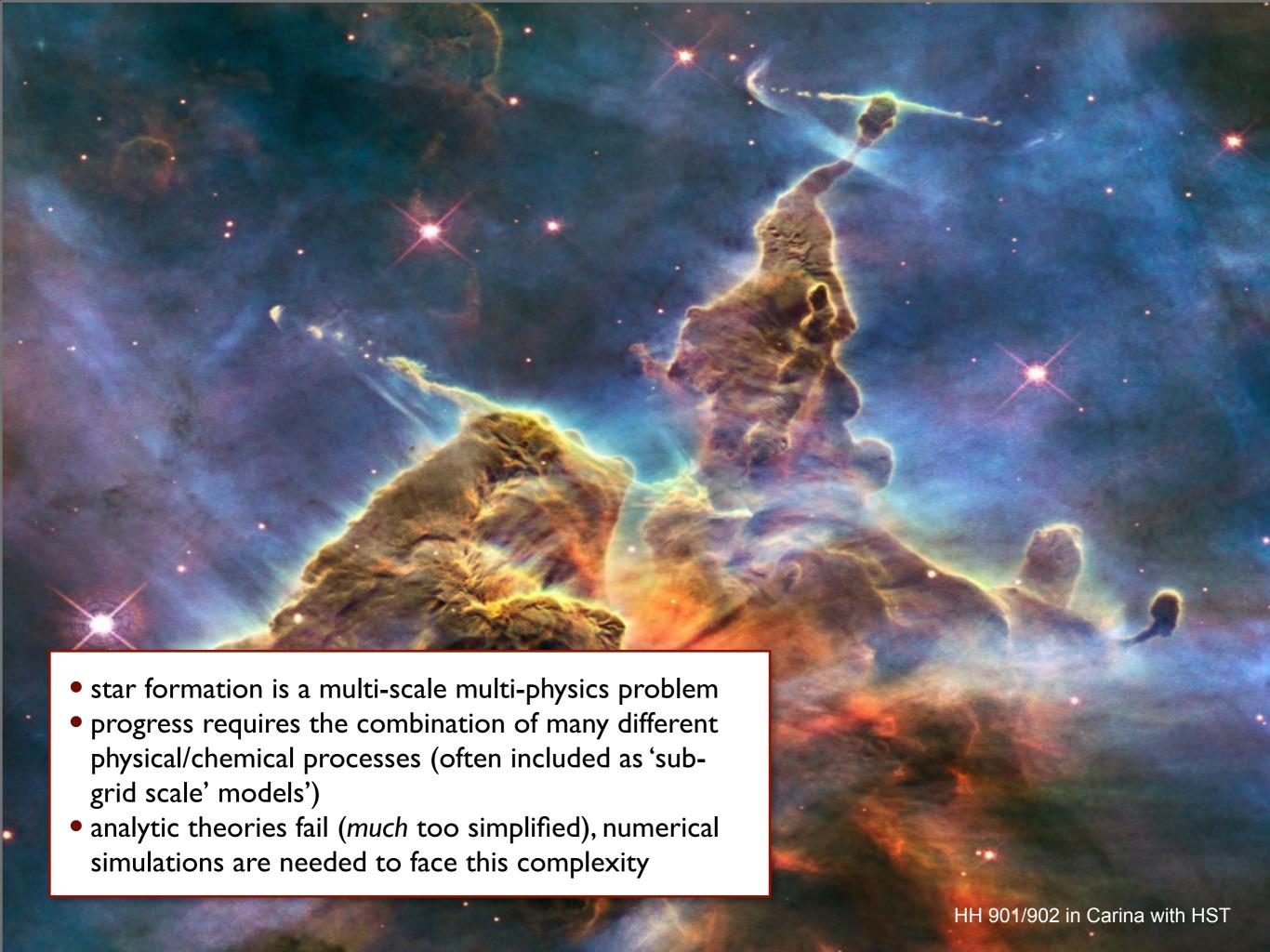
face

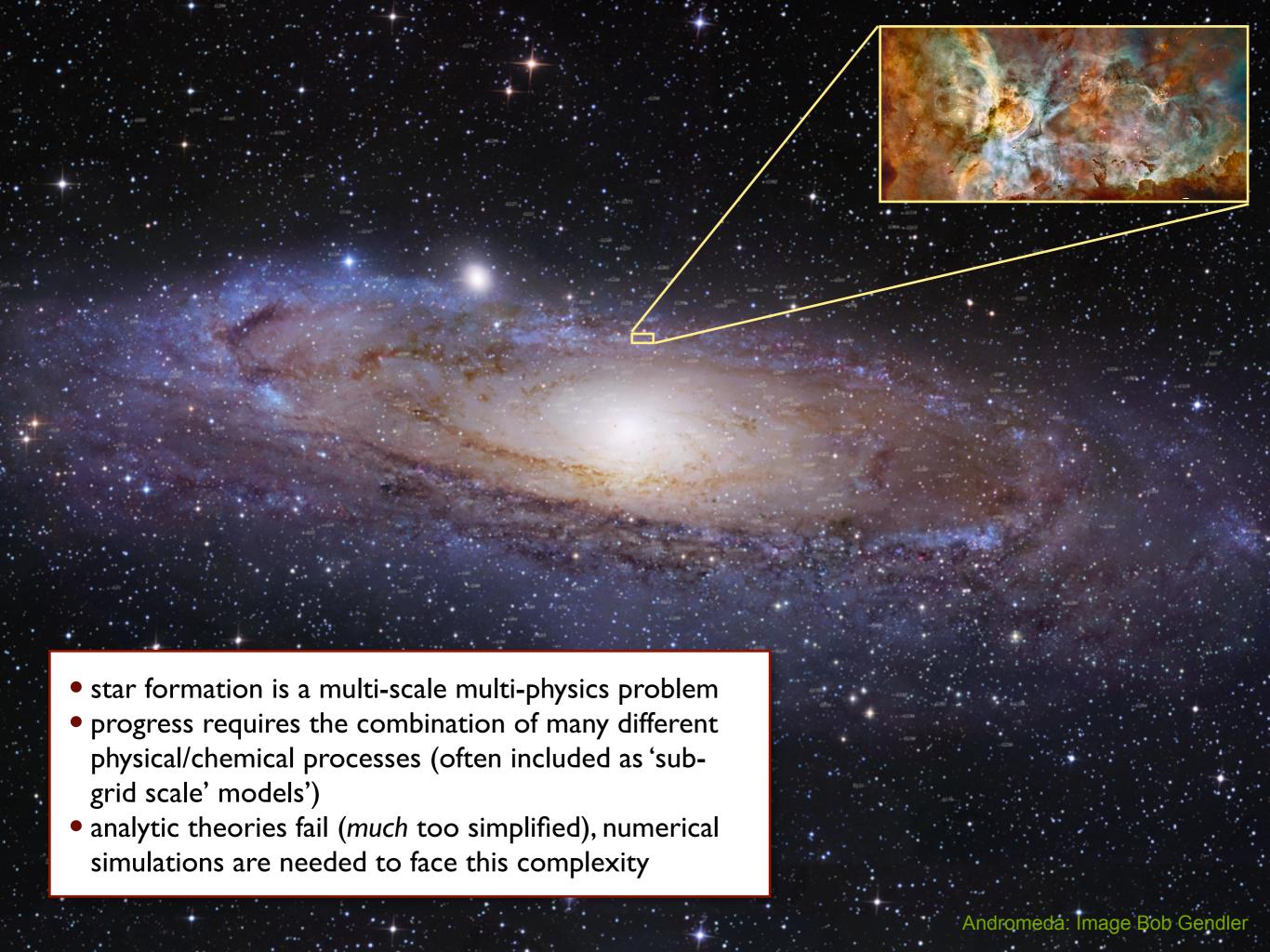


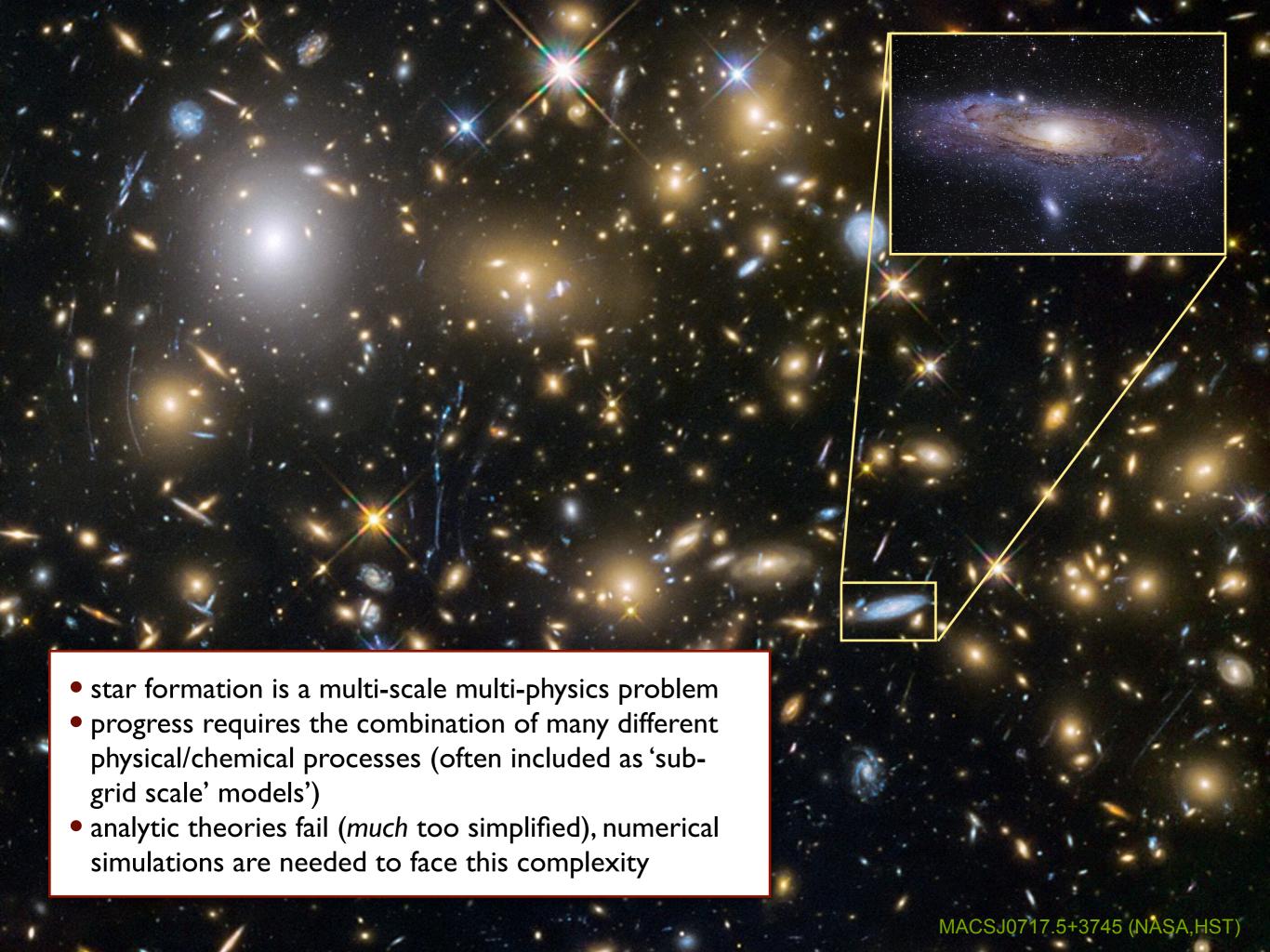












decrease in spatial scale / increase in density













- star formation is a multi-scale multi-physics problem
- progress requires the combination of many different physical/chemical processes (often included as 'subgrid scale' models')
- analytic theories fail (*much* too simplified), numerical simulations are needed to face this complexity

early theoretical models

- Jeans (1902): Interplay between self-gravity and thermal pressure
 - stability of homogeneous spherical density enhancements against gravitational collapse
 - dispersion relation:

$$\omega^2 = c_s^2 k^2 - 4\pi G \rho_0$$

- instability when

$$\omega^2 < 0$$

- minimal mass:

Sir James Jeans, 1877 - 1946

$$M_J = \frac{1}{6}\pi^{-5/2}G^{-3/2}\rho_0^{-1/2}c_s^3 \propto \rho_0^{-1/2}T^{+3/2}$$

first approach to turbulence

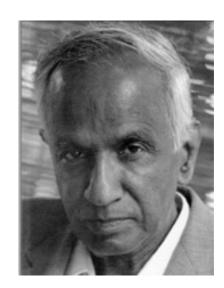
- von Weizsäcker (1943, 1951) and Chandrasekhar (1951): concept of MICROTURBULENCE
 - BASIC ASSUMPTION: separation of scales between dynamics and turbulence

$$\ell_{\text{turb}} \ll \ell_{\text{dyn}}$$

- then turbulent velocity dispersion contributes to effective soundspeed:

$$c_c^2 \mapsto c_c^2 + \sigma_{rms}^2$$

- → Larger effective Jeans masses → more stability
- BUT: (1) turbulence depends on k: $\sigma_{rms}^2(k)$
 - (2) supersonic turbulence $\rightarrow \sigma_{rms}^2(k) >> c_s^2$ usually





S. Chandrasekhar, 1910 - 1995

C.F. von Weiszäcker, 1912 - 2007

problems of early dynamical theory

- molecular clouds are highly Jeans-unstable, yet, they do NOT form stars at high rate and with high efficiency (Zuckerman & Evans 1974 conundrum) (the observed global SFE in molecular clouds is ~5%)
 - → something prevents large-scale collapse.
- all throughout the early 1990's, molecular clouds had been thought to be long-lived quasi-equilibrium entities.
- molecular clouds are magnetized

magnetic star formation

- Mestel & Spitzer (1956): Magnetic fields can prevent collapse!!!
 - Critical mass for gravitational collapse in presence of B-field

$$M_{cr} = \frac{5^{3/2}}{48\pi^2} \frac{B^3}{G^{3/2} \rho^2}$$

Critical mass-to-flux ratio
 (Mouschovias & Spitzer 1976)

$$\left[\frac{M}{\Phi}\right]_{cr} = \frac{\zeta}{3\pi} \left[\frac{5}{G}\right]^{1/2}$$

- Ambipolar diffusion can initiate collapse



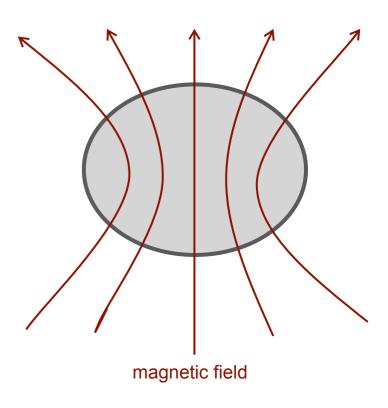
Lyman Spitzer, Jr., 1914 - 1997

"standard theory" of star formation

- BASIC ASSUMPTION: Stars form from magnetically highly subcritical cores
- Ambipolar diffusion slowly increases (M/ Φ): $\tau_{AD} \approx 10\tau_{ff}$
- Once $(M/\Phi) > (M/\Phi)_{crit}$: dynamical collapse of SIS
 - Shu (1977) collapse solution
 - $dM/dt = 0.975 c_s^3/G = const.$
- Was (in principle) only intended for isolated, low-mass stars



Frank Shu, 1943 -



problems of "standard theory"

- Observed B-fields are weak, at most marginally critical (Crutcher 1999, Bourke et al. 2001)
- Magnetic fields cannot prevent decay of turbulence (Mac Low et al. 1998, Stone et al. 1998, Padoan & Nordlund 1999)
- Structure of prestellar cores (e.g. Bacman et al. 2000, Alves et al. 2001)
- Strongly time varying dM/dt (e.g. Hendriksen et al. 1997, André et al. 2000)
- More extended infall motions than predicted by the standard model (Williams & Myers 2000, Myers et al. 2000)
- Most stars form as binaries (e.g. Lada 2006)

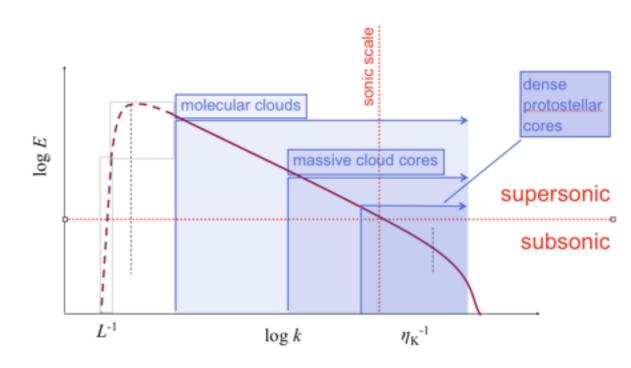
- As many prestellar cores as protostellar cores in SF regions (e.g. André et al 2002)
- Molecular cloud clumps are chemically young (Bergin & Langer 1997, Pratap et al 1997, Aikawa et al 2001)
- Stellar age distribution small ($\tau_{\rm ff}$ << $\tau_{\rm AD}$) (Ballesteros-Paredes et al. 1999, Elmegreen 2000, Hartmann 2001)
- Strong theoretical criticism of the SIS as starting condition for gravitational collapse (e.g. Whitworth et al 1996, Nakano 1998, as summarized in Klessen & Mac Low 2004)
- Standard AD-dominated theory is incompatible with observations (Crutcher et al. 2009, 2010ab, Bertram et al. 2011)

gravoturbulent star formation

BASIC ASSUMPTION:

star formation is controlled by interplay between supersonic turbulence and self-gravity

- turbulence plays a dual role:
- on large scales it provides support
- on small scales it can trigger collapse
- some predictions:
 - dynamical star formation timescale $\tau_{\rm ff}$
 - high binary fraction
 - complex spatial structure of embedded star clusters
- and many more . . .



properties of turbulence

laminar flows turn turbulent at high Reynolds numbers

$$Re = \frac{\text{advection}}{\text{dissipation}} = \frac{VL}{\nu}$$

V= typical velocity on scale L, $v = \eta/\rho$ = kinematic viscosity, turbulence for Re > 1000 \rightarrow typical values in ISM 10⁸-10¹⁰

Navier-Stokes equation (transport of momentum)

$$\rho \frac{d\vec{v}}{dt} = \rho \left(\frac{\partial \vec{v}}{\partial t} + (\vec{v} \cdot \vec{\nabla}) \vec{v} \right) = -\vec{\nabla} P + \eta \vec{\nabla}^2 \vec{v} + \left(\frac{\eta}{3} + \zeta \right) \vec{\nabla} (\vec{\nabla} \cdot \vec{v})$$

$$\text{shear viscosity}$$

$$\text{bulk viscosity}$$

$$\sigma_{ij} \equiv \eta \left(\frac{\partial v_i}{\partial x_j} + \frac{\partial v_j}{\partial x_i} - \frac{2}{3} \delta_{ij} \frac{\partial v_k}{\partial x_k} \right) + \zeta \delta_{ij} \frac{\partial v_k}{\partial x_k}$$

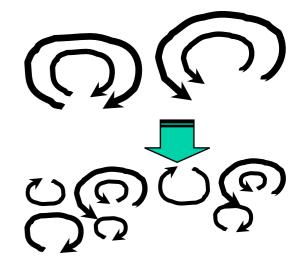
viscous stress tensor

properties of turbulence

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 vortex streching --> turbulence is intrinsically anisotropic (only on large scales you may get

homogeneity & isotropy in a statistical sense; see Landau & Lifschitz, Chandrasekhar, Taylor, etc.)

(BUT: ISM turbulence: shocks & B-field cause additional inhomogeneity)



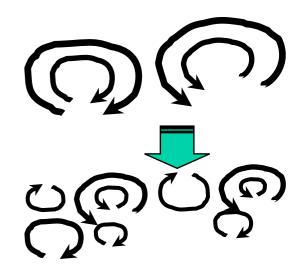
Tornado over Portofino

properties of turbulence

laminar flows turn *turbulent* at *high* Reynolds numbers

$$Re = \frac{\text{advection}}{\text{dissipation}} = \frac{VL}{\nu}$$

V= typical velocity on scale L, $\nu = \eta/\rho$ = kinematic viscosity, turbulence for Re > 1000 → typical values in ISM 108-1010

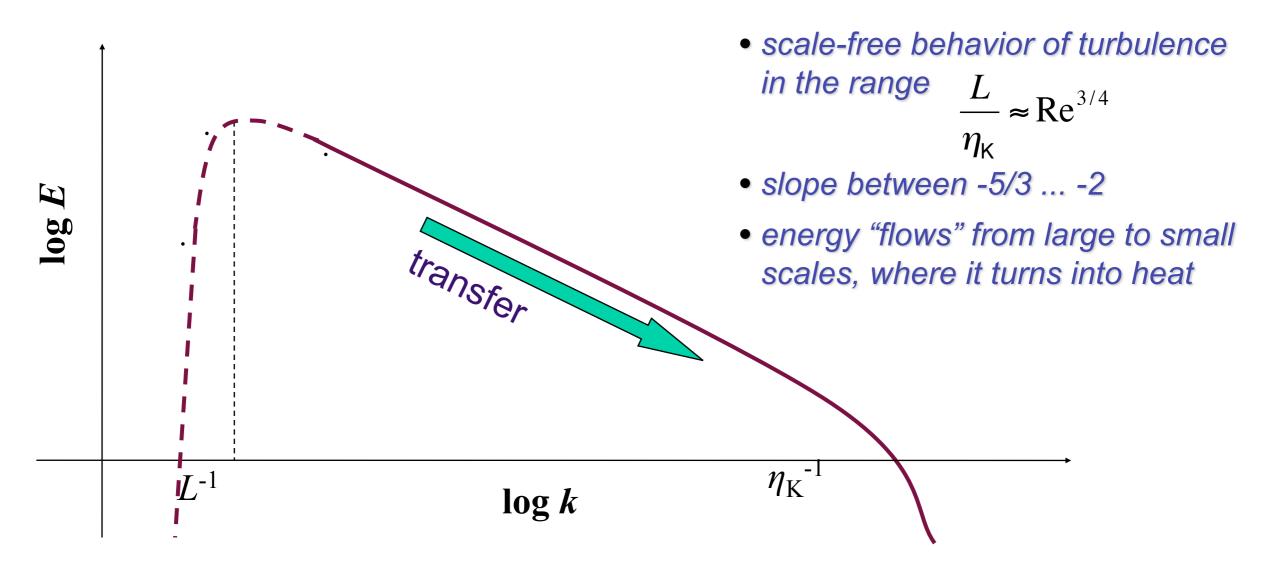


 vortex streching --> turbulence is intrinsically anisotropic (only on large scales you may get homogeneity & isotropy in a statistical sense; see Landau & Lifschitz, Chandrasekhar, Taylor, etc.)

(BUT: ISM turbulence: shocks & B-field cause additional inhomogeneity)



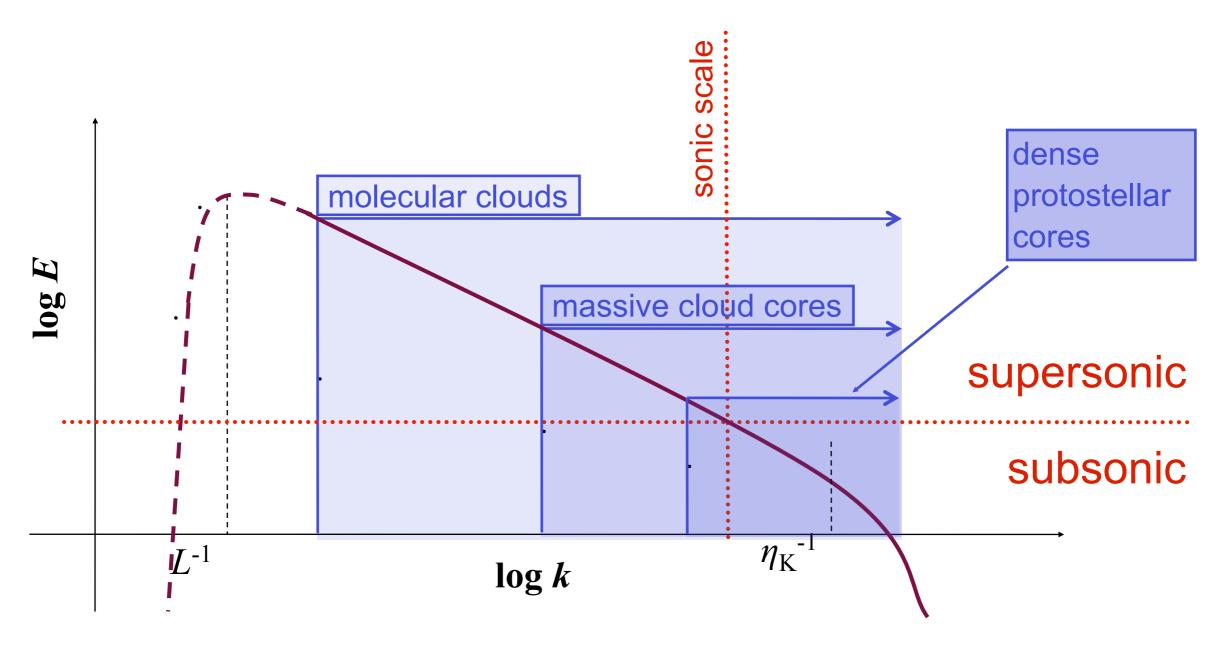
turbulent cascade in the ISM



energy source & scale NOT known (supernovae, winds, spiral density waves?)

dissipation scale not known (ambipolar diffusion, molecular diffusion?)

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$$\sigma_{rms} << 1 \text{ km/s}$$

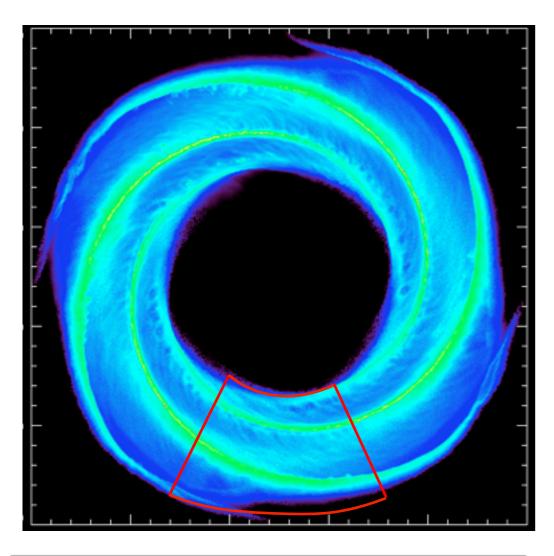
$$M_{rms} \le 1$$

$$L \approx 0.1 \text{ pc}$$

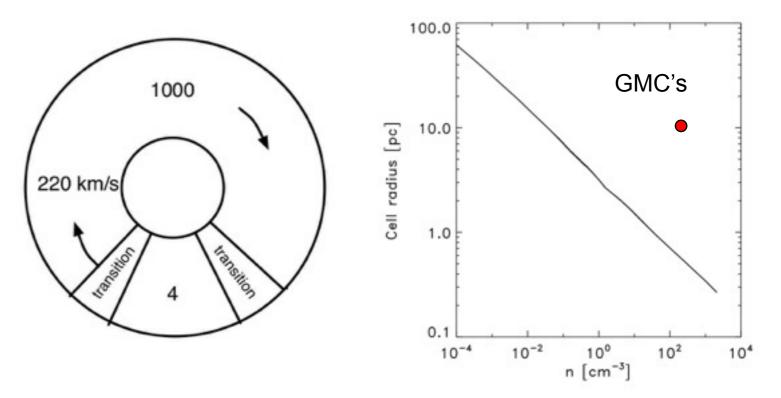
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modeling molecular cloud formation

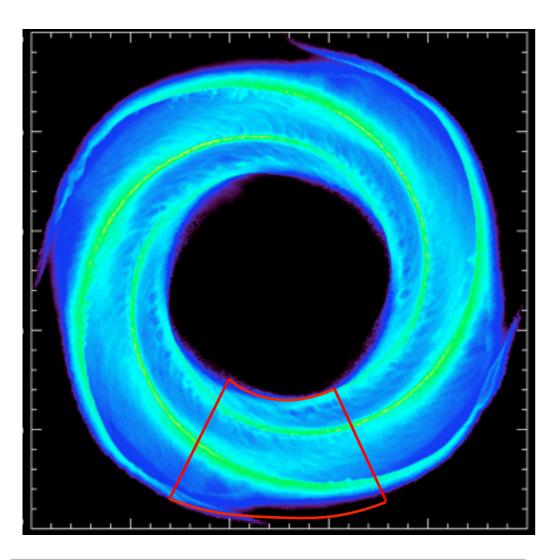


Simulation	Surface Density ${\rm M}_{\odot}~{\rm pc}^{-2}$	Radiation Field G_0
Milky Way	10	1
Low Density	4	1
Strong Field	10	10
Low & Weak	4	0.1

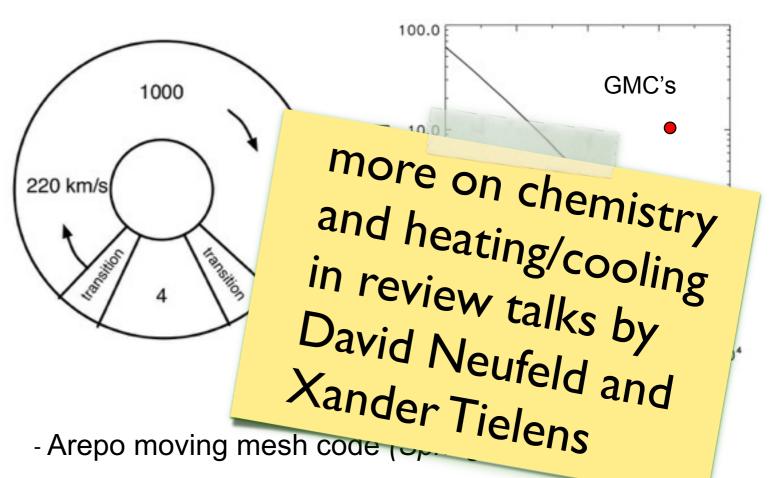


- Arepo moving mesh code (Springel 2010)
- time dependent chemistry (Glover et al. 2007) gives heating & cooling in a 2 phase medium
- two layers of refinement with mass resolution down to 4 M_☉ in full Galaxy simulation
- UV field and cosmic rays
- TreeCol (Clark et al. 2012)
- external spiral potential (Dobbs & Bonnell 2006)

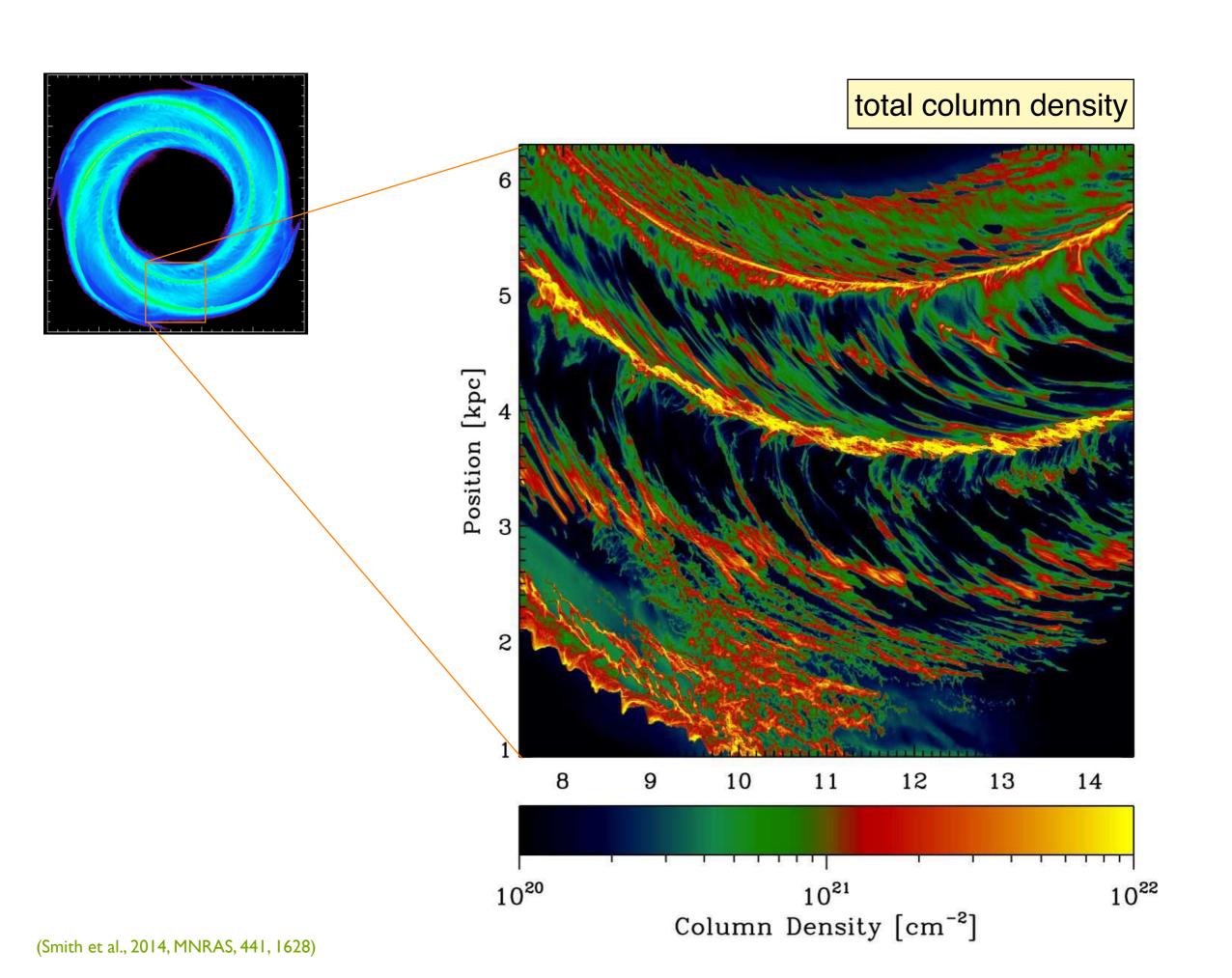
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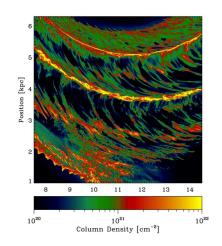


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HI column density

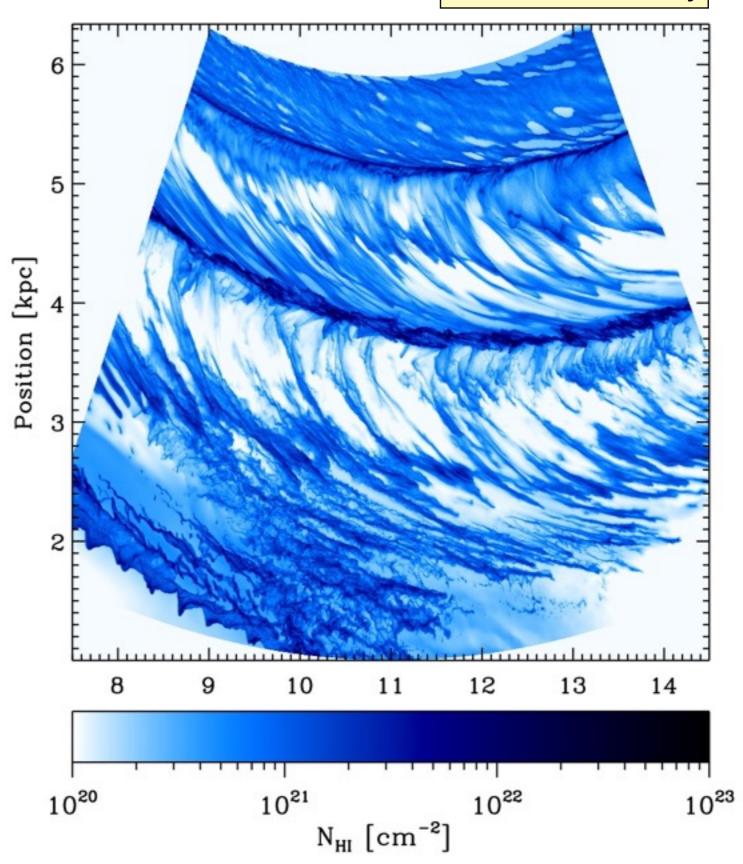
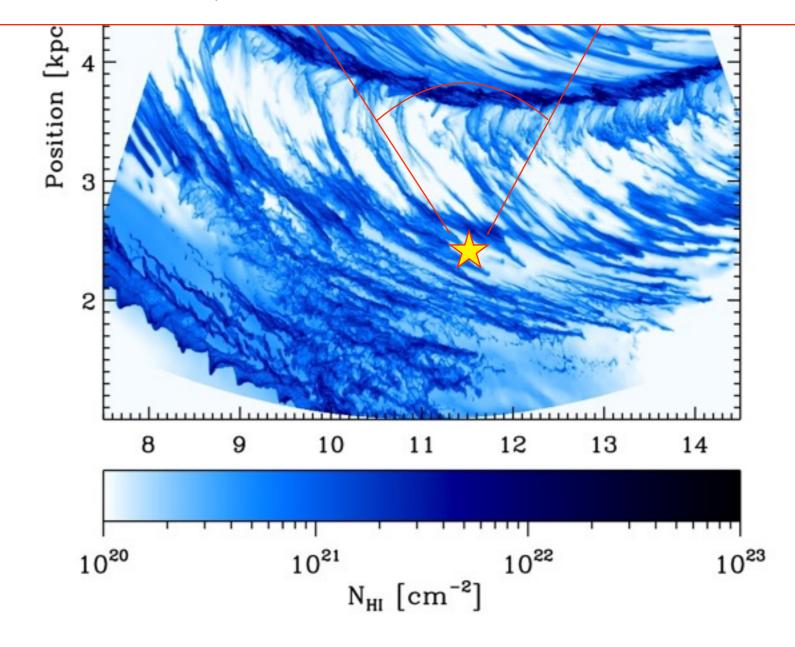
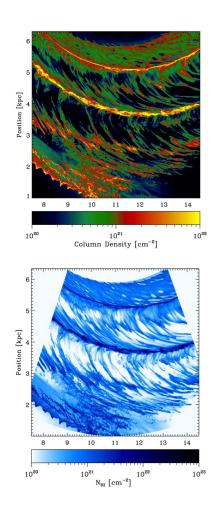


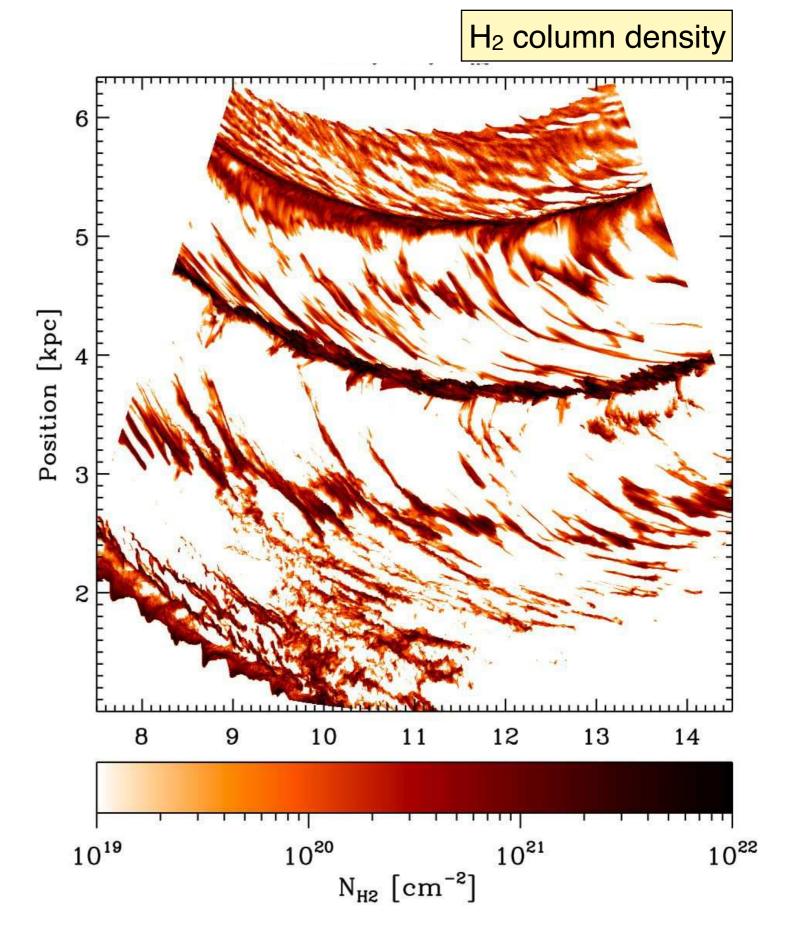
image from THOR Galactic plane survey (PI H. Beuther): continuum emission around 21 cm

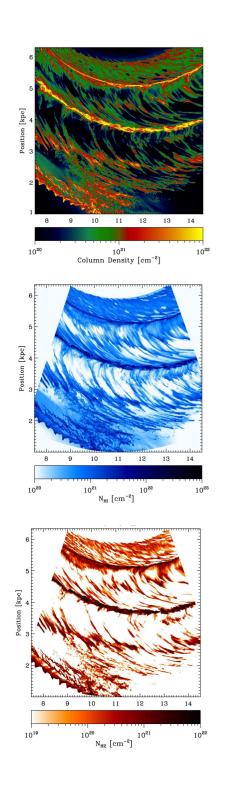
next step: produce all sky maps at various positions in the model galaxy (use RADMC-3D)

(Beuther et al., 2016, A&A, in press, arXiv:1609.03329, Bihr et al. 2016, A&A, 588, A97)

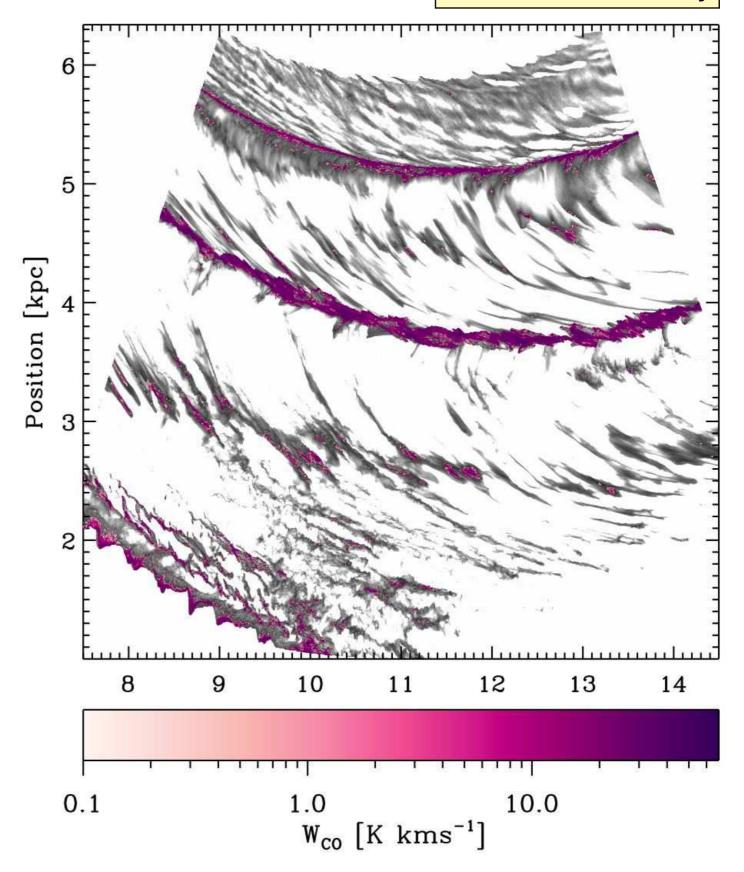


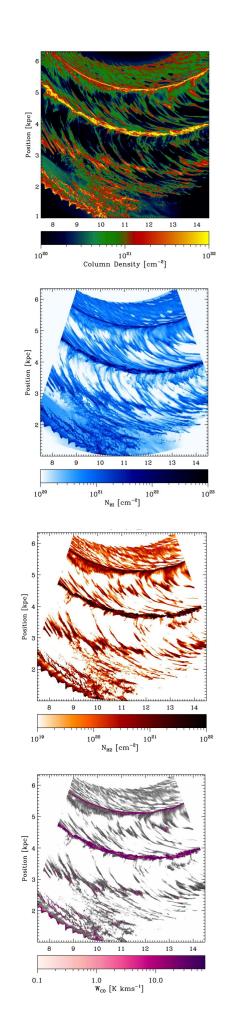


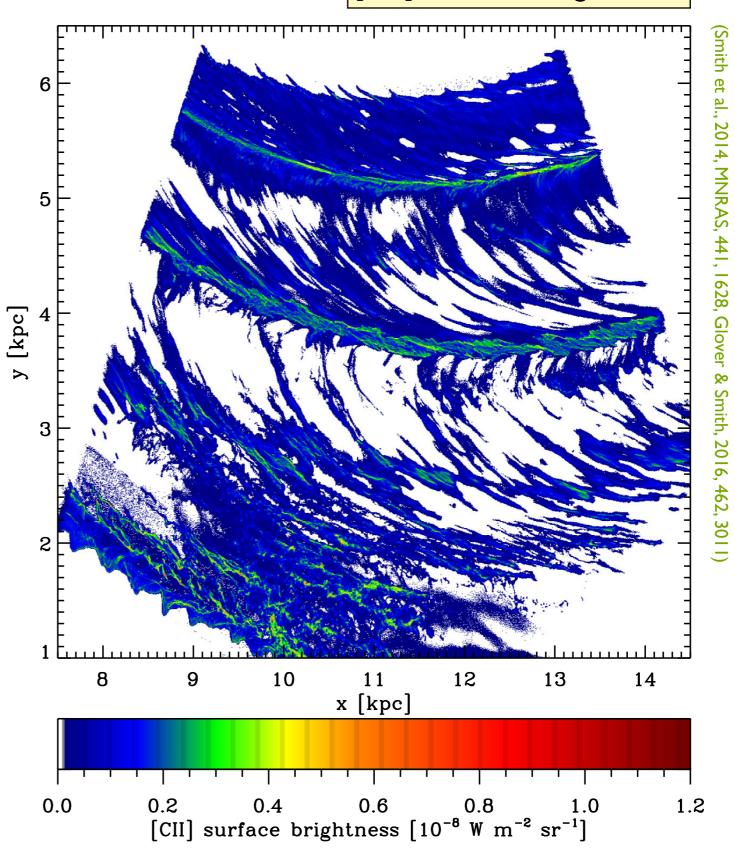




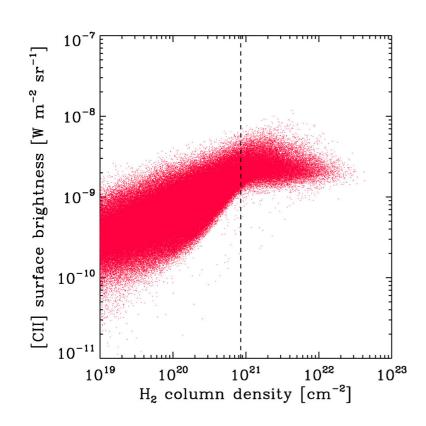
CO column density

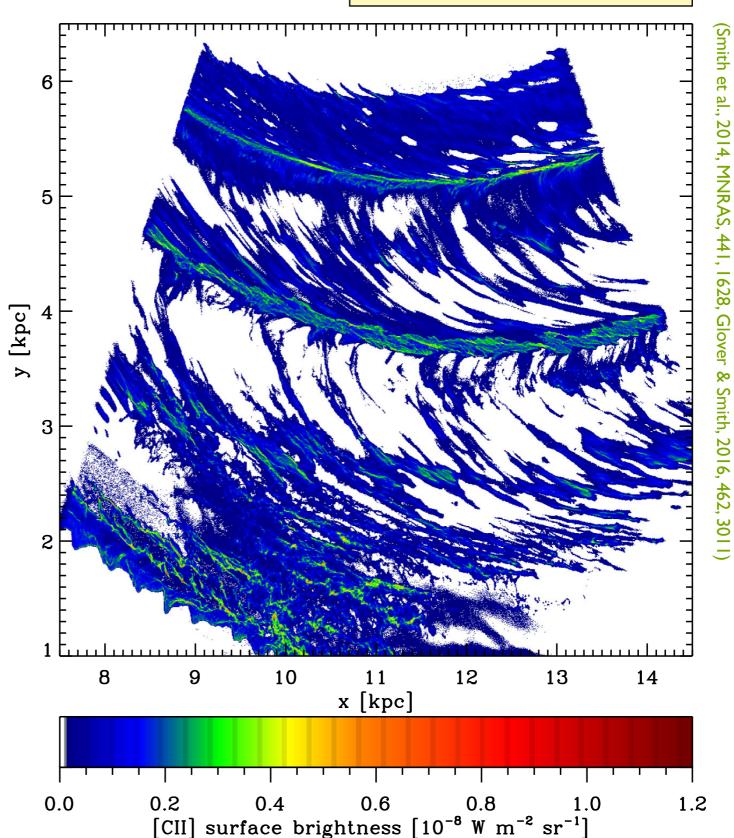


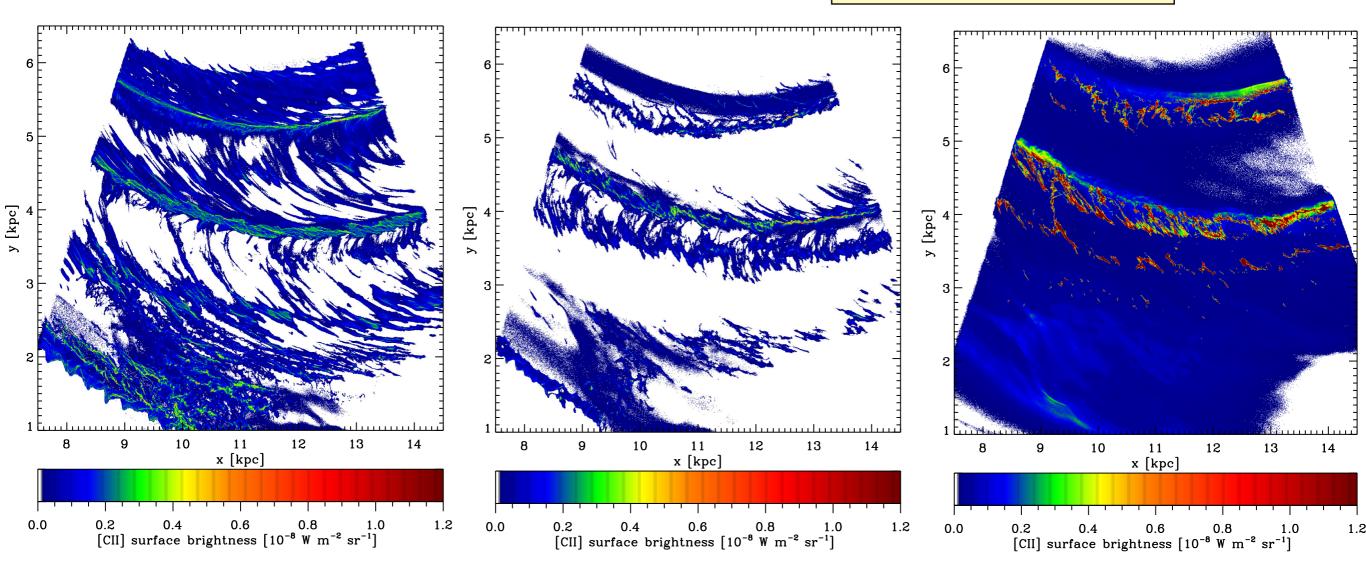




- weak correlation between [CII] emission and H₂ column density (saturation at large columns)
- CO-bright component is cold (T ≤ 30 K) and gas is almost 100% molecular (clouds)
- CO-dark gas has range of temperatures (30 K ≤ T ≤ 100 K), H₂ fraction varies strongly



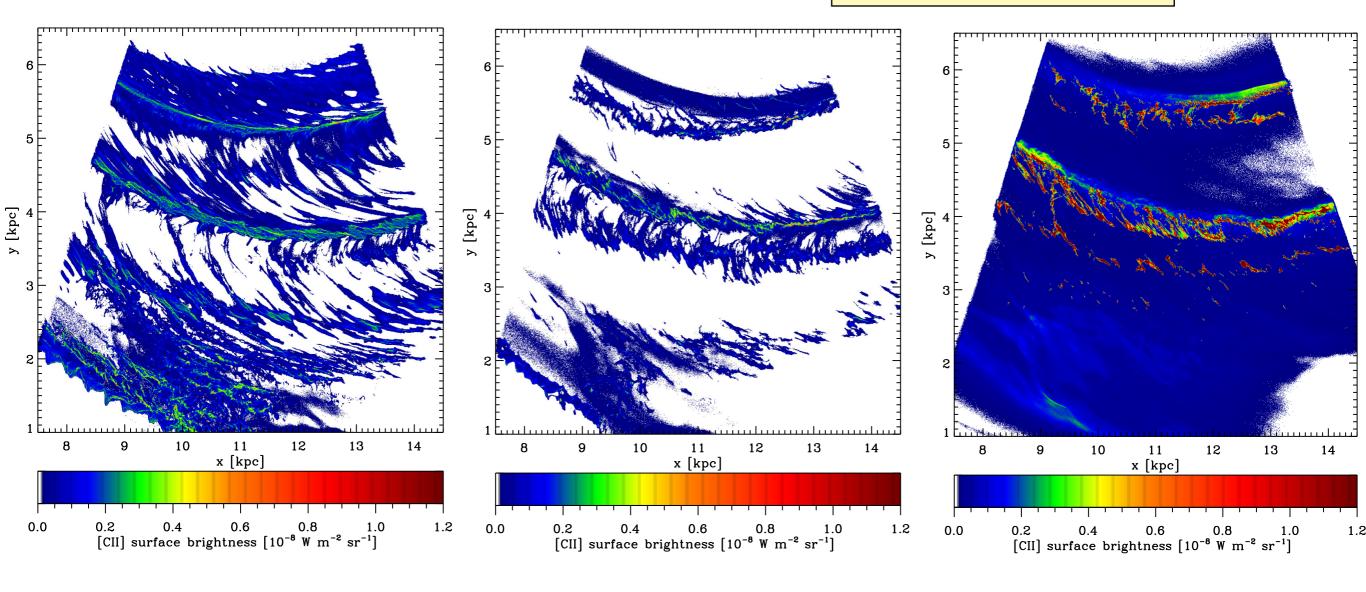




standard MW case

low surface density

standard surface density with high radiation field $(G_0 = 17)$



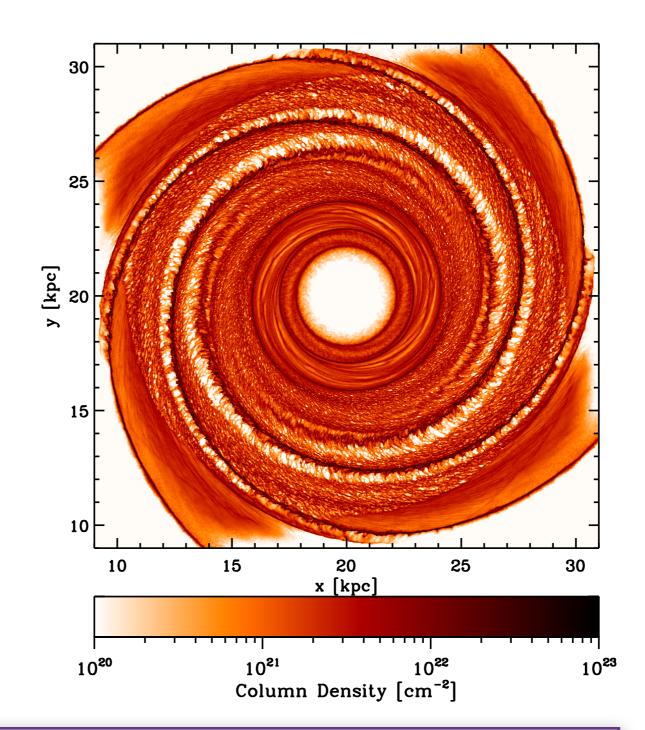
comparison with data from SOFIA large program on M51

surface density radiation field

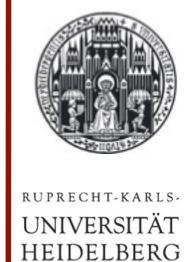
new models (full disk)

same as in Smith et al. (2014), but improved:

- more realistic potential (better disk scale height)
- larger disk area
- with self-gravity and supernovae feedback!
- two types now:
 - high resolution (4 M_☉)
 wedge as previously
 - new runs with 200 M_☉ resolution everywhere



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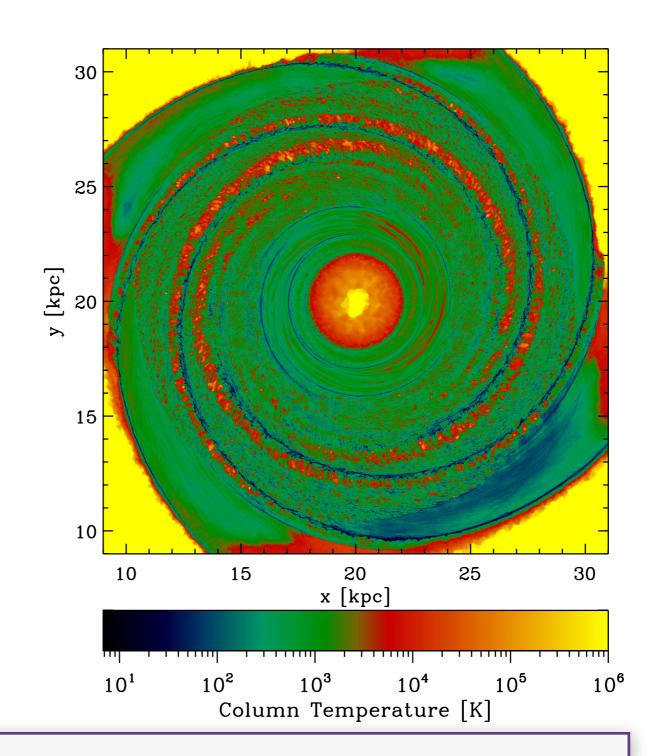


Ralf Klessen 17.10.2016

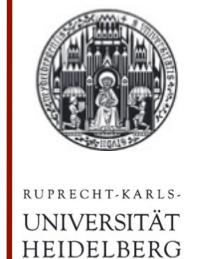
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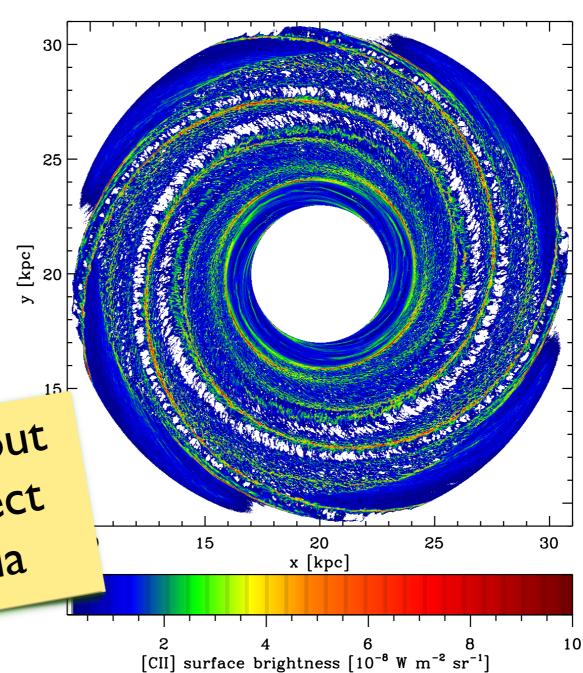


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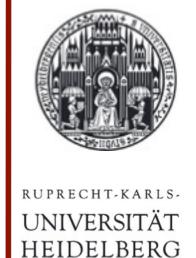
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- more information about the M51 SOFIA project in talk by Jorge Pineda

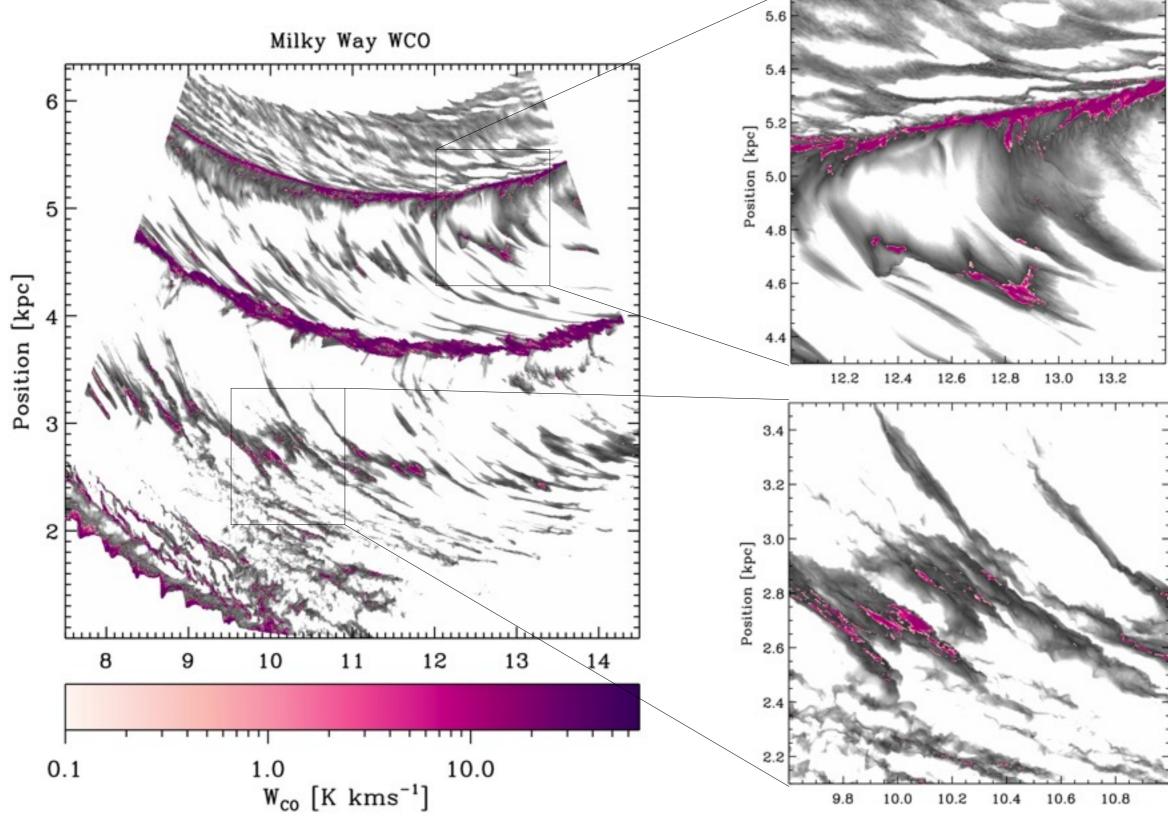


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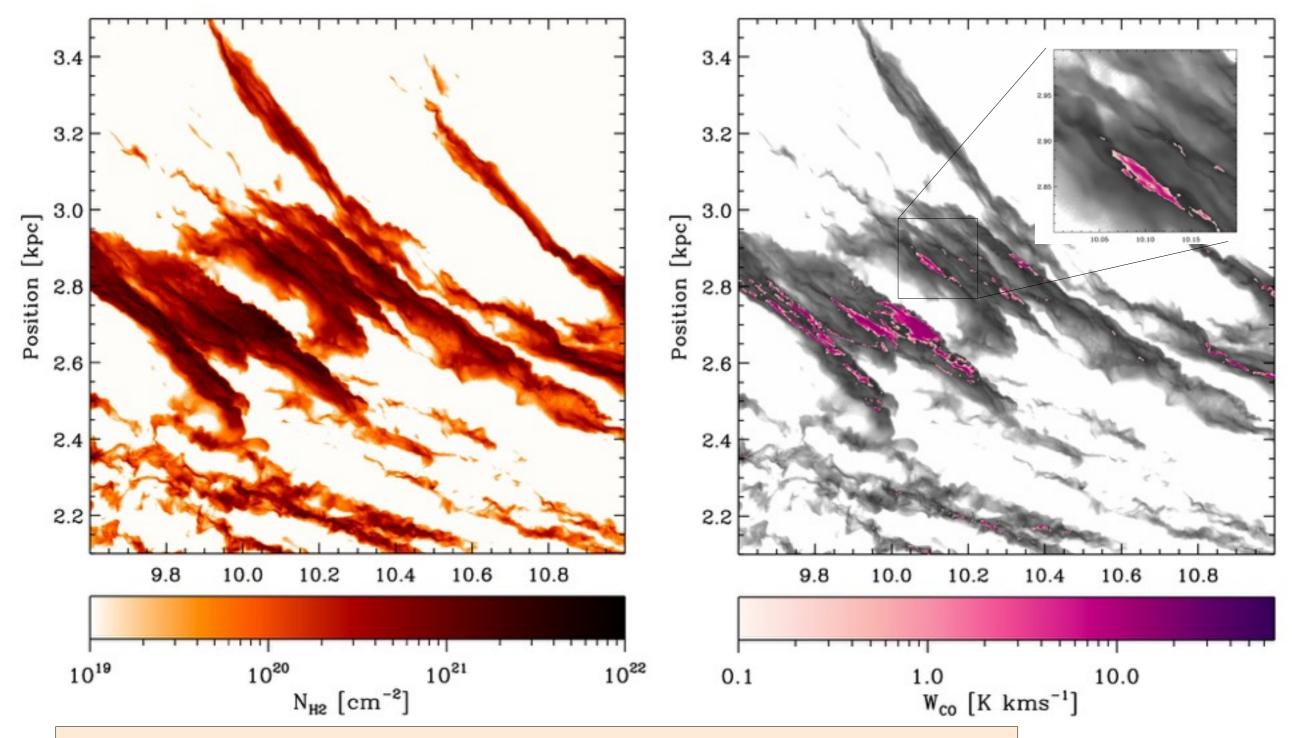


Ralf Klessen 17.10.2016

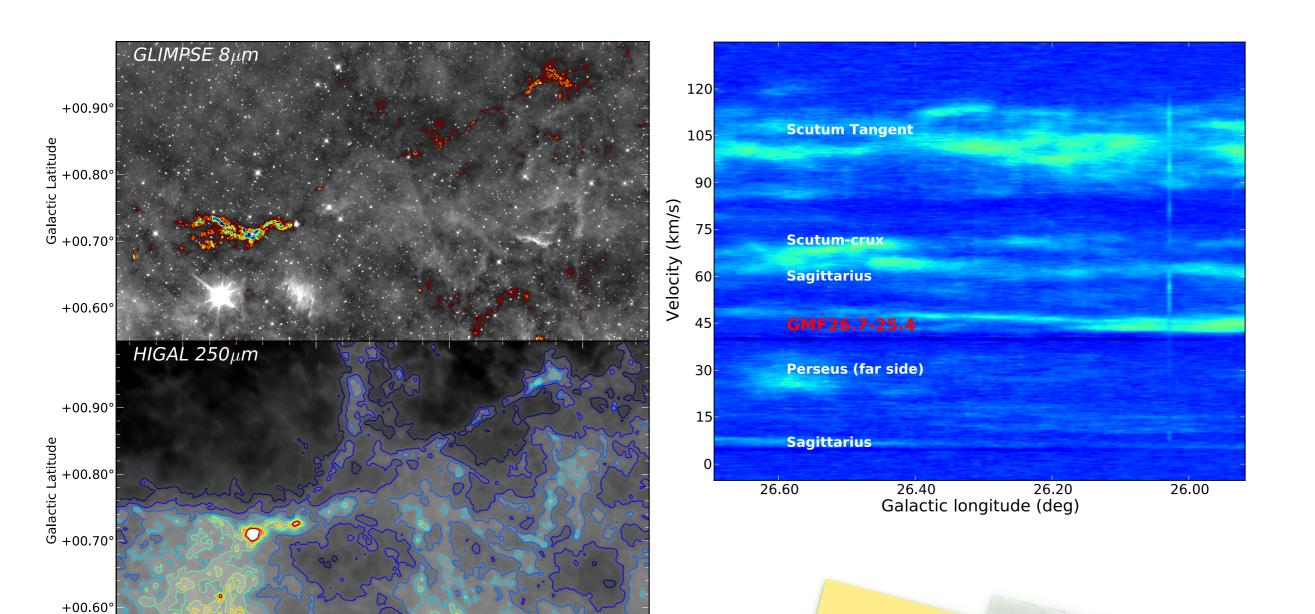
details of CO emission



relation between H₂ and CO



Filamentary molecular clouds in inter-arm regions are likely only the observable parts of much larger structures.



Filamentary molecular clouds in inter-arm regions are likely only the observable parts of much larger structures.

26.40°

26.20°

Galactic Longitude

26.00°

26.60°

26.00°		m	ore	On.	CI			
Name	Cloud mass [M _G	4.8e+4 1.3e+4 3.7e+4 ^a 5.0e+3 ^a	ks by	Sa	rilam rah	ent	s in	
GMF 18.0-16.8 GMF 20.0-17.9	1.5e+5 4.0e+5	4.8e+4	Edit	th F	alga	Mag	an	rm
GMF 26.7-25.4 GMF 38.1-32.4a	2.0e+5 7.0e+5	1.3e+4 3.7e+4 ^a	6.5 5.3	5.9	"हव।	on	е	
GMF 38.1-32.4b GMF 41.0-41.3 GMF 54.0-52.0	7.7e+4 4.9e+4 6.8e+4	5.0e+3 ^a 7.7e+2 2.4e+3	6.5 1.6 3.5	6.2 6.5 7.3	11.5 12.5 11.2	19 25	W 52	
Nessie G32.02+0.06	- 2.0e+5	3.9e+5 3.0e+4	15.0	5.6 4.7	-7.8 19.7	- 48	SC-arm	

dark gas fraction

Observational estimates:

Grenier et al. (2005) $f_{DG} = 0.33-0.5$

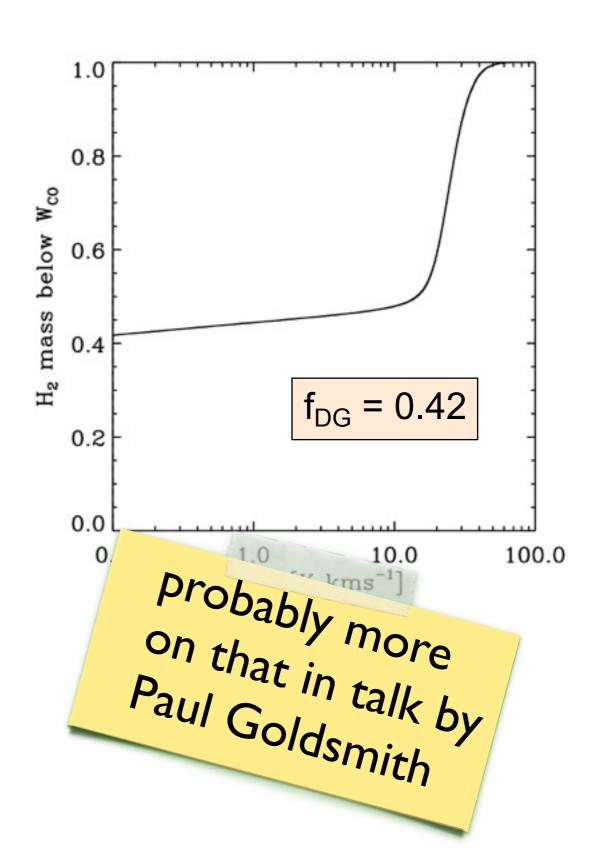
Planck coll. (2011)* $f_{DG} = 0.54$

Paradis et al. $(2012)^*$ $f_{DG} = 0.62$

(inner $f_{DG} = 0.71$, outer $f_{DG} = 0.43$)

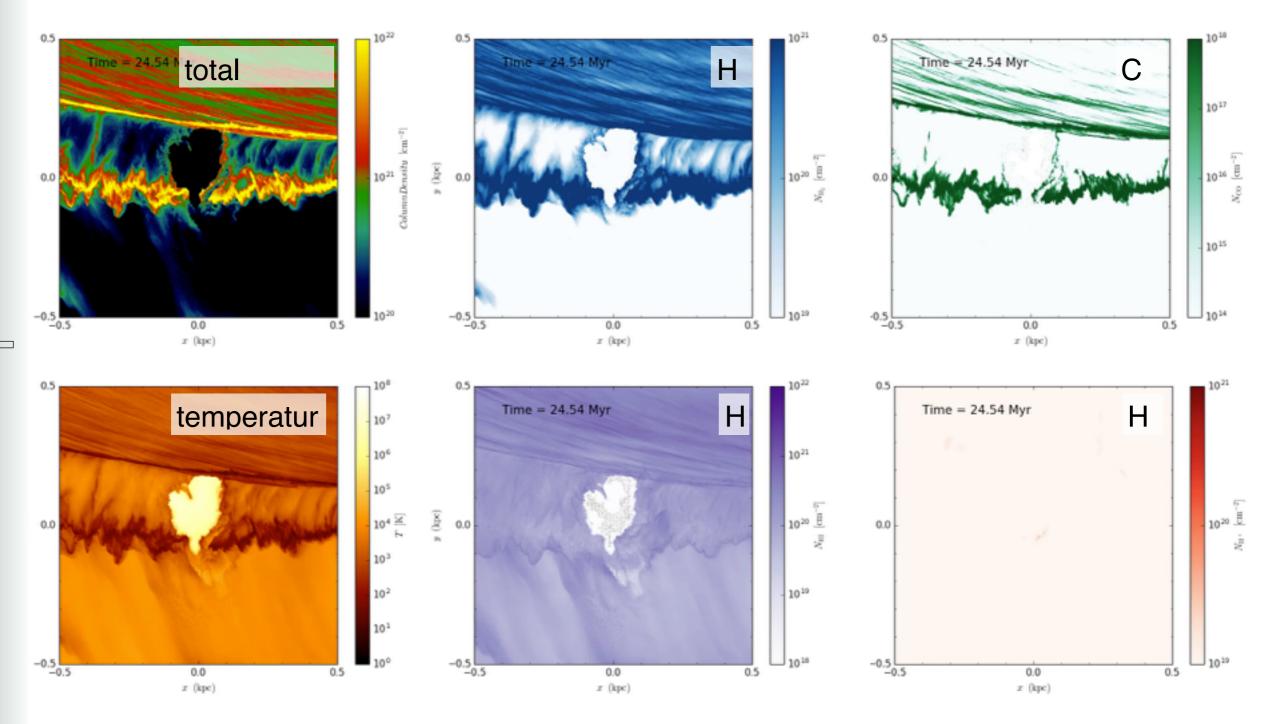
Pineda et al. (2013) $f_{DG} = 0.3$

Roman-Duval et al. $f_{DG} \sim 0.5$ (in prep.)

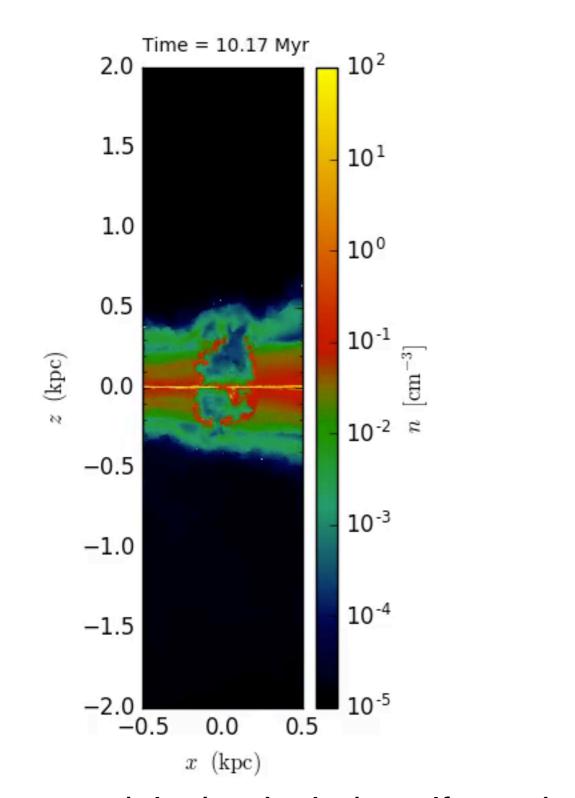


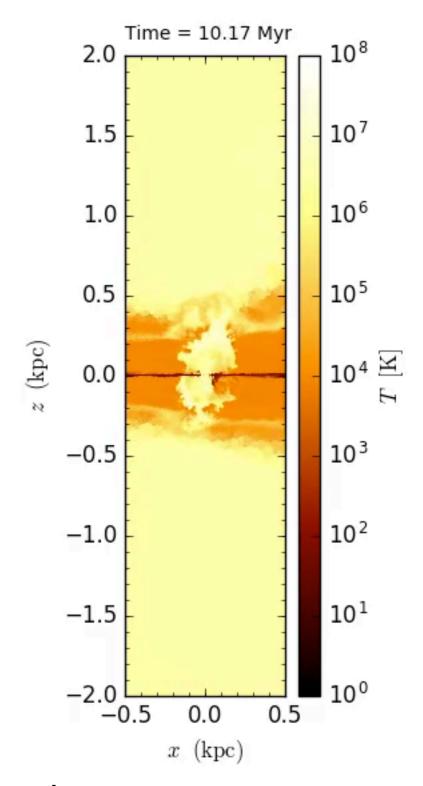
^{*} dust methods have large uncertainties.



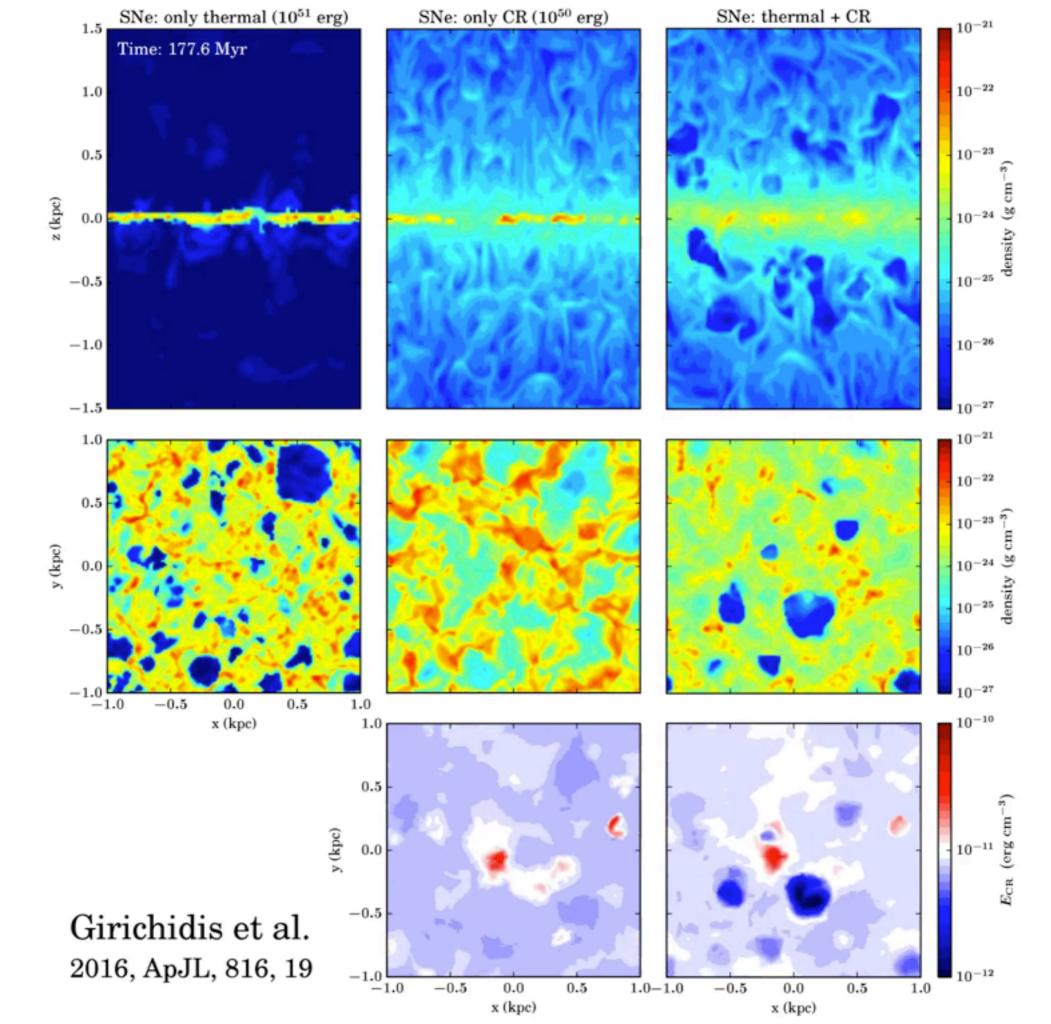


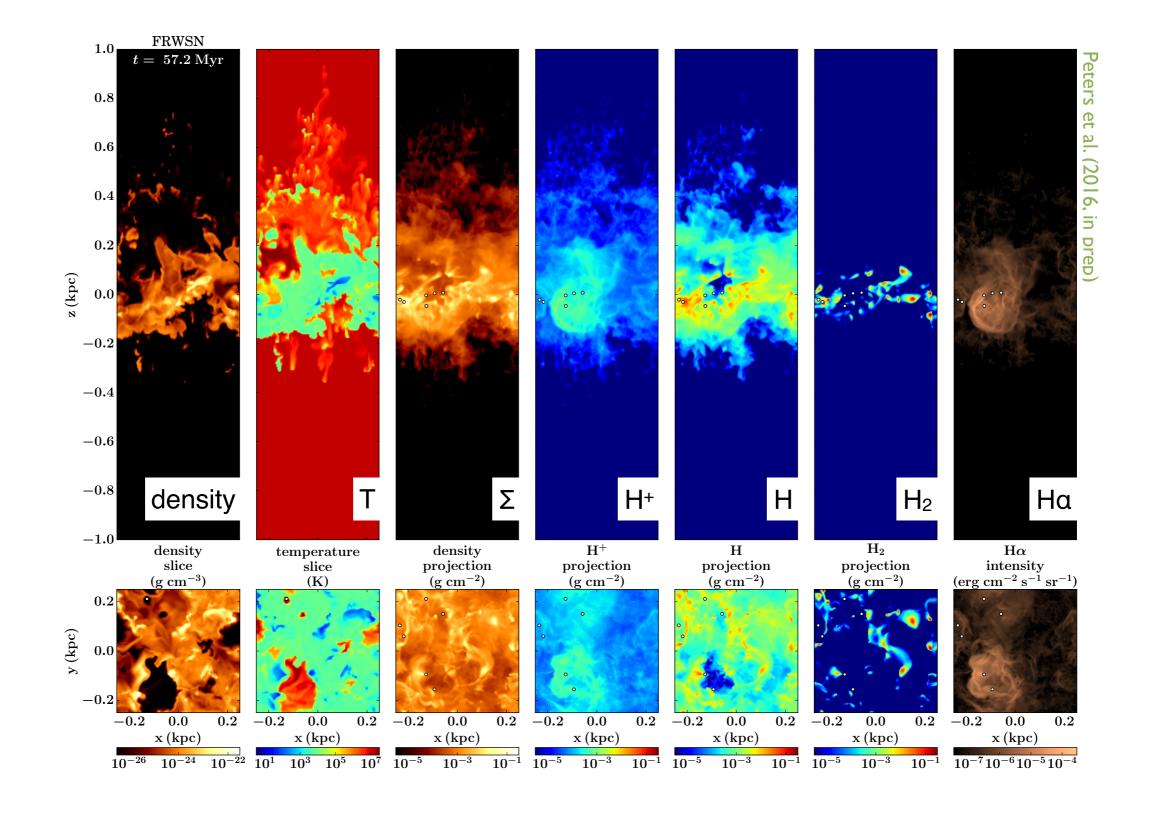
new models that include self-consistent star formation





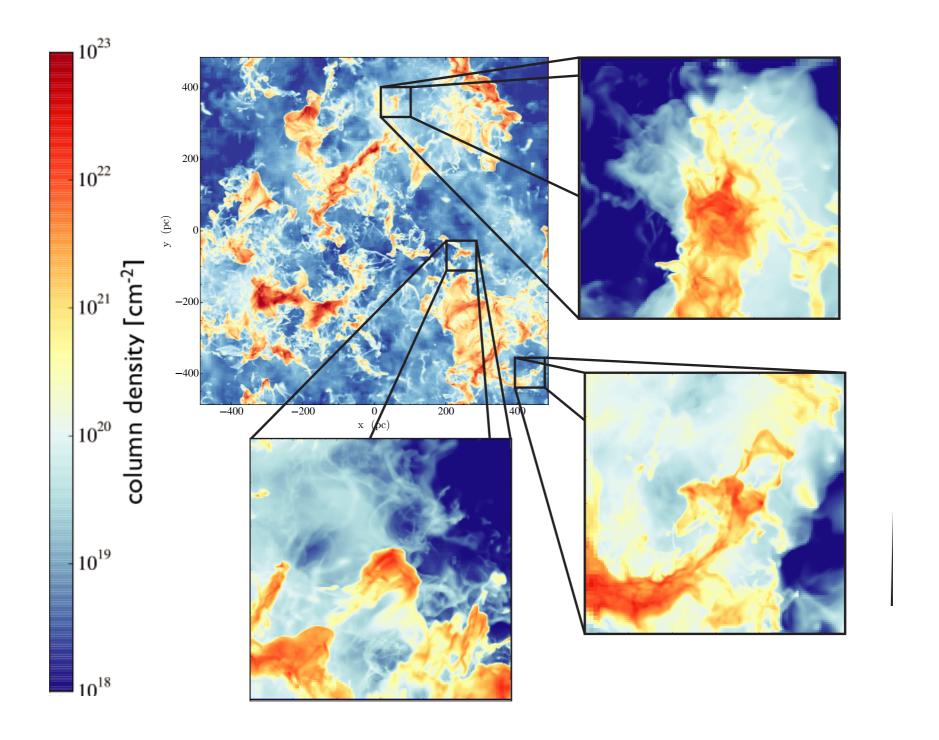
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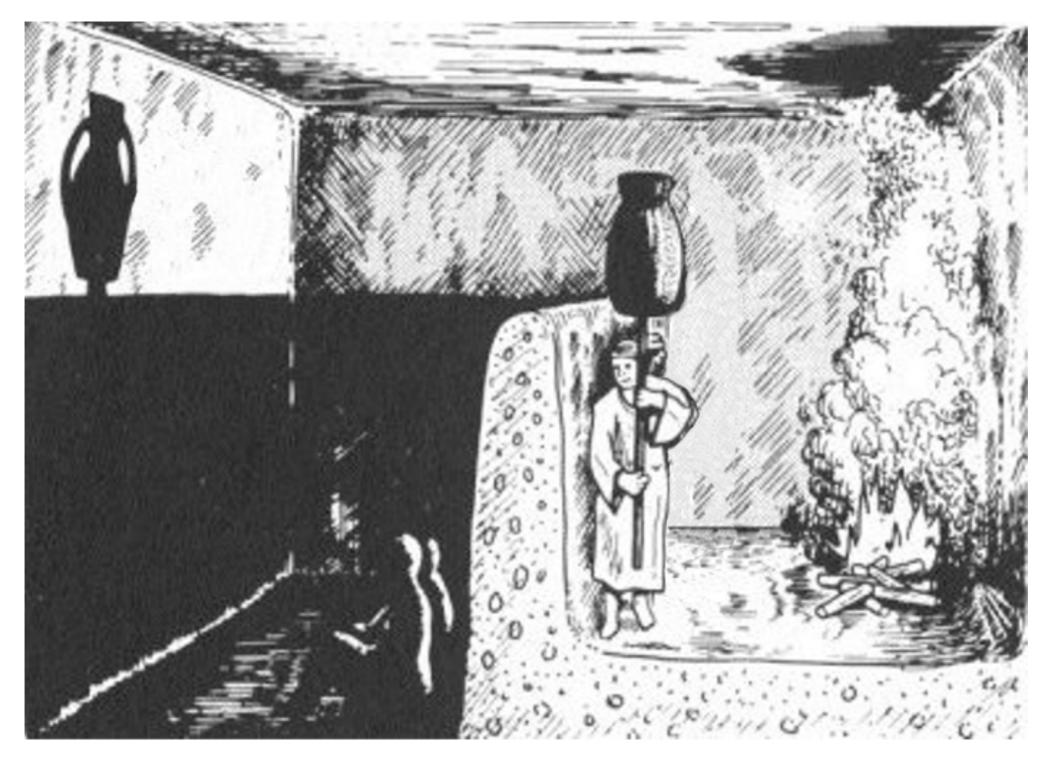
synthetic maps in further observables: HI, Halpha, other radio recombination lines

SILCC collaboration: http://hera.ph1.uni-koeln.de/~silcc/



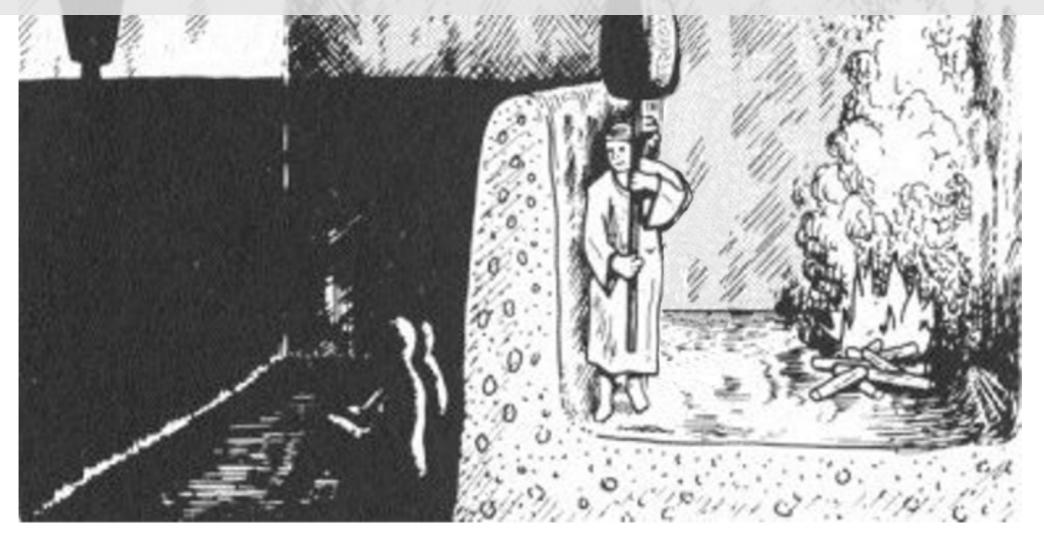
zoom-in calculations to provide better boundary conditions for star cluster formation simulations





^{*} The Republic (514a-520a)

 ISM dynamics and star formation is governed by complex interplay of self-gravity and a large number of competing processes (such as turbulence, B-field, feedback, thermal pressure, CR pressure, and other processes)



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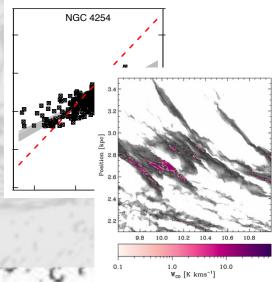
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MARCHANTIAN CONTROL

- example 1: hierarchical Bayesian statistics indicated complex relation between ISM properties and star formation on galactic scales (Kennicutt Schmidt relation)
- example 2: detailed (M)HD calculations with time-dependent chemistry allow us to study the properties of CO-dark H2 gas



- (personal) WISH LIST for SOFIA: help quantifying amount of CO-dark ISI sel turk
- H₂ in the Milky Way and in other othe identify and characterize convergent therr cooli
- flows in turbulent ISM that form closing genera
- —> help determining initial conditions molecular clouds example relation l on galact
 - for star (cluster) formation
- example 2 arations with time-dependent chemistry to study the properties of CO-dark H2 gas CONTRACTOR CONTROL

and c

